



Funded by the Horizon 2020 Framework Programme  
of the European Union

## Real space imaging of the superconducting vortex lattice: Recent results and prospects

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Laboratorio de Bajas Temperaturas  
Departamento de Física de la Materia Condensada  
Instituto de Ciencia de Materiales Nicolás Cabrera

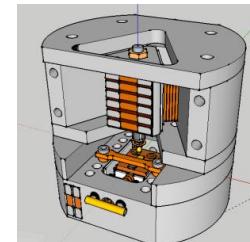
Center for Condensed Matter Physics (IFIMAC)

Universidad Autónoma de Madrid (UAM)

## Real space imaging of the superconducting vortex lattice: Recent results and prospects

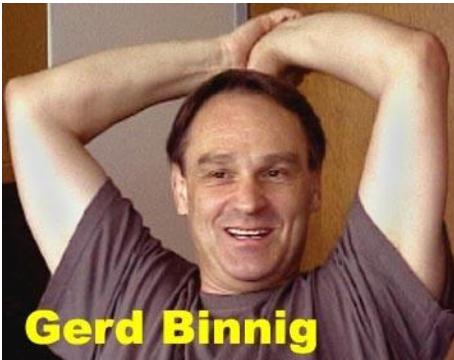


- Real space imaging and spectroscopy of the vortex lattice
- Superconducting CaKFe<sub>4</sub>As<sub>4</sub>, bandstructure, vortex core, vortex lattice and pinning
- Vortex lattice structure: interactions vs disorder
- Observation of vortex creep on cooling
- Josephson effect at atomic scale



## The birth of scanning probe microscopy

The invention of the scanning tunneling microscope  
Gerd Binnig and Heinrich Rohrer, Zurich, 1982



**Gerd Binnig**



**Heinrich Rohrer**

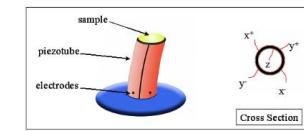
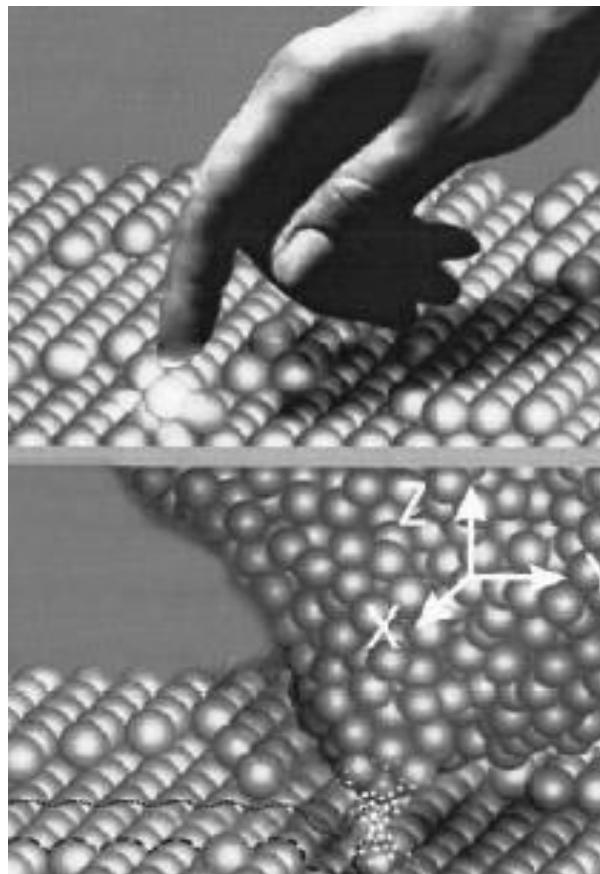


Figure 2.5. A piezotube scheme. The cross section drawing shows the inner electrode (Z) and the four outer electrodes (+X, -X, +Y, -Y).

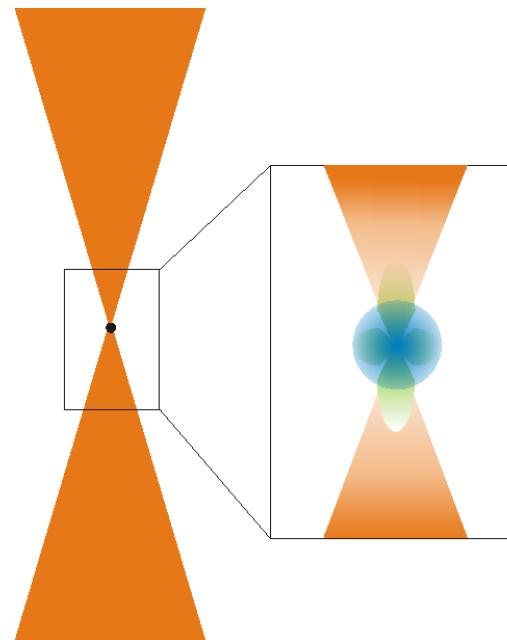
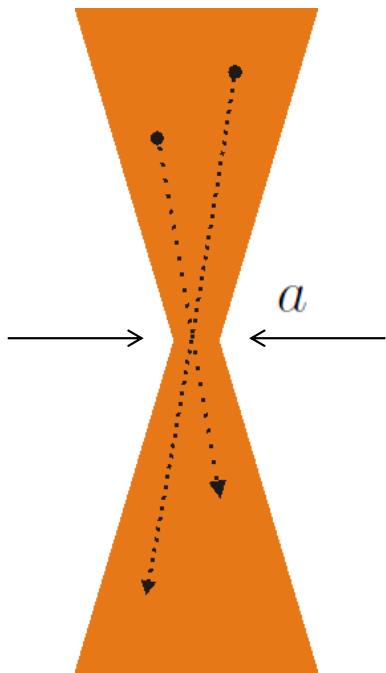
$$\Delta x, y \approx 20 \frac{A}{V}$$

$1\text{\AA} \sim 250\text{ mV}$

Tunneling  
Piezoelectrics

## Tunneling

Conduction through constrictions



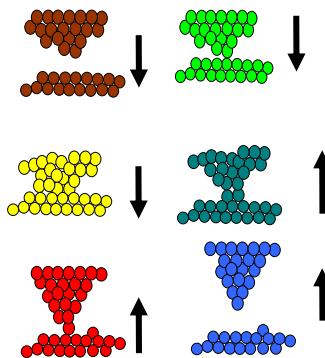
$$G = \frac{2e^2}{h} \left( \frac{k_F a}{2} \right)^2$$

$$G = \frac{2e^2}{h} \sum_i T_i$$

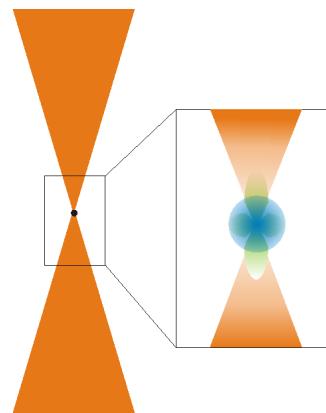
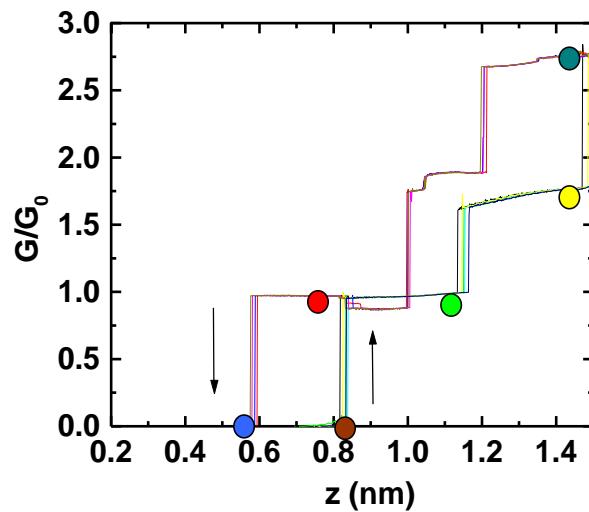
Imaging superconducting vortex cores and lattices with a STM  
H. Suderow, I. Guillamón, J.G. Rodrigo and S. Vieira  
Superc. Sci. Techn., 27, 063001 (2014)

## Tunneling

Conduction through constrictions



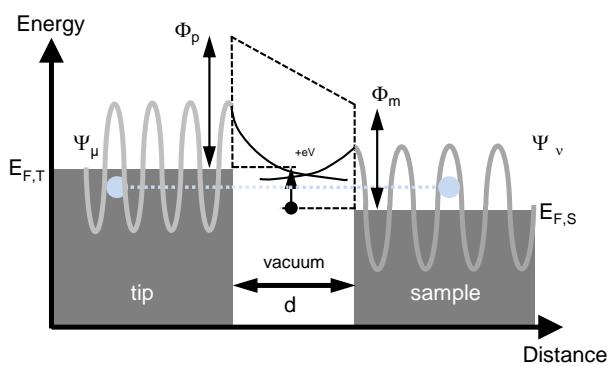
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## Tunneling

The tunneling current: Bardeen's formalism



Tunnelling from a Many-Particle Point of View  
 J. Bardeen  
 PRL, 6, 57 (1961)

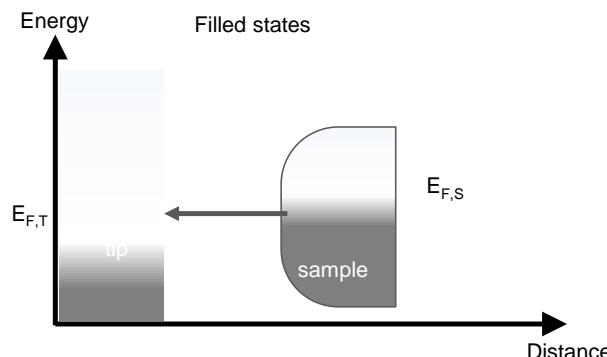
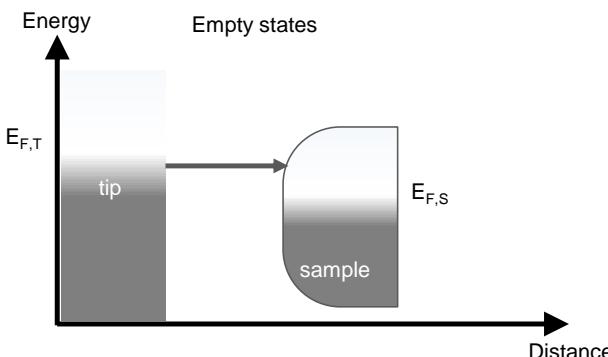
$$I = \frac{2\pi e}{\hbar} \sum_{\mu\nu} f(E_\mu) (1 - f(E_\nu + eV)) |M_{\mu\nu}|^2 \delta(E_\mu - E_\nu)$$

$$f(E) = \frac{1}{1 + e^{\frac{(E-E_F)}{k_B T}}}$$

$$E_\mu; \Psi_\mu$$

$$E_\nu; \Psi_\nu$$

$$M_{\mu\nu} = \frac{\hbar^2}{2m} \int dS (\Psi_\mu^* \nabla \Psi_\nu - \Psi_\nu \nabla \Psi_\mu^*)$$

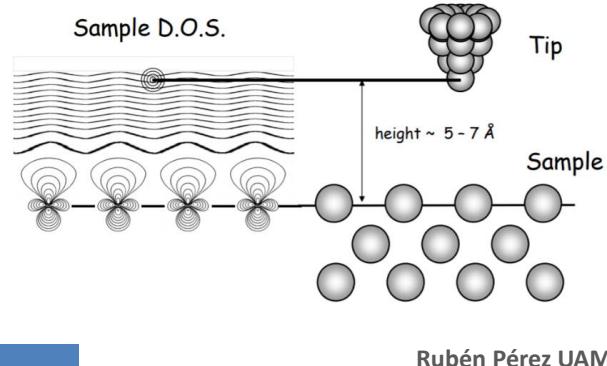
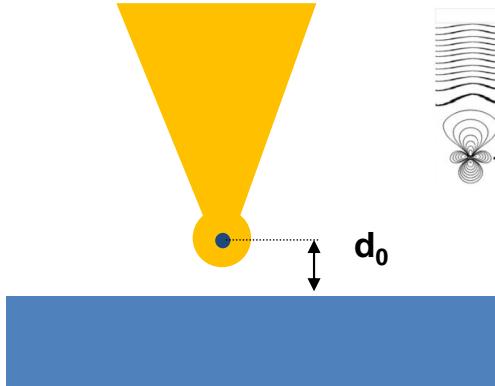


## Tunneling

### The tunneling current: Tersoff-Haman

*Theory and Application for the Scanning Tunneling Microscope*  
 J. Tersoff and D. R. Hamann  
 PRL, 50, 1998 (1983)

Ideal tip, low bias voltage ( $V=0$ ) :



Rubén Pérez UAM

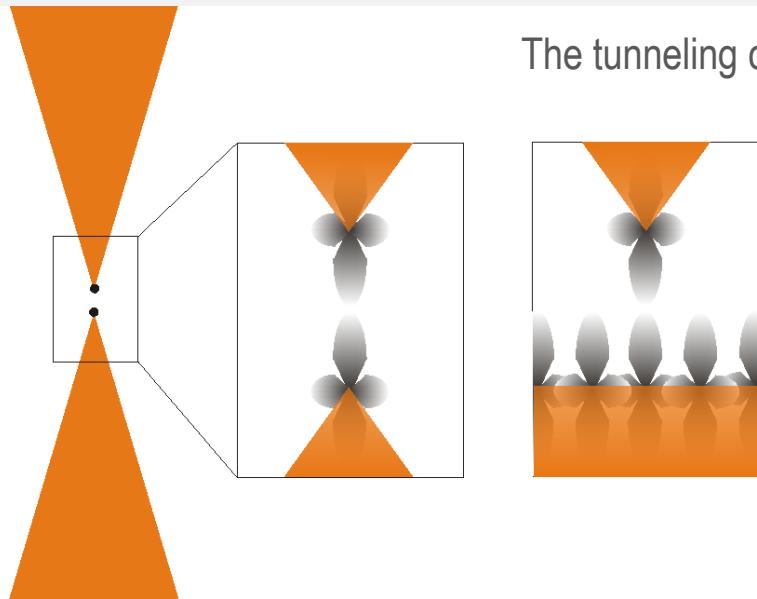
$$I \propto \sum_{\nu} |\psi_{\nu}(\mathbf{d}_0)|^2 \delta(E_{\nu} - E_F) = \rho(\mathbf{d}_0, E_F)$$

STM: Contour map of the local DOS

$$I(r_0, V) \propto \int_{E_F}^{E_F + eV} dE \rho_s(\mathbf{d}_0, E) T(\mathbf{d}_0, E, eV)$$

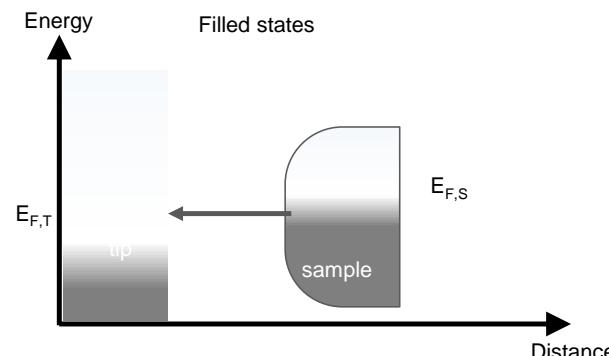
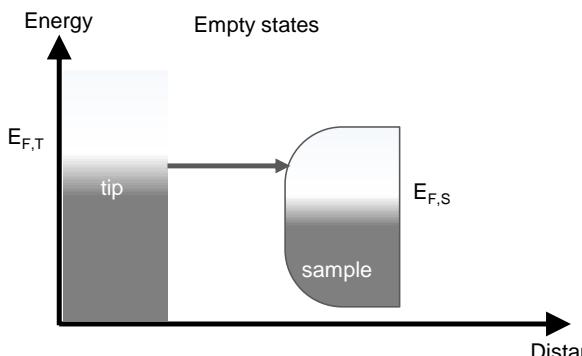
## Tunneling

The tunneling current: Chen    Directional tunneling



Microscopic ( $C_p, \kappa, \rho, \dots$ )  
Empty + filled states (ARPES, dHvA)  
Bandstructure (ARPES)  
Real space (neutrons, x-ray)

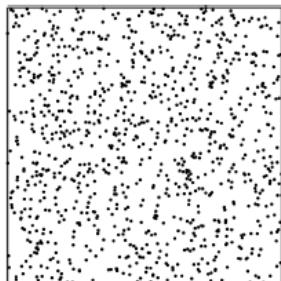
$$I(r_0, V) \propto \int_{E_F}^{E_F + eV} dE \rho_s(\mathbf{d}_0, E) T(\mathbf{d}_0, E, eV)$$



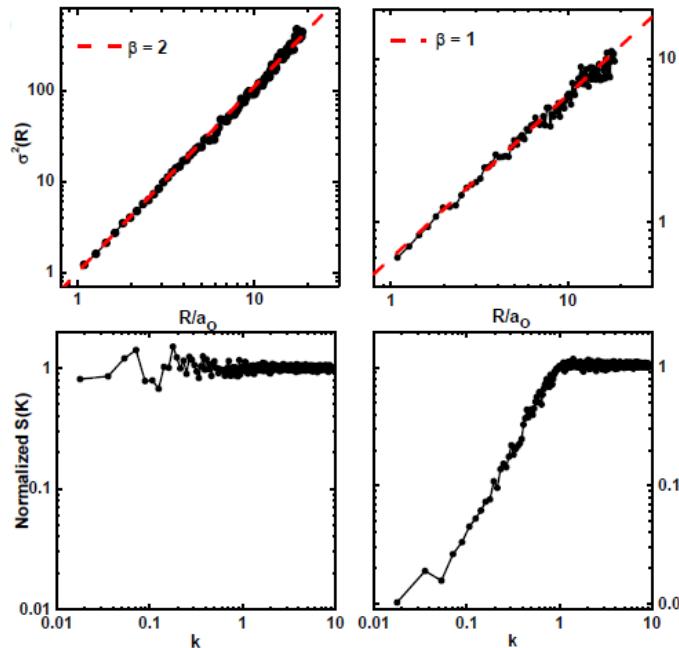
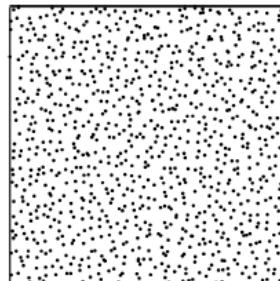
## Imaging

Working in real space

Random



Hyperuniform



Vortex lattice ?

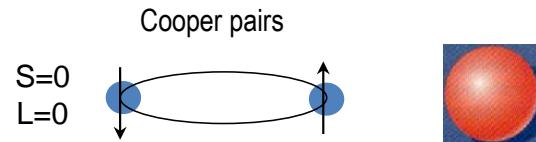
## Tunneling

### Superconductor

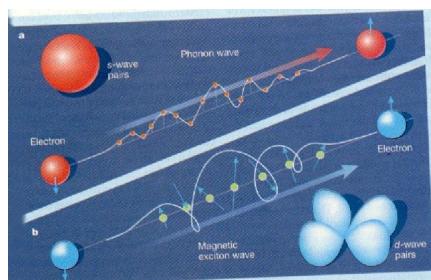
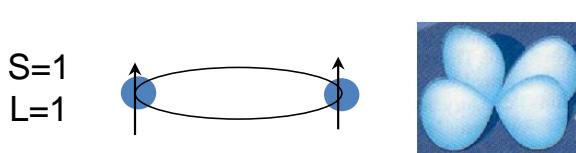


**Microscopic theory**  
 J. Bardeen, L. Cooper, J. Schrieffer (BCS, 1957)

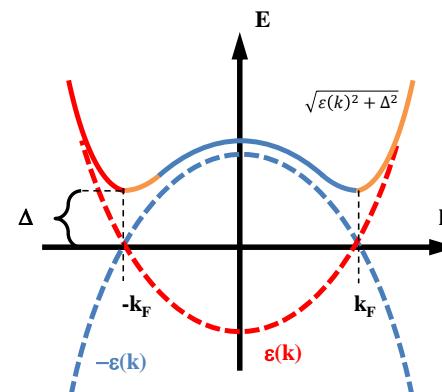
$$\psi = |\psi| e^{i\varphi}$$



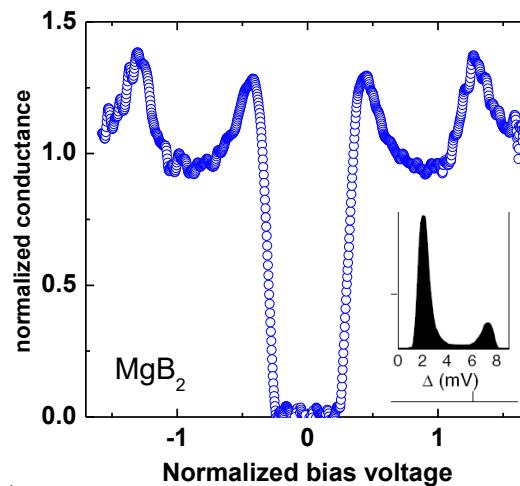
Kohn, Luttinger (1965)



$$\Delta_k = - \sum_{k'} V_{kk'} \frac{\Delta_{k'}}{2E_{k'}}$$



## STM Tunneling



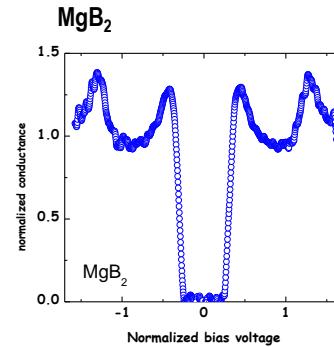
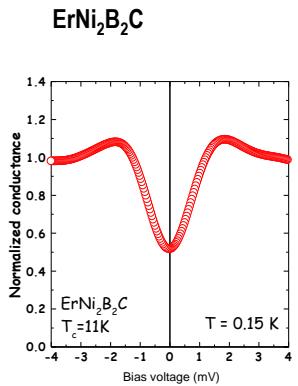
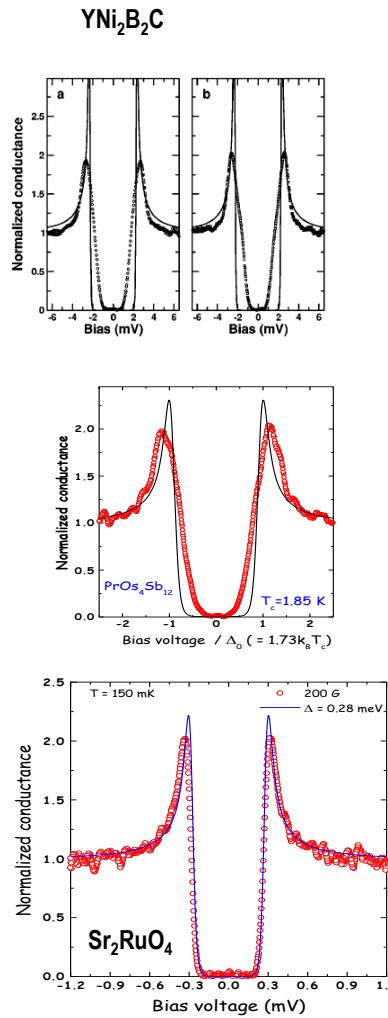
**Bardeen's formalism**  $I \propto \sum_{\nu} |\psi_{\nu}(\mathbf{d}_0)|^2 \delta(E_{\nu} - E_F) = N(\mathbf{d}_0, E_F)$

**The local density of states**  $\frac{dI}{dV} \propto \int N(E) \frac{\partial f(E - eV)}{\partial V} dE$

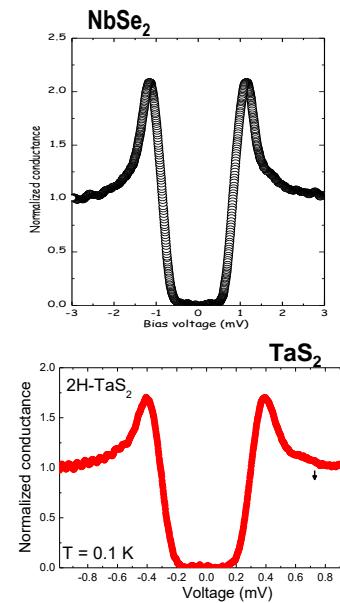
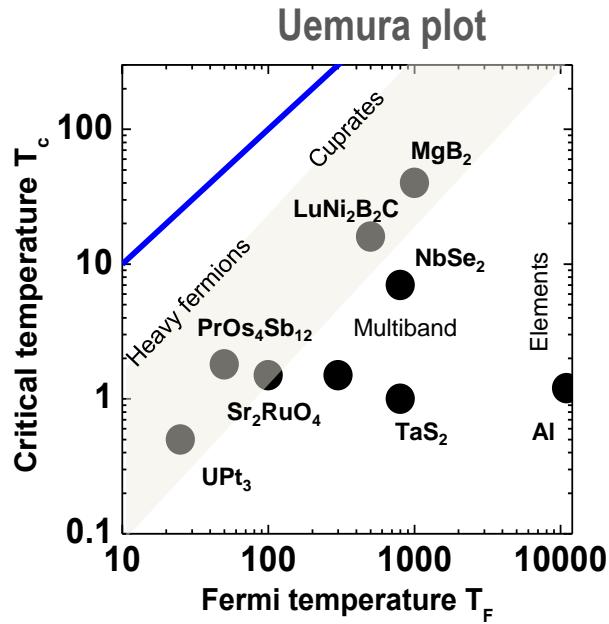
$$N_s(E) = \sum_i \gamma_i \frac{E}{\sqrt{E^2 - \Delta_i(\mathbf{k})^2}}$$

P. Martínez-Samper et al. Physica C 2003

## The superconducting gap through scanning tunneling spectroscopy at very low temperatures



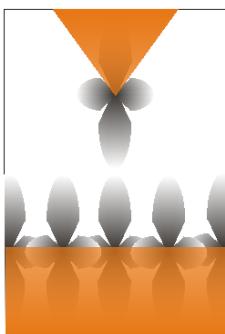
$$\frac{dI}{dV} \propto \int N(E) \frac{\partial f(E - eV)}{\partial V} dE$$



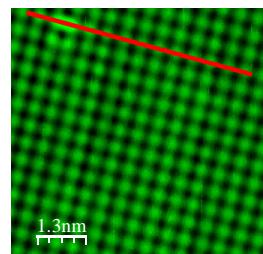
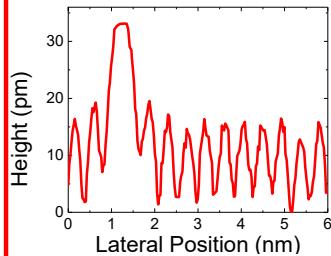
Madrid

## Can we find atomic resolution in a metal ?

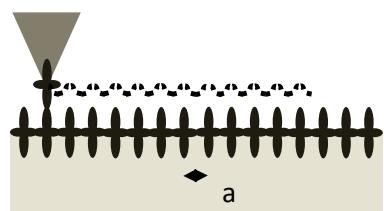
STM



Atomic resolution



Here: URu<sub>2</sub>Si<sub>2</sub> at 150 mK

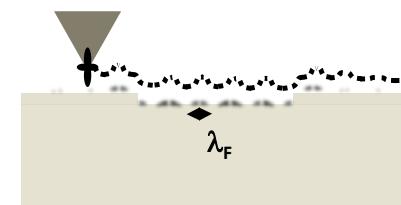
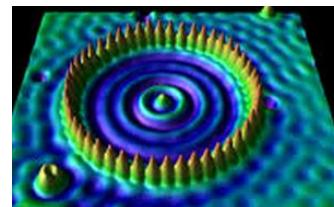


Quasiparticle interference

Electronic standing waves

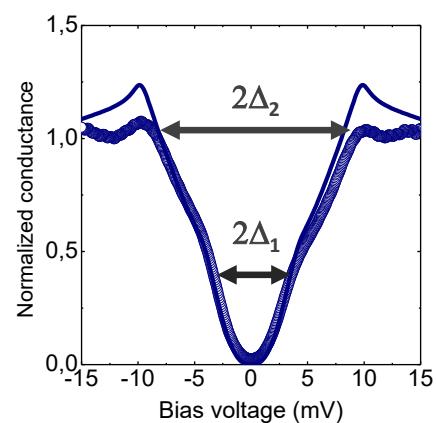
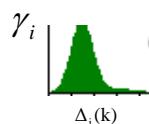
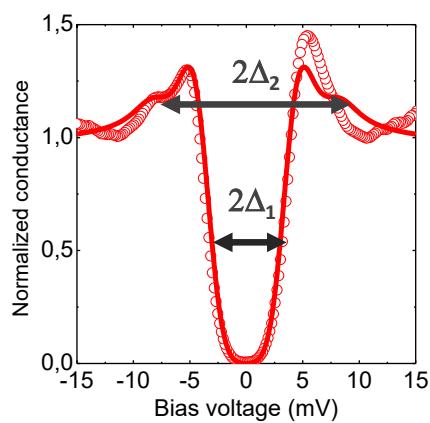
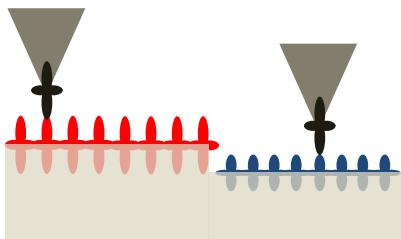


$$\rho(r) = \frac{\cos(2k_Fr)}{r^3}$$



$$I \propto \sum_{\nu} |\psi_{\nu}(d_0)|^2 \delta(E_{\nu} - E_F) = N(d_0, E_F)$$

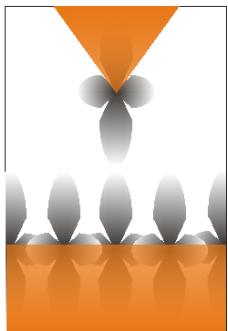
## Atomic resolution



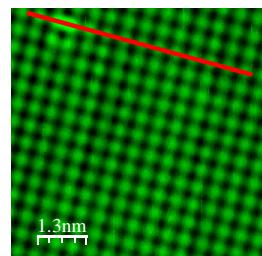
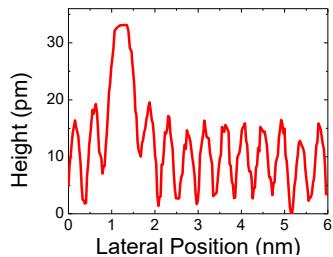
$$N_s(E) = \sum_i \gamma_i \frac{E}{\sqrt{E^2 - \Delta_i(k)^2}}$$

## Can we find atomic resolution in a metal ?

STM

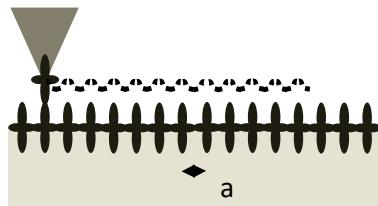


Atomic resolution



Here:  $\text{URu}_2\text{Si}_2$  at 150 mK

$$I \propto \sum_v |\psi_v(d_0)|^2 \delta(E_v - E_F) = N(d_0, E_F)$$

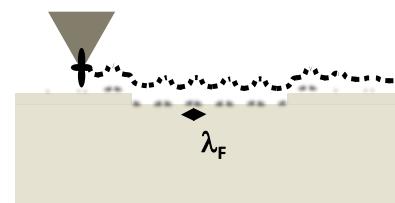
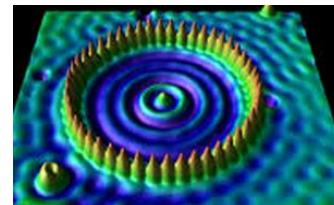


Quasiparticle interference

Electronic standing waves



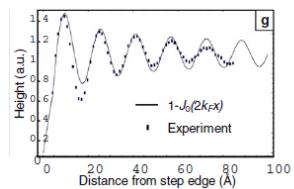
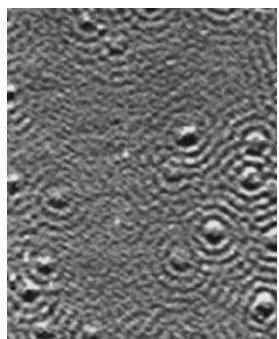
$$\rho(r) = \frac{\cos(2k_F r)}{r^3}$$



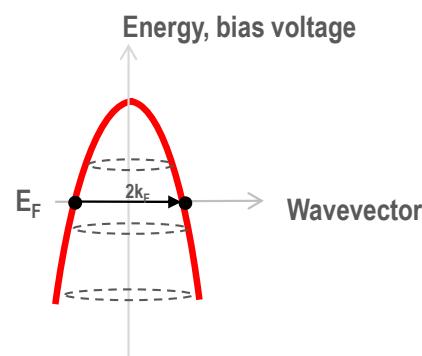
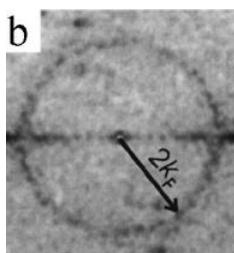
# QPI

## Quasiparticle interference Bandstructure of empty and filled states

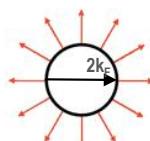
Real space conductance map



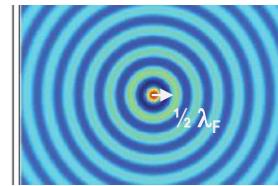
Fourier transform



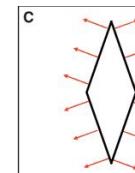
Fermi surface



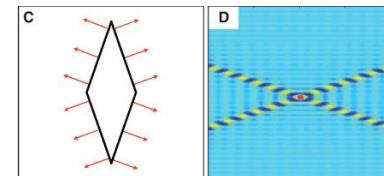
Real space DOS



Fermi surface



Real space DOS



Direct imaging of the two-dimensional Fermi contour: Fourier-transform STM  
 L. Petersen et al.

Phys. Rev. B, R6858, 57 (1998)

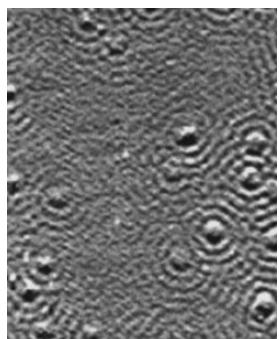
A phenomenological approach of joint density of states for the determination of  
 bandstructure in the case of a semi-metal studied by FT-STS

L. Simon, F. Vonau and D. Aubel  
 J. Phys. Cond. Matt. 19 (2007) 355009

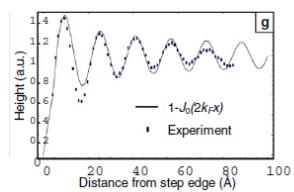
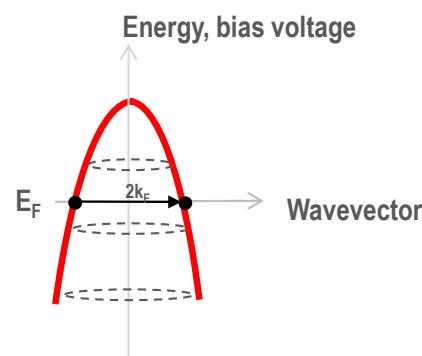
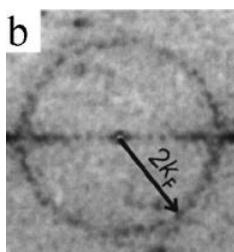
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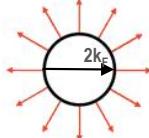
Real space conductance map



Fourier transform



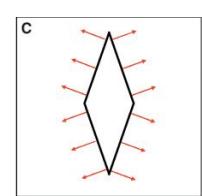
Fermi surface



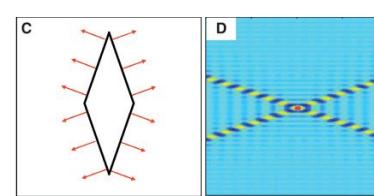
Real space DOS



Fermi surface



Real space DOS



$$g(E, \vec{q}) \propto |V_s(\vec{q})|^2 J(E, \vec{q})$$

$$J(E, \vec{q}) \propto N_1(E, \vec{k}_1) N_2(E, \vec{k}_2)$$

Direct imaging of the two-dimensional Fermi contour: Fourier-transform STM

L. Petersen et al.

Phys. Rev. B, R6858, 57 (1998)

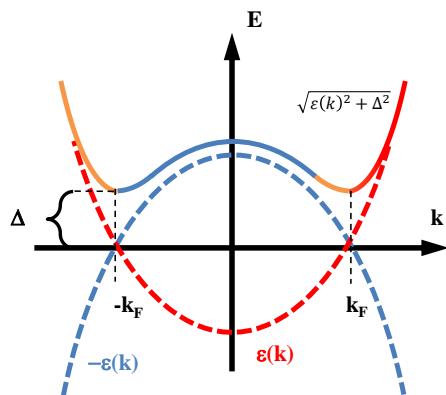
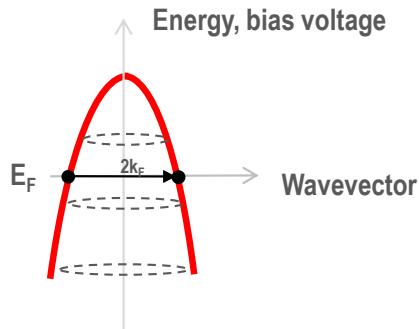
A phenomenological approach of joint density of states for the determination of bandstructure in the case of a semi-metal studied by FT-STS

L. Simon, F. Vonau and D. Aebel

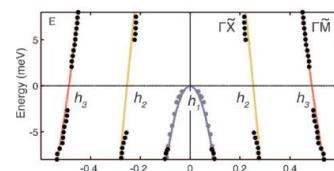
J. Phys. Cond. Matt. 19 (2007) 355009

# QPI

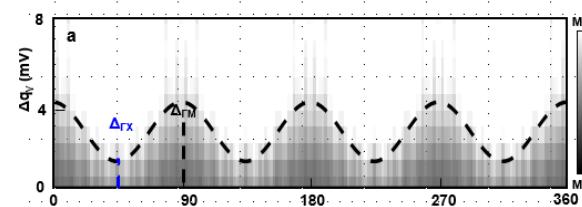
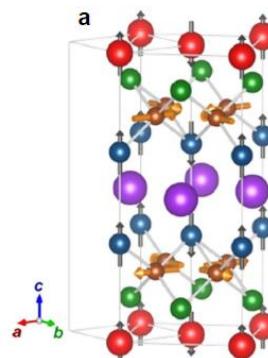
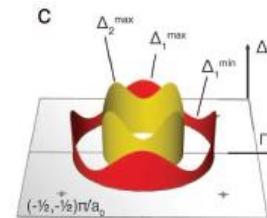
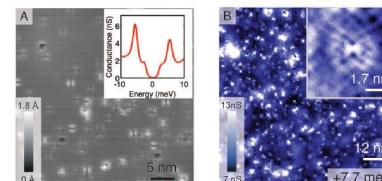
## Quasiparticle interference Superconducting gap anisotropy



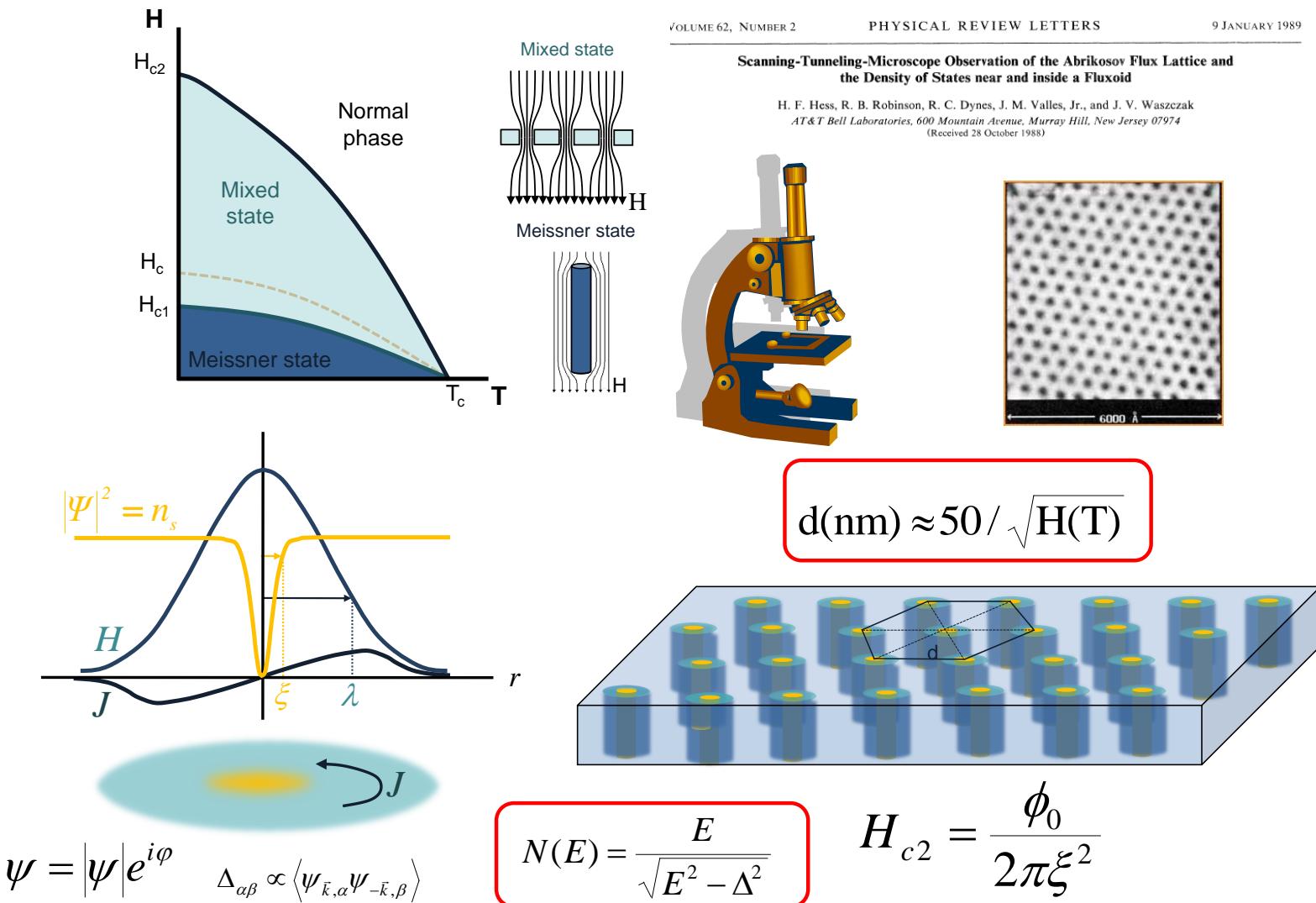
Spectroscopic scanning tunneling microscopy  
insights into Fe-based superconductors  
J. Hoffman Rep. Prog. Phys. 74 (2011) 124103



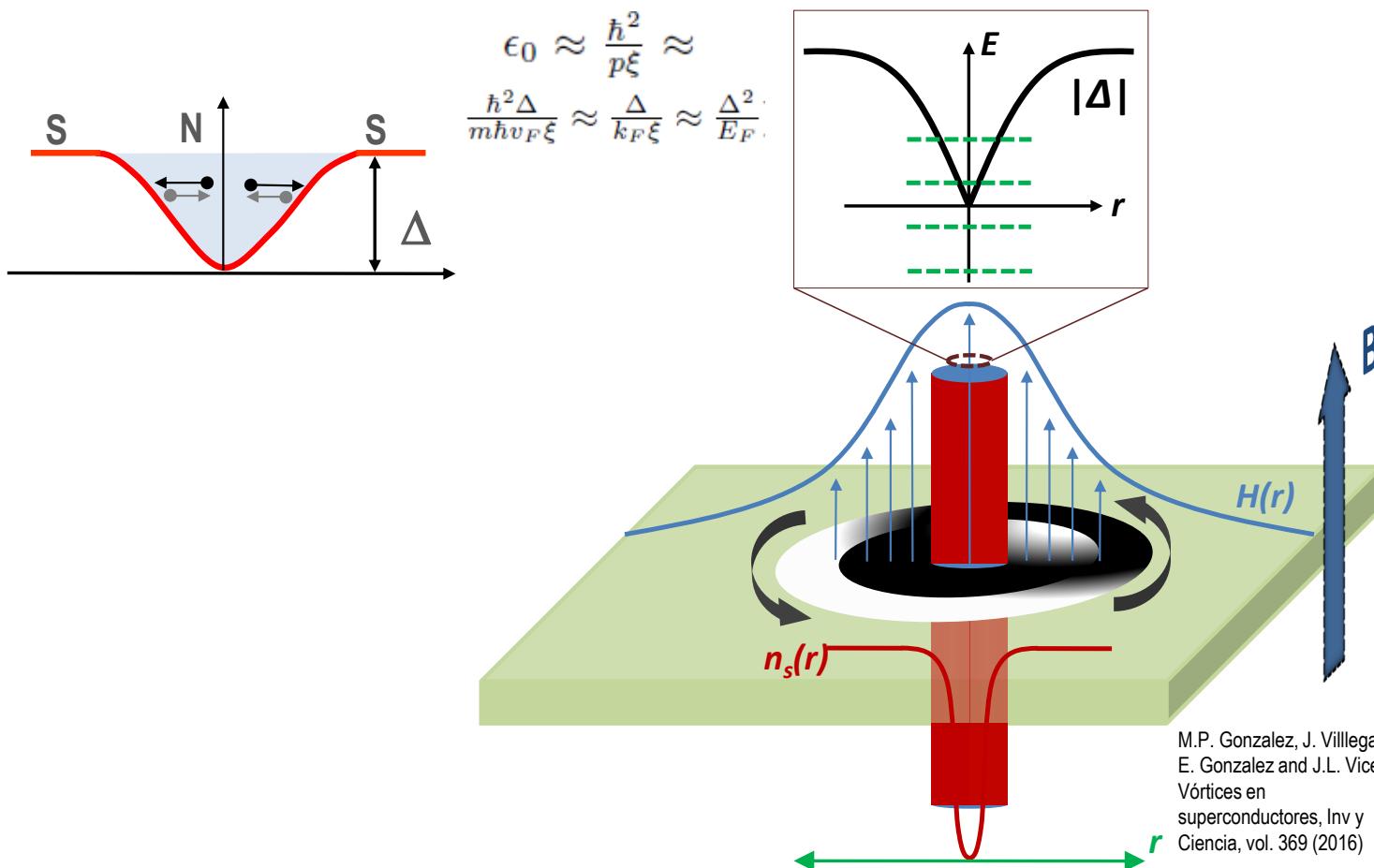
Anisotropic energy gaps of iron-based superconductivity from intraband quasiparticle scattering in LiFeAs  
Science 336 (2012) 563



## Type II superconductors : high upper critical field + the vortex lattice



## Bound core states: Caroli-de Gennes-Matricon states

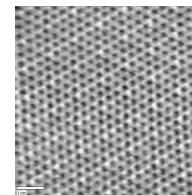


M.P. Gonzalez, J. Villegas,  
 E. Gonzalez and J.L. Vicent,  
 Vórtices en  
 superconductores, Inv y  
 Ciencia, vol. 369 (2016)

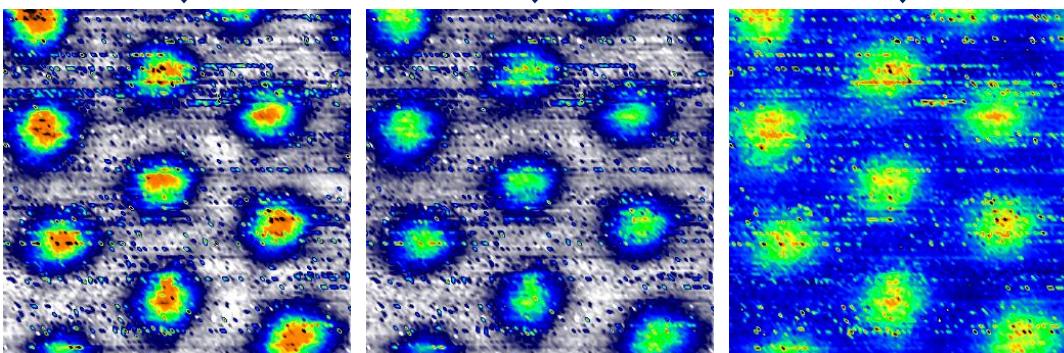
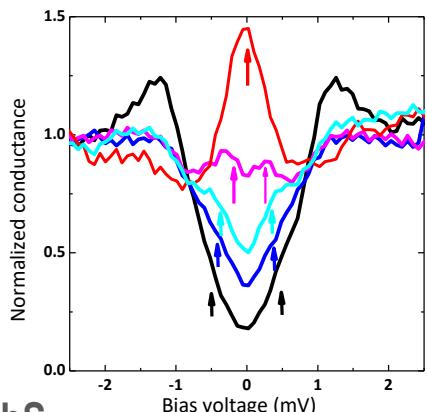
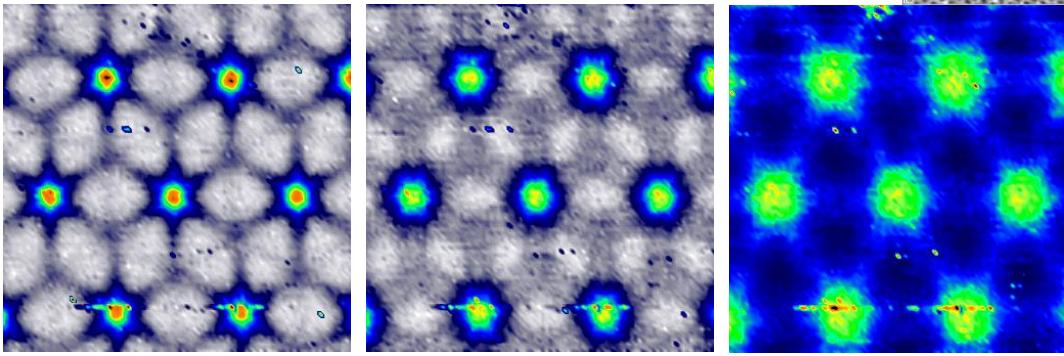
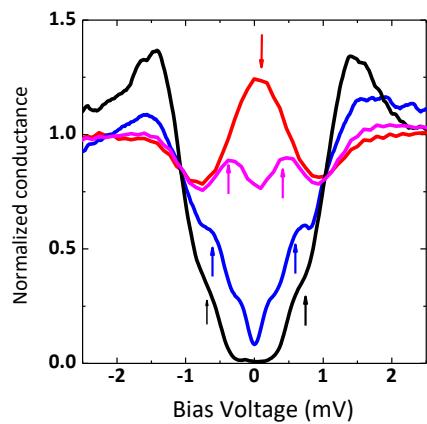
The vortex physics which comes from the vortex core  
 N.B. Kopnin  
*Physica. B*, 280, 231, (2000)

Current carried by bound states of superconducting vortices  
 D. Rainer, J.A. Sauls and D. Waxmann  
*Phys. Rev. B*, 54, (1996)

## Internal electronic structure of the vortex cores: $\text{NbSe}_2$ vs. $\text{NbS}_2$

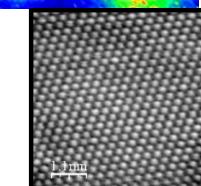


$\text{NbSe}_2$



$\text{NbS}_2$

*Superconducting density of states and vortex cores in NbS<sub>2</sub>*  
I. Guillamon, H. Suderow, S. Vieira and P. Rodiere  
Phys Rev Lett, 101, 166407 (2008)

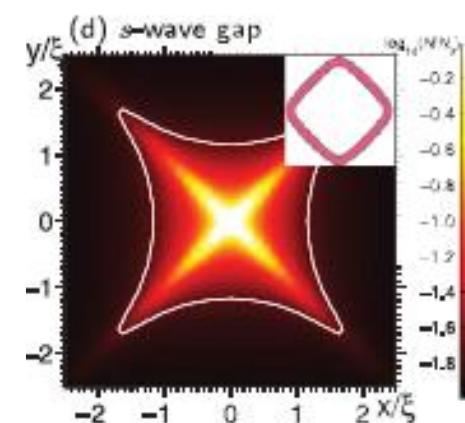
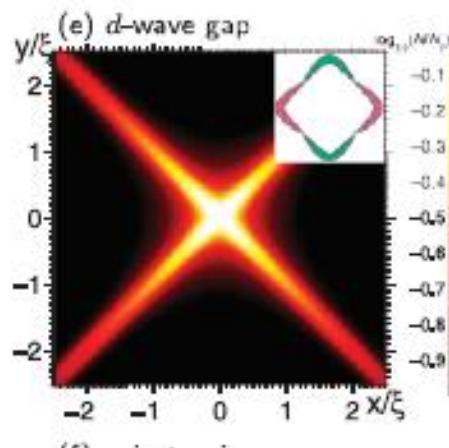
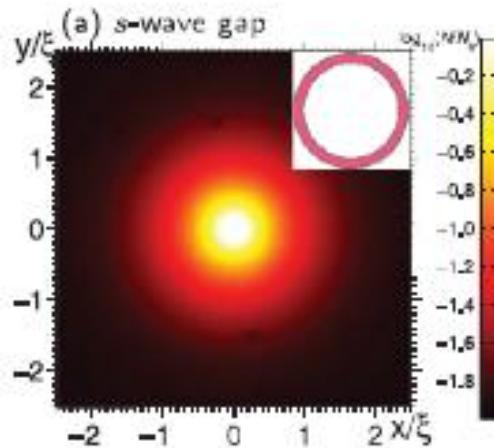


## Shape of the vortex core depends on bandstructure and gap symmetry

$$\xi = \frac{\hbar v_F}{\pi \Delta}$$

$$\Delta = \Delta(\theta)$$

$$v_F = \langle v_F \rangle_{FS} = \int \int dz d\theta N(z, \theta) v_F$$



Angular band structure of a vortex line in a type II superconductor

F. Gygi and M. Schlüter

Phys. Rev. Lett., 65, 1820 (1990)

Star shaped local density of states around vortices in a type II superconductor

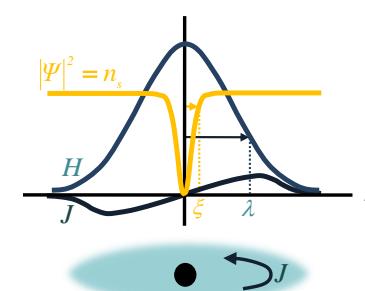
M. Ichioka, N. Hayashi and K. Machida

Phys. Rev. Lett. 77, 4074 (1996)

Theory of quasiparticle bound states in iron-based superconductors: application to scanning tunneling spectroscopy of LiFeAs

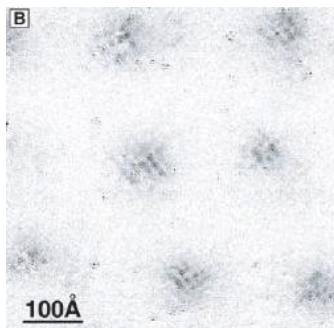
Y. Wang, P.J. Hirschfeld and I. Vekhter

Phys Rev B, 85, 020506(R) (2012)

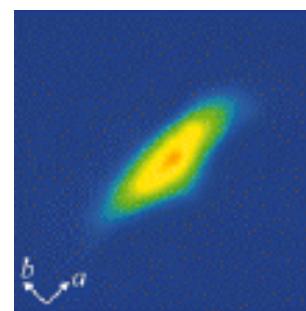


## Vortex cores in perpendicular fields by STM

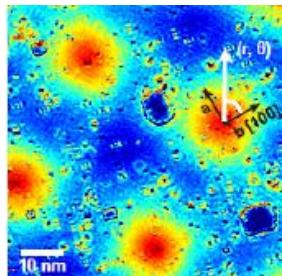
High  $T_c$  materials:  
Geneva , Paris, Tokyo/Nagoya, Cornell , Princeton, ...



FeAs, FeSe  
Beijing, Harvard, Tokyo ...

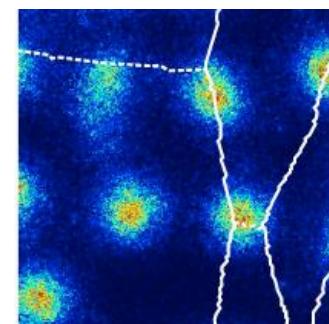


Heavy fermions  
Princeton, Cornell

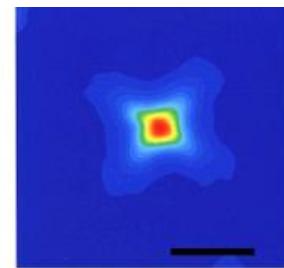


Imaging superconducting vortex cores and lattices with a STM  
H. Suderow, I. Guillamón, J.G. Rodrigo and S. Vieira  
Supercond. Sci. Techn., 27, 063001 (2014)

Very thin films  
Paris, ...

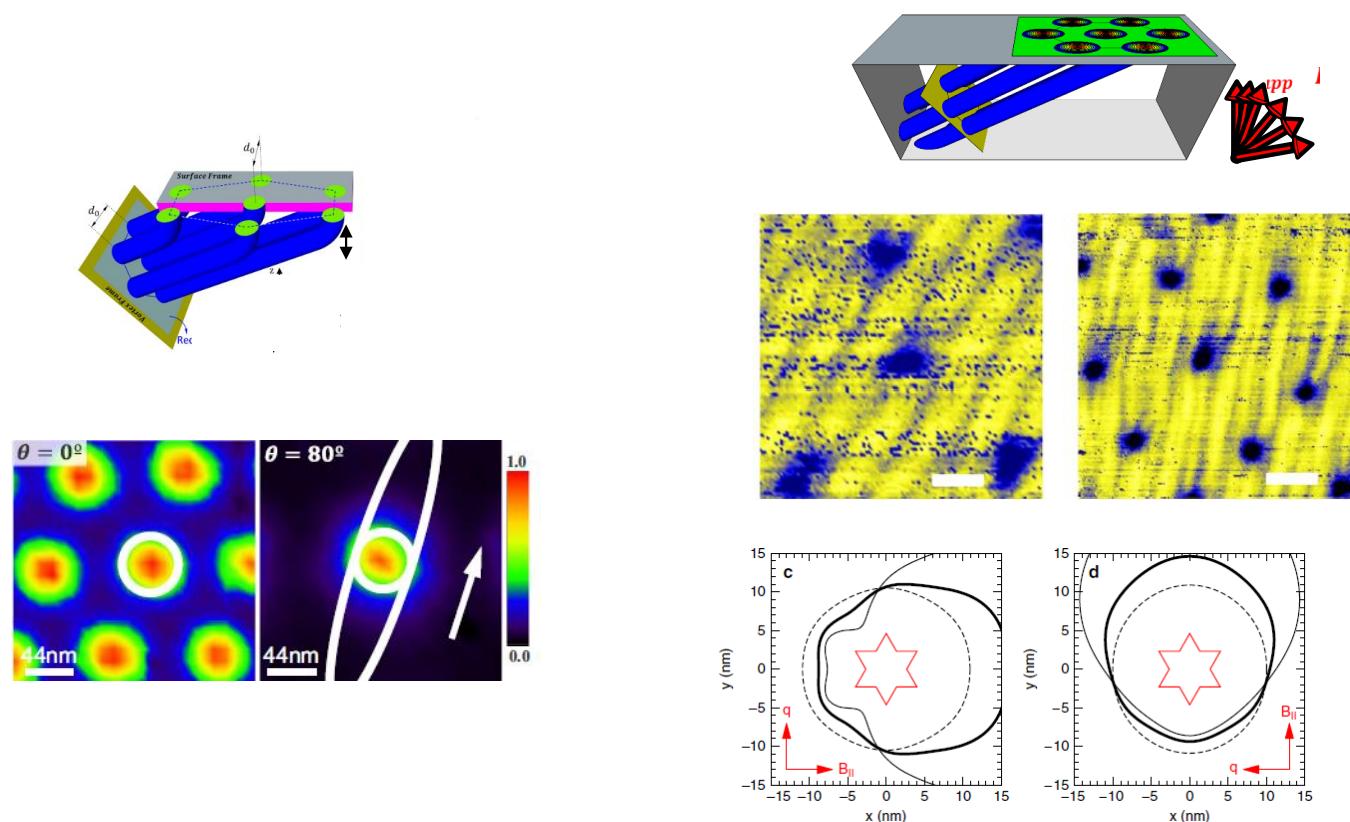


Nickel borocarbides  
Tokyo/Nagoya, Argonne, Madrid



Scanning tunneling spectroscopy of high-temperature superconductors  
O. Fischer, M. Kugler, I. Maggio-Aprile, Ch. Berthod  
Rev Mod Phys, 79, 373 (2007)

## Vortex cores in tilted fields by STM

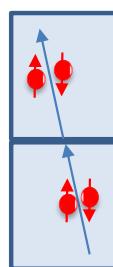
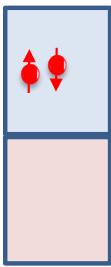
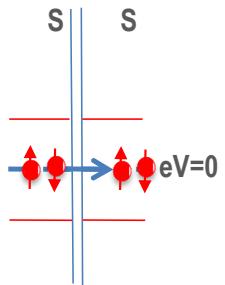


Subsurface bending and reorientation of tilted vortex lattices in bulk isotropic superconductors due to Coulomb-like repulsion at the surface  
**E. Herrera et al, Phys Rev B, 96, 184502 (2017)**

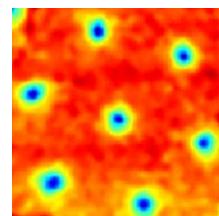
Tilted vortex cores and superconducting gap anisotropy in 2H-NbSe<sub>2</sub>  
**J.A. Galvis et al, Comm. Phys. 2:31 (2019)**

Lateral imaging of the superconducting vortex lattice using Doppler-modulated scanning tunneling microscopy, **I. Fridman et al, Appl. Phys. Lett. 99, 192505 (2011)**

## Superconducting tip: Josephson



Josephson



$$\psi = |\psi| e^{i\varphi}$$

$$\varphi_L \quad | \quad \varphi_R$$

$$\delta\varphi = \varphi_L - \varphi_R$$

$$I = I_0 \sin(\delta\varphi)$$

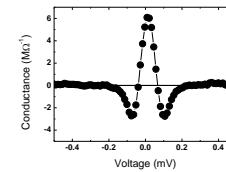
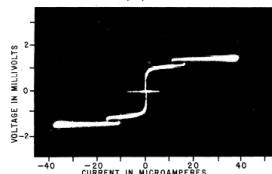
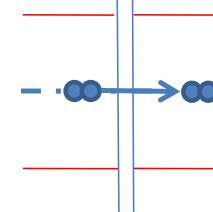
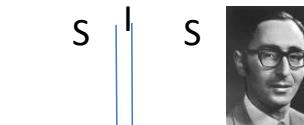
$$\frac{\partial \delta\varphi}{\partial t} = \frac{2e}{\hbar} V$$

Phase ?

$$I = I_0 \boxed{\sin(\delta\varphi)}$$

Density of Cooper  
pairs  $n_{s1}$  and  $n_{s2}$

### Cooper pair tunneling



## Quantum phase as a measurable quantity: DC and AC Josephson effect



$$\Psi_1 = \psi_{0,1} e^{i\phi_1}$$

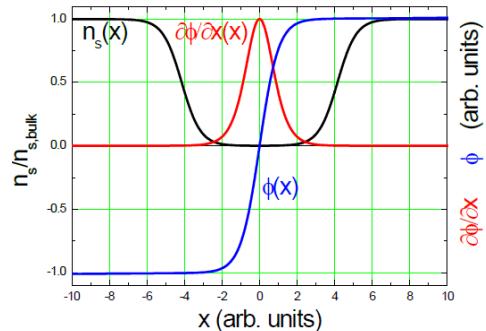
$$\Psi_2 = \psi_{0,2} e^{i\phi_2}$$

$$n_{S1} = |\psi_{0,1}|^2$$

$$n_{S2} = |\psi_{0,2}|^2$$

$$j_s = 2e n_s v_s = \frac{\hbar e}{m} n_s \nabla \phi$$

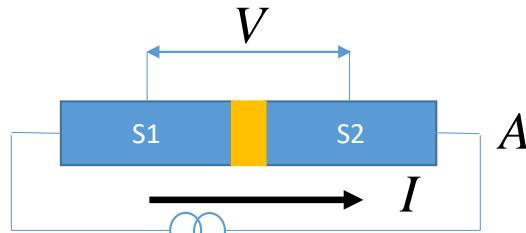
$$j_s \propto n_s \nabla \phi$$



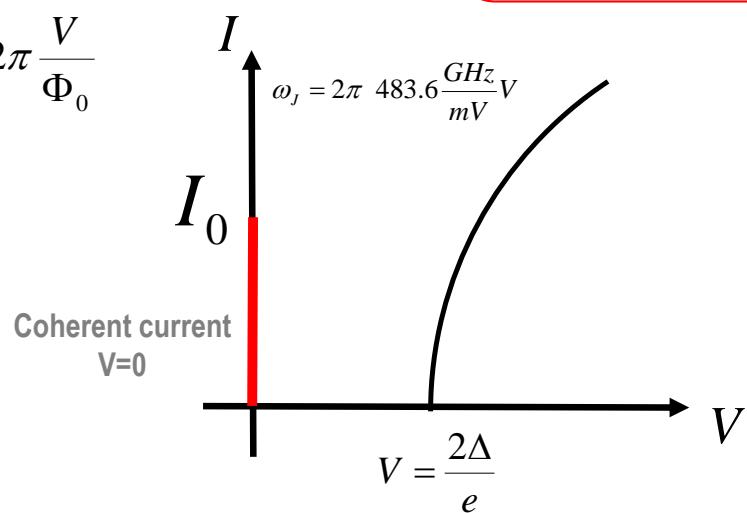
$$j_s = \frac{I}{A} = cte$$

$$n_s \frac{\partial \phi(x)}{\partial x} = cte$$

## Quantum phase as a measurable quantity: DC and AC Josephson effect



$$\omega_J = 2\pi \frac{V}{\Phi_0}$$



$$I + I_n(t) = I_0 \sin \delta + \frac{\Phi_0}{2\pi R} \dot{\delta} + \frac{\Phi_0 C}{2\pi} \ddot{\delta}$$

$$I = I_0 \sin(\delta)$$

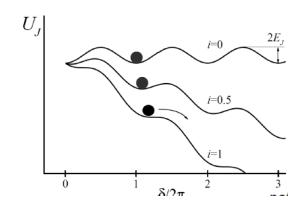
First Josephson equation

$$\frac{\partial \delta}{\partial t} = \frac{2e}{\hbar} V$$

Second Josephson equation

$$\delta(t) = \delta_0 + 2\pi \frac{V}{\Phi_0} t$$

$$I = I_0 \sin(\delta_0 + 2\pi \frac{V}{\Phi_0} t)$$



## Josephson effect at atomic level

Eur. Phys. J. B 40, 483–488 (2004)  
DOI: 10.1140/epjb/e2004-00273-y

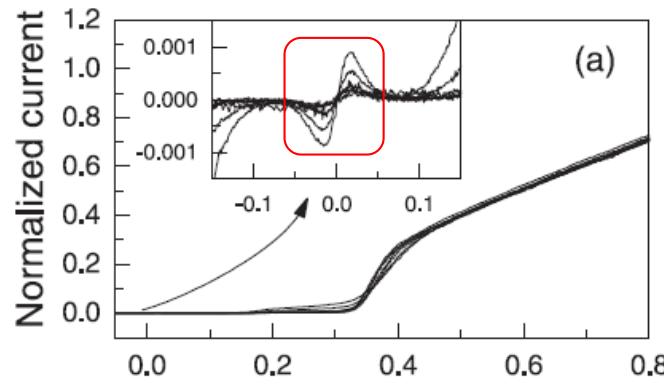
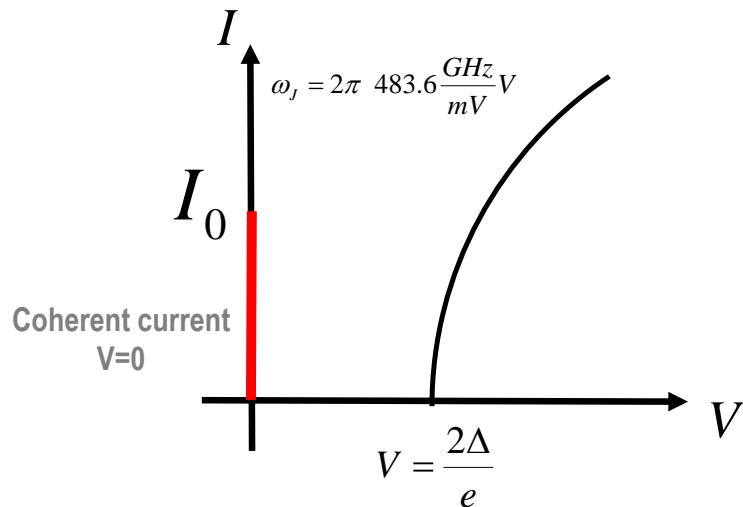
THE EUROPEAN  
PHYSICAL JOURNAL B

On the use of STM superconducting tips at very low temperatures

$E_J \rightarrow$  Junction resistance

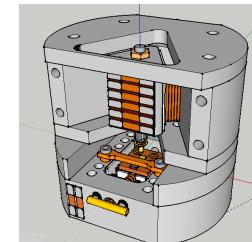
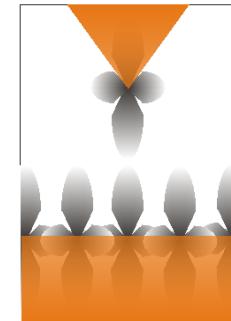
<<

$k_B T$  about 0.1 K



## Real space visualization of electronic correlations

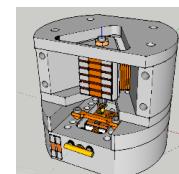
- Real space imaging and spectroscopy
  - Atomic structure
  - Electronic bandstructure
  - Superconducting gap structure
  - Vortex lattice structure
  - Josephson effect (phase)
  - Pairing interaction
  - Electronic correlations (Coulomb, Kondo)



## Real space imaging of the superconducting vortex lattice: Recent results and prospects

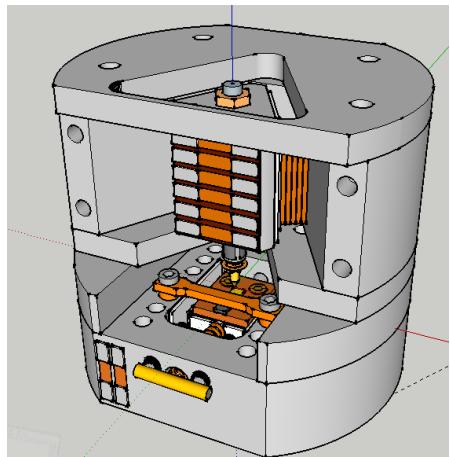


- Real space imaging and spectroscopy of the vortex lattice
- Superconducting CaKFe<sub>4</sub>As<sub>4</sub>, bandstructure, vortex core, vortex lattice and pinning
- Vortex lattice structure: interactions vs disorder
- Observation of vortex creep on cooling
- Josephson effect at atomic scale

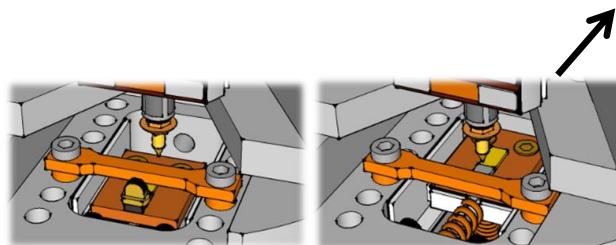


## Small compact dil fridge STM/S

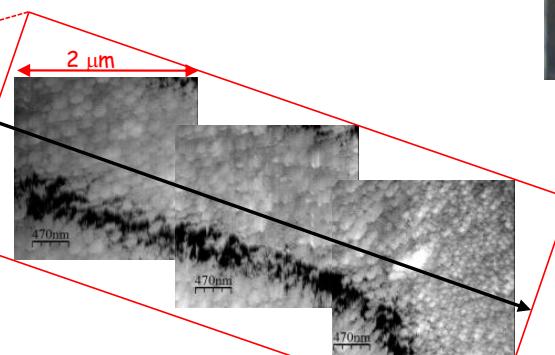
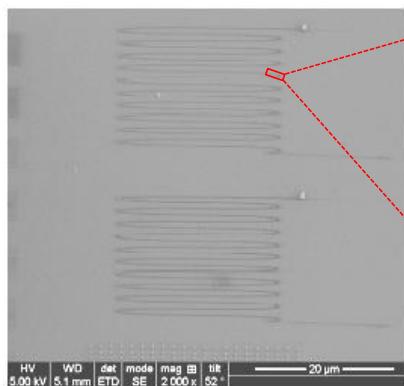
### Tip and sample preparation and scanning window



100 mK  
13  $\mu$ eV



100mK, 50 mK, 10 mK, 30 mK



REVIEW OF SCIENTIFIC INSTRUMENTS 82, 033711 (2011)

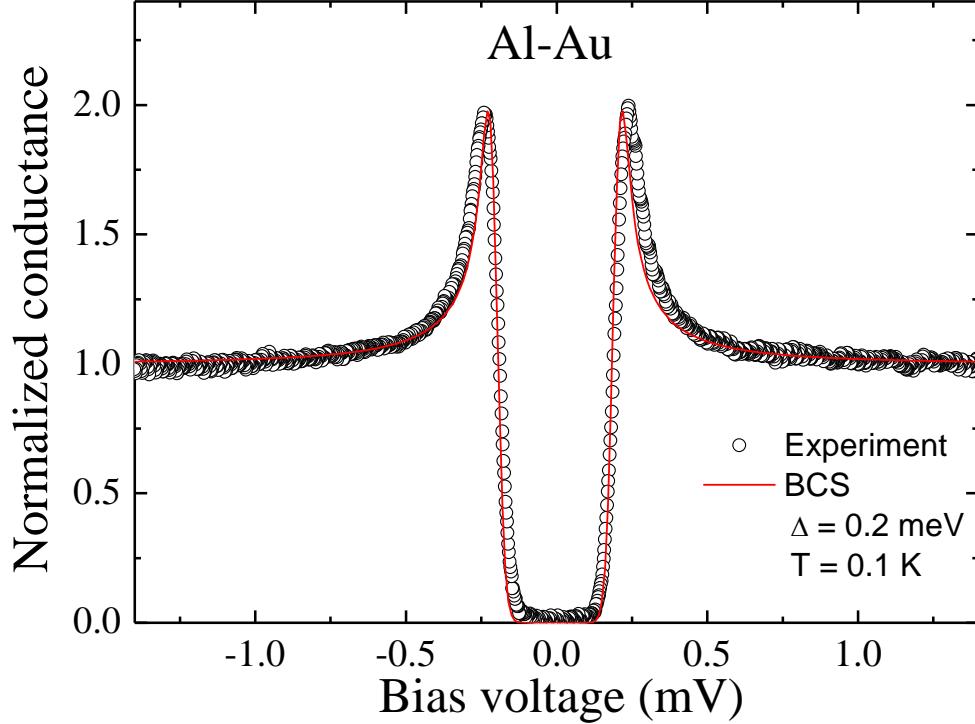
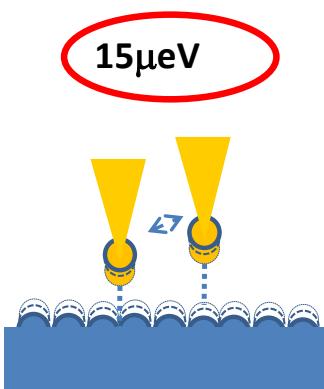
Compact very low temperature scanning tunneling microscope with mechanically driven horizontal linear positioning stage

H. Suderow,<sup>a)</sup> I. Guillamon, and S. Vieira  
Laboratorio de Bajas Temperaturas, Departamento de Física de la Materia Condensada Instituto de Ciencia  
de Materiales Nicolás Cabrera, Facultad de Ciencias Universidad Autónoma de Madrid, 28049 Madrid, Spain  
(Received 3 December 2010; accepted 19 February 2011; published online 23 March 2011)

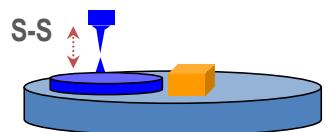
Patente

## Temperature

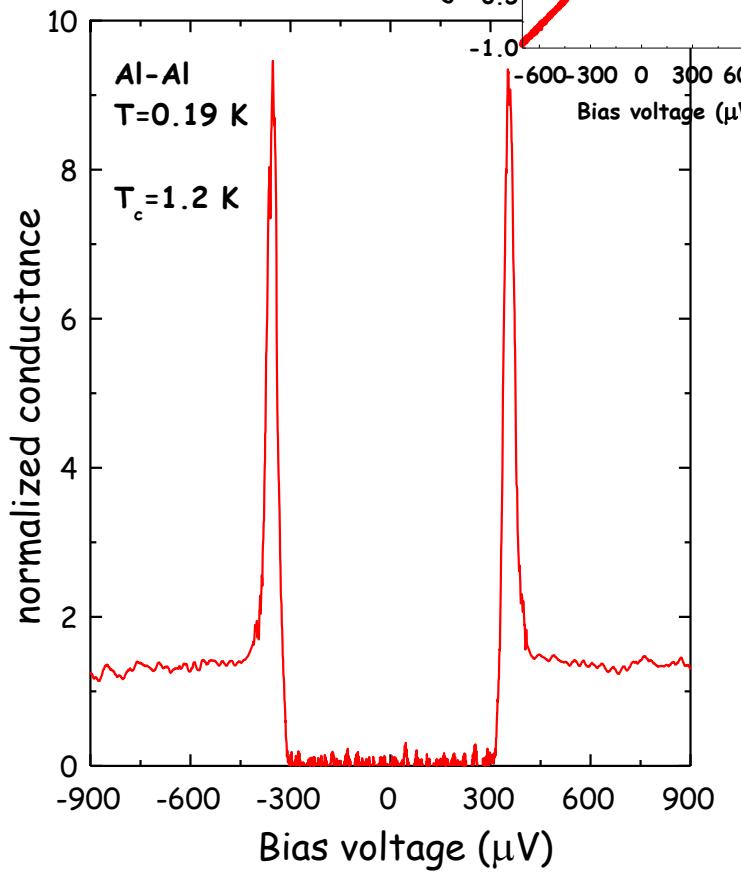
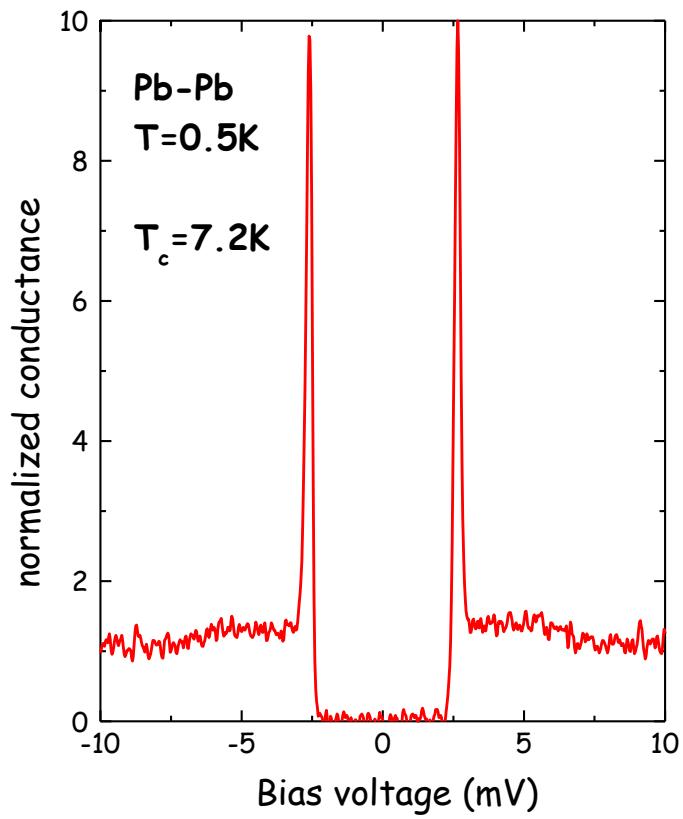
$$\frac{dI}{dV} \propto \int N_s(E) \frac{\partial f(E - eV)}{\partial V} dE$$



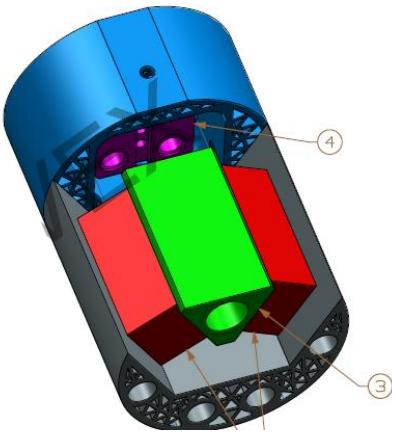
## Energy resolution



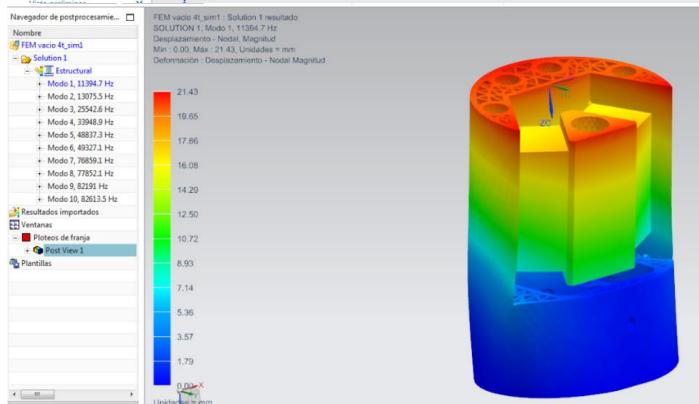
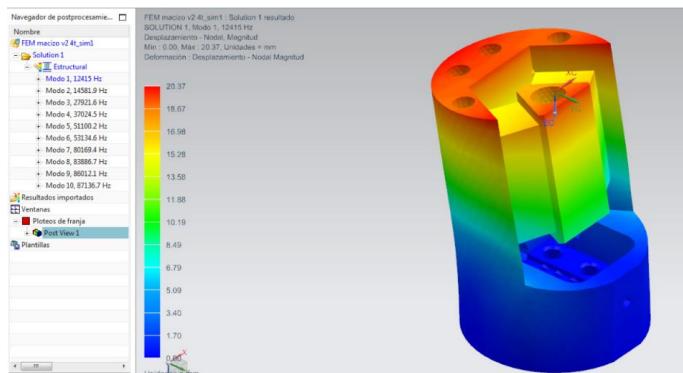
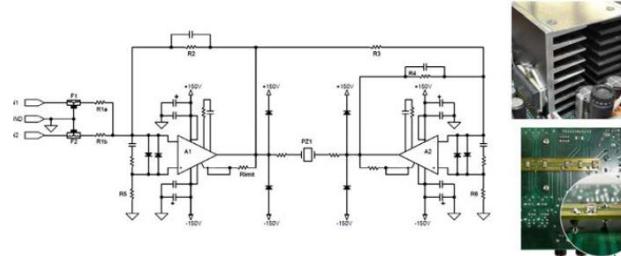
$15 \mu\text{V} \sim 183\text{mK}$



## New smaller STM/S



3D printed in grade 5 Titanium alloy



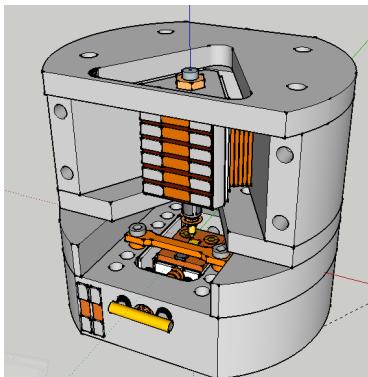
REVIEW OF SCIENTIFIC INSTRUMENTS 82, 033711 (2011)

Compact very low temperature scanning tunneling microscope with mechanically driven horizontal linear positioning stage

H. Suderow<sup>a)</sup>, I. Guillamon, and S. Vieira  
*Laboratorio de Bajas Temperaturas, Departamento de Física de la Materia Condensada Instituto de Ciencia de Materiales Nicolás Cabrera, Facultad de Ciencias Universidad Autónoma de Madrid, 28049 Madrid, Spain*  
 (Received 3 December 2010; accepted 19 February 2011; published online 23 March 2011)

## STM/S in Madrid

**22 T superconducting coil, hundreds of field ramps and no single quench  
in Madrid**

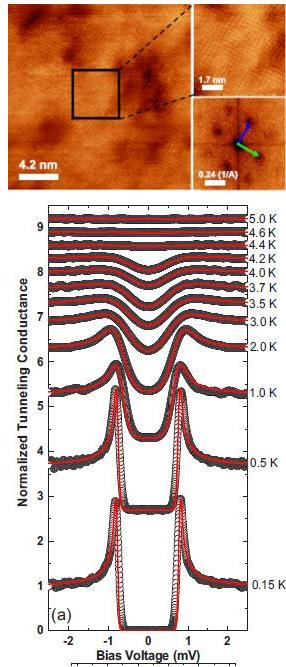


*Very high magnetic field STM (22 T), Isabel Guillamón*

[isabel.guillamon@uam.es](mailto:isabel.guillamon@uam.es)

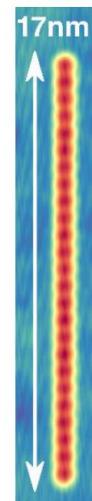
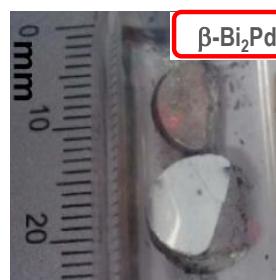


# QUAMaterials

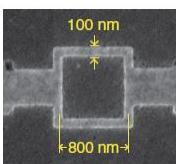


Superconductivity at 5.4 K in  $\beta\text{-Bi}_2\text{Pd}$   
 Y. Imai JPSJ 81, 113708

Full gap superconductivity in spin-polarised surface states of topological semimetal  $\beta\text{-Bi}_2\text{Pd}$   
 K. Iwaya Nat Comm, 8, 976 (2017)



J. Zaldivar, J.I. Pascual



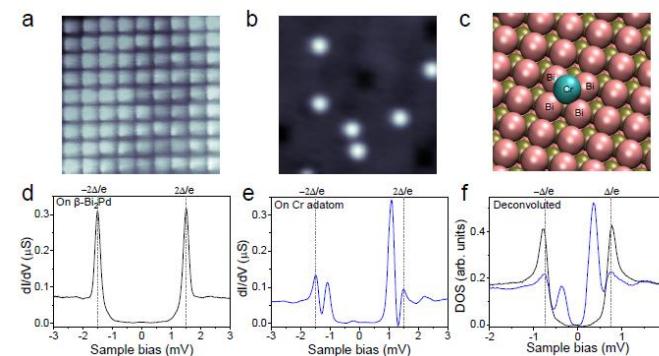
## RESEARCH

### SUPERCONDUCTIVITY

#### Observation of half-quantum flux in the unconventional superconductor $\beta\text{-Bi}_2\text{Pd}$

Yufan Li<sup>1\*</sup>, Xiaoying Xu<sup>1\*</sup>, M.-H. Lee<sup>2</sup>, M.-W. Chu<sup>2</sup>, C. L. Chien<sup>1,3,4†</sup>

## Shiba states on Cr adatoms

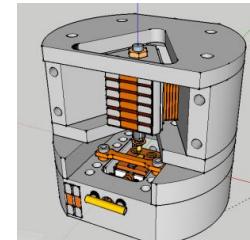


Influence of Magnetic Ordering between Cr Adatoms on the Yu-Shiba-Rusinov States of the  $\beta\text{-Bi}_2\text{Pd}$  Superconductor  
 DJ Choi et al PRL 120 (2018) 167001

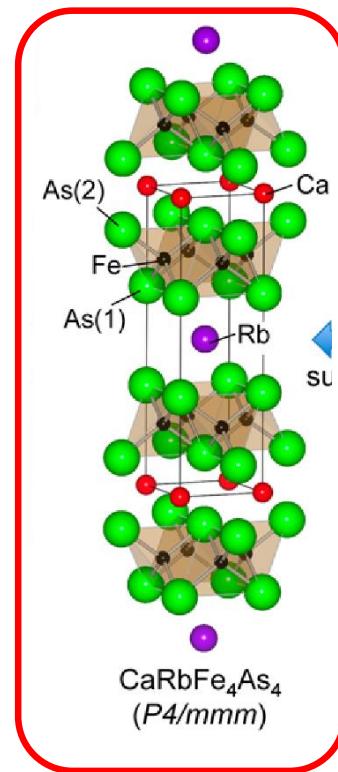
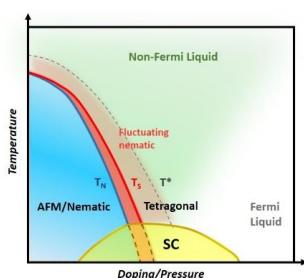
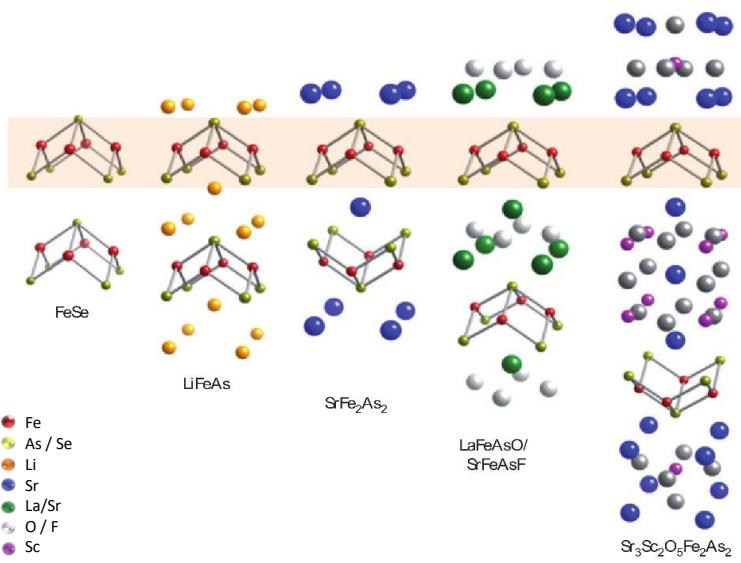
## Real space imaging of the superconducting vortex lattice: Recent results and prospects



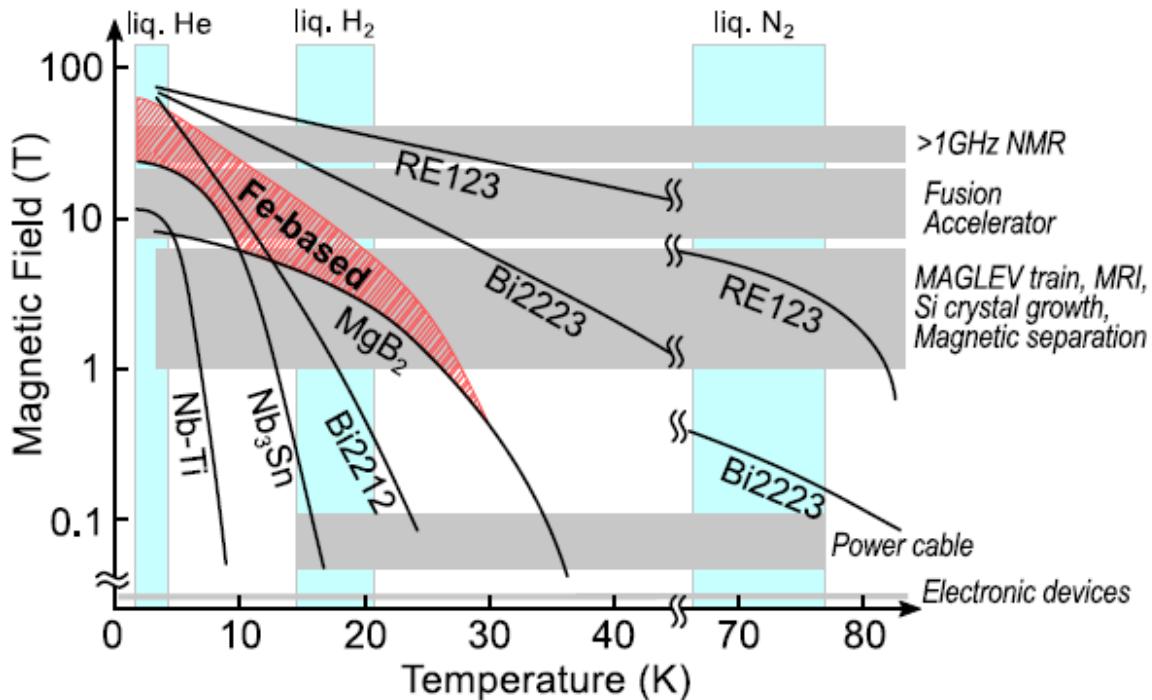
- Real space imaging and spectroscopy of the vortex lattice
- Superconducting CaKFe<sub>4</sub>As<sub>4</sub>, bandstructure, vortex core, vortex lattice and pinning
- Vortex lattice structure: interactions vs disorder
- Observation of vortex creep on cooling
- Josephson effect at atomic scale



## Crystal structures in the iron pnictides

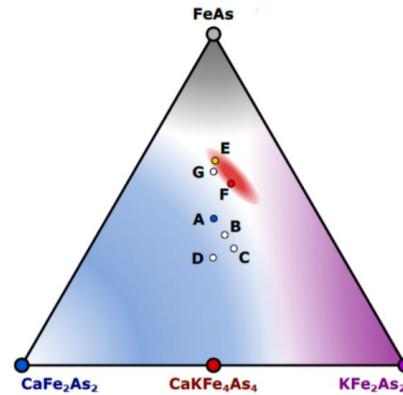
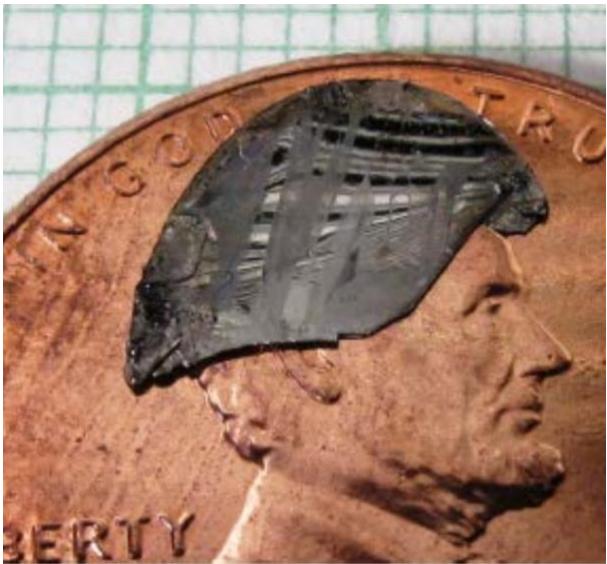


## Critical temperature and magnetic field



Potentials of iron-based superconducting alloys for practical future materials  
J.I. Shimoyama  
Superc. Sci. Tech., 044002, 27 (2014)

## The $\text{CaKFe}_4\text{As}_4$ system



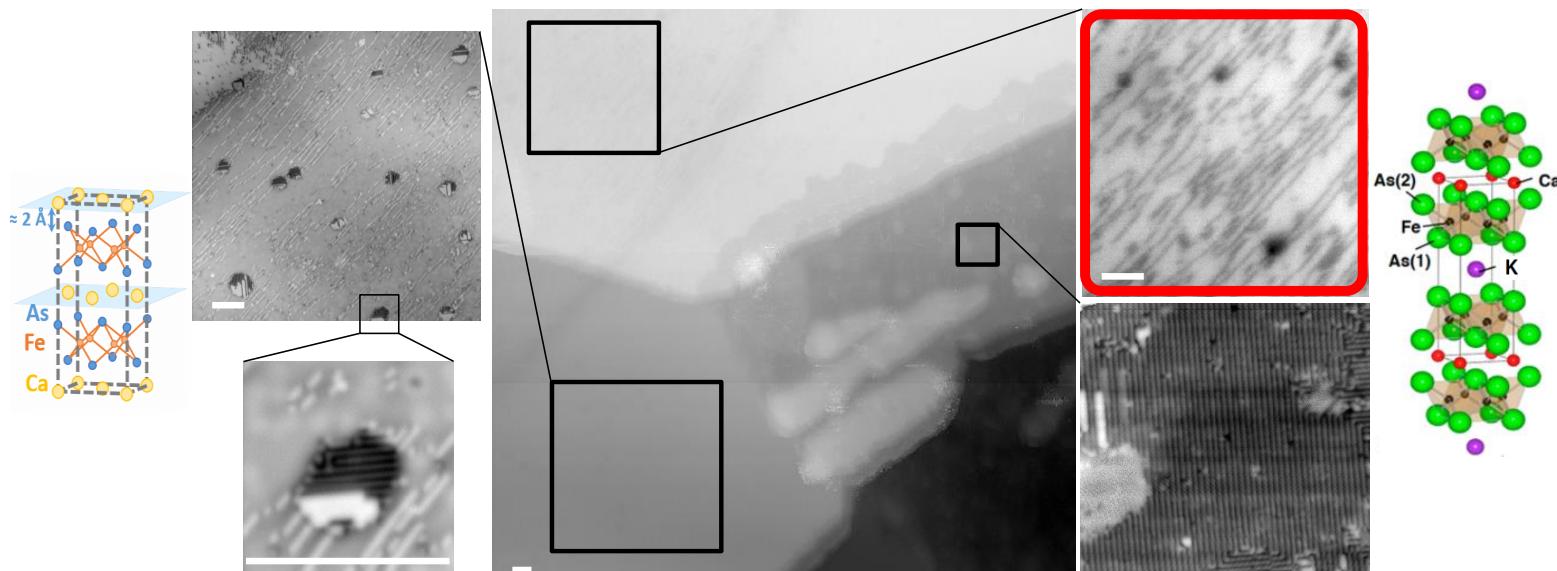
Wed. Afternoon

PC5-4, Unique defect structure and advantageous vortex pinning properties in  $\text{CaKFe}_4\text{As}_4$ , Ishida

PC5-5, Critical current density and its enhancement by particle irradiation in  $\text{CaK}_2\text{Fe}_4\text{As}_4\text{F}_2$ , Tamegai

## The CaKFe<sub>4</sub>As<sub>4</sub> system

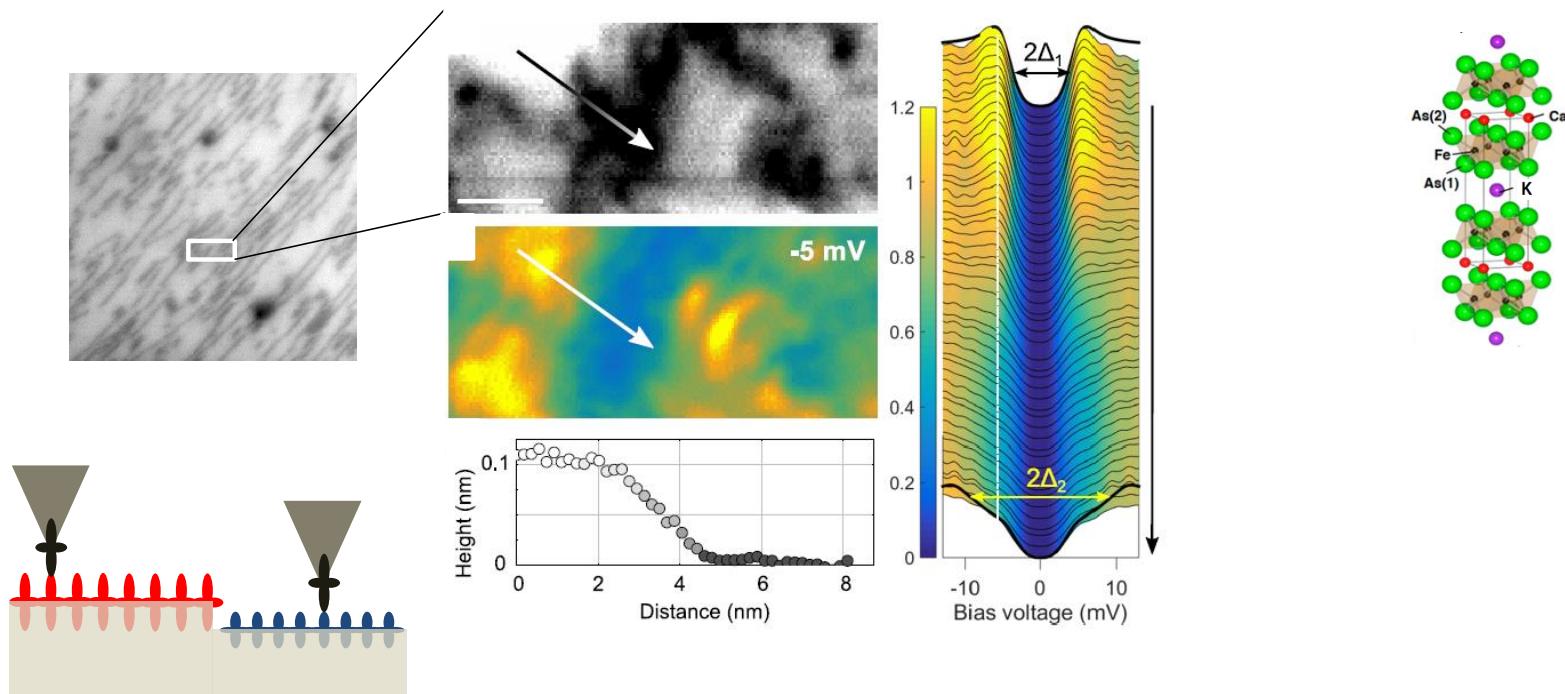
Topography of CaKFe<sub>4</sub>As<sub>4</sub>: No surface reconstruction  
Ca-K surface with stripes



White scale bar: 20 nm

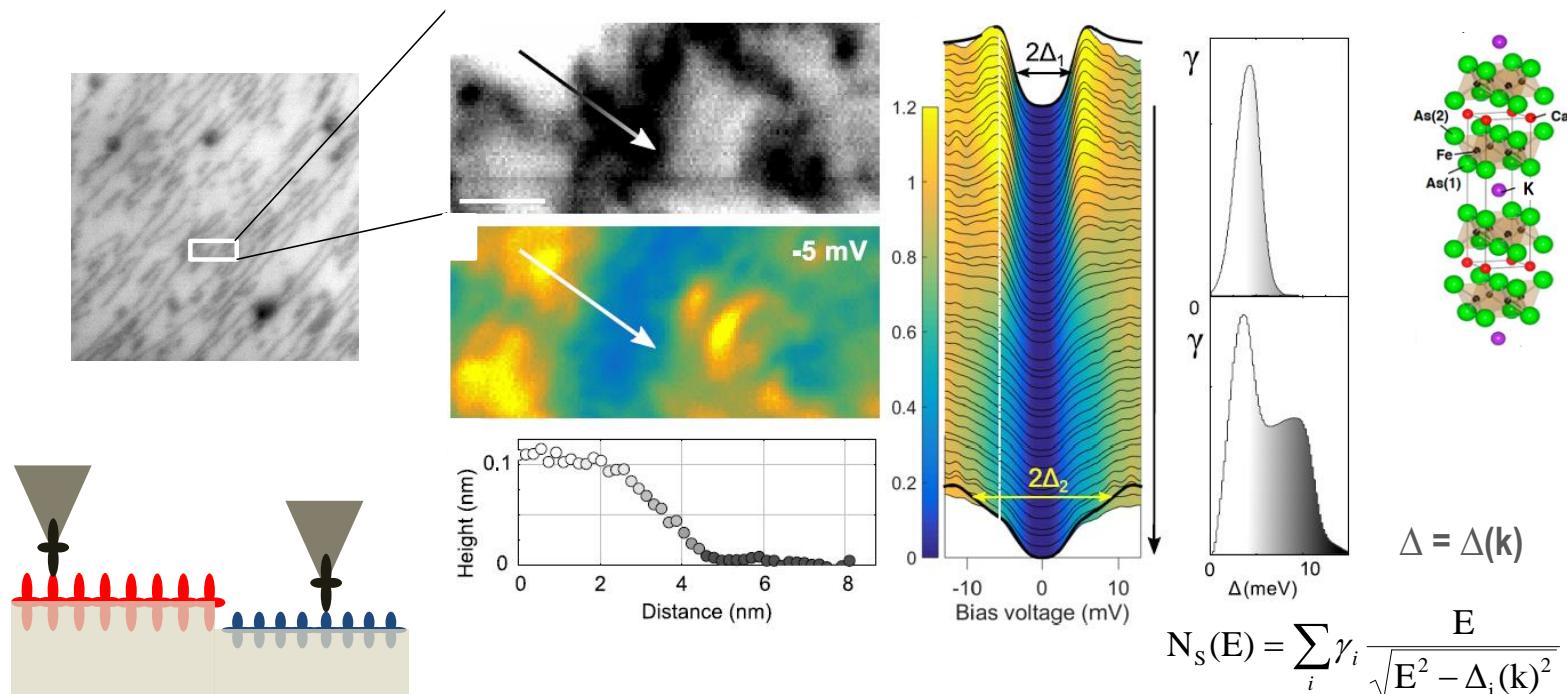
## The CaKFe<sub>4</sub>As<sub>4</sub> system

### Two gap superconductivity in CaKFe<sub>4</sub>As<sub>4</sub>, T<sub>c</sub> = 35 K Stripes



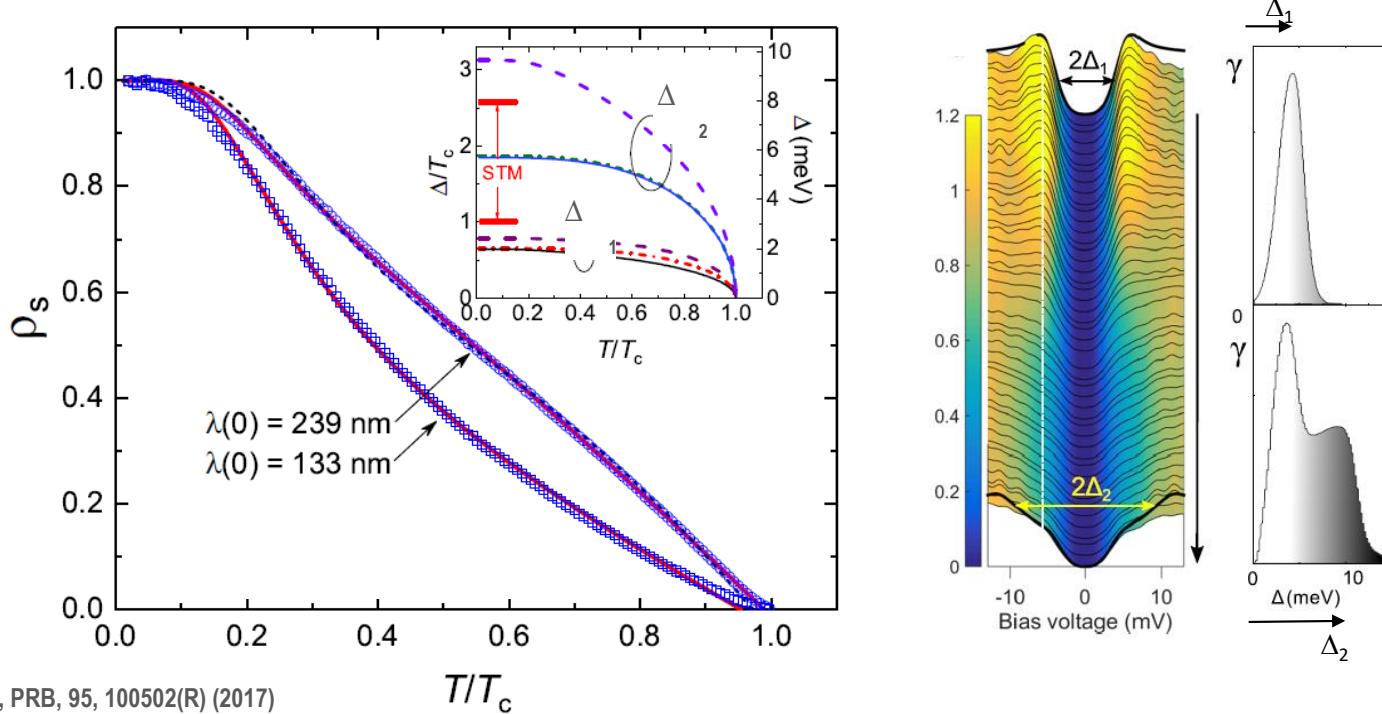
## The CaKFe<sub>4</sub>As<sub>4</sub> system

Two gap superconductivity in CaKFe<sub>4</sub>As<sub>4</sub>, T<sub>c</sub> = 35 K  
Stripes



## The CaKFe<sub>4</sub>As<sub>4</sub> system

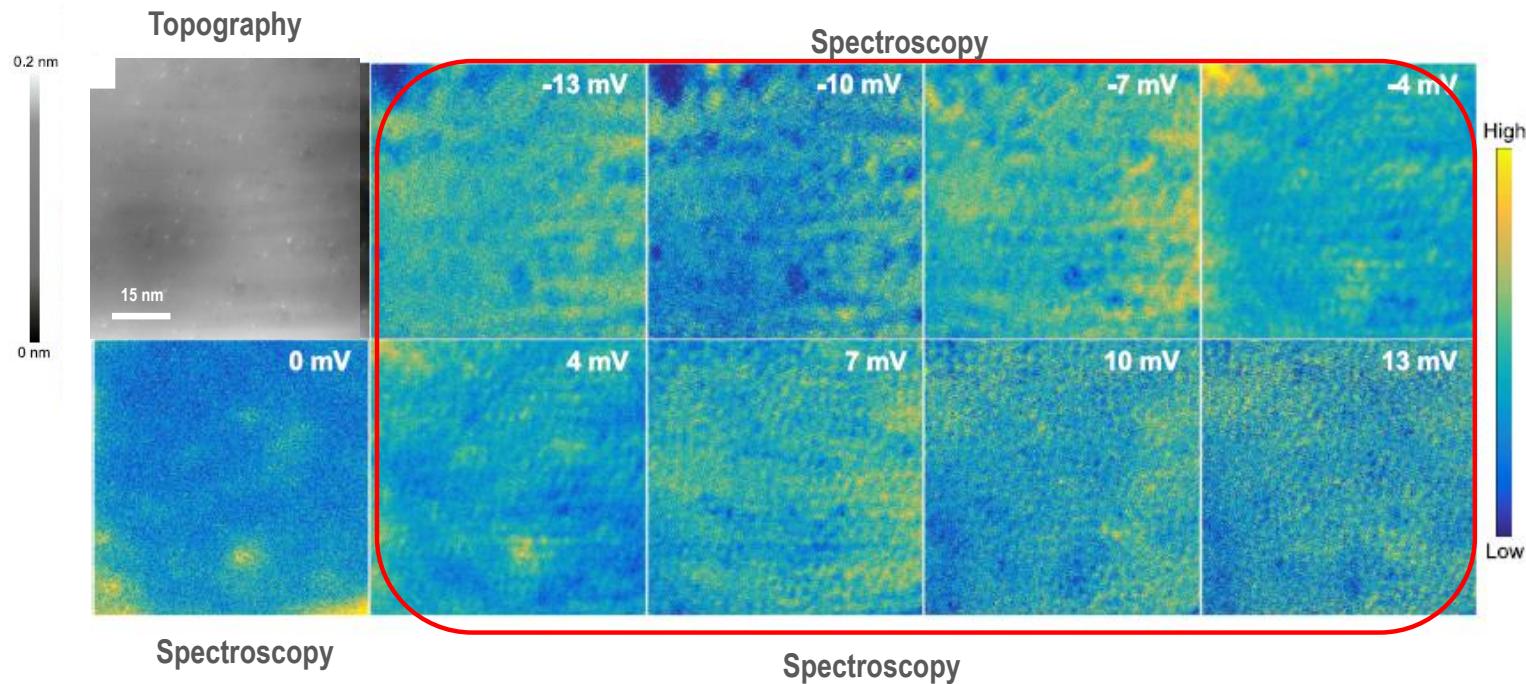
Two gap superconductivity in CaKFe<sub>4</sub>As<sub>4</sub>, T<sub>c</sub> = 35 K  
London penetration depth, R. Prozorov et al



K. Cho et al, PRB, 95, 100502(R) (2017)

## The CaKFe<sub>4</sub>As<sub>4</sub> system

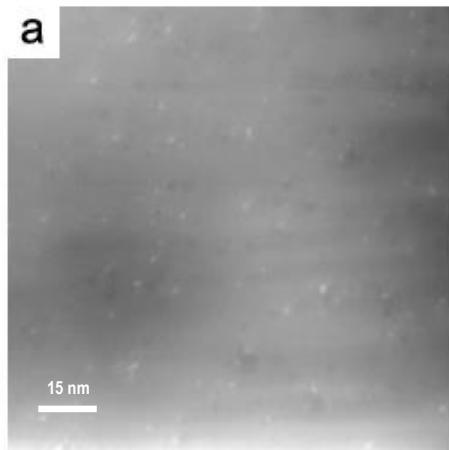
### Quasiparticle interference scattering



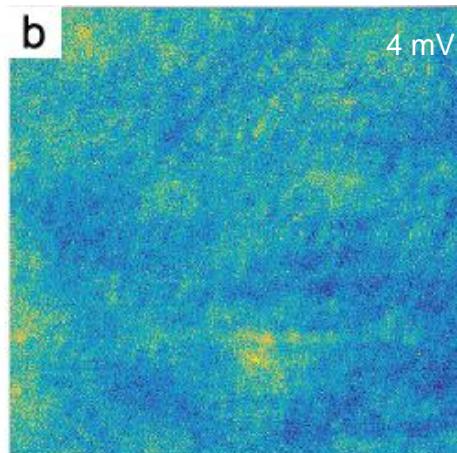
## The CaKFe<sub>4</sub>As<sub>4</sub> system

### Quasiparticle interference scattering

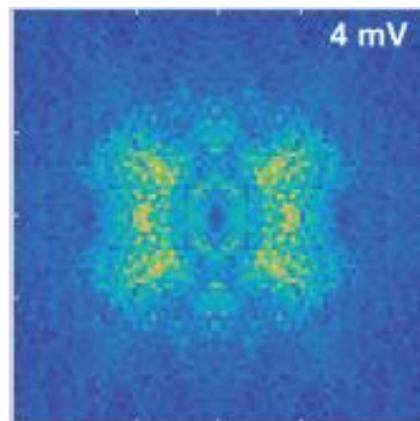
Topography



Spectroscopy

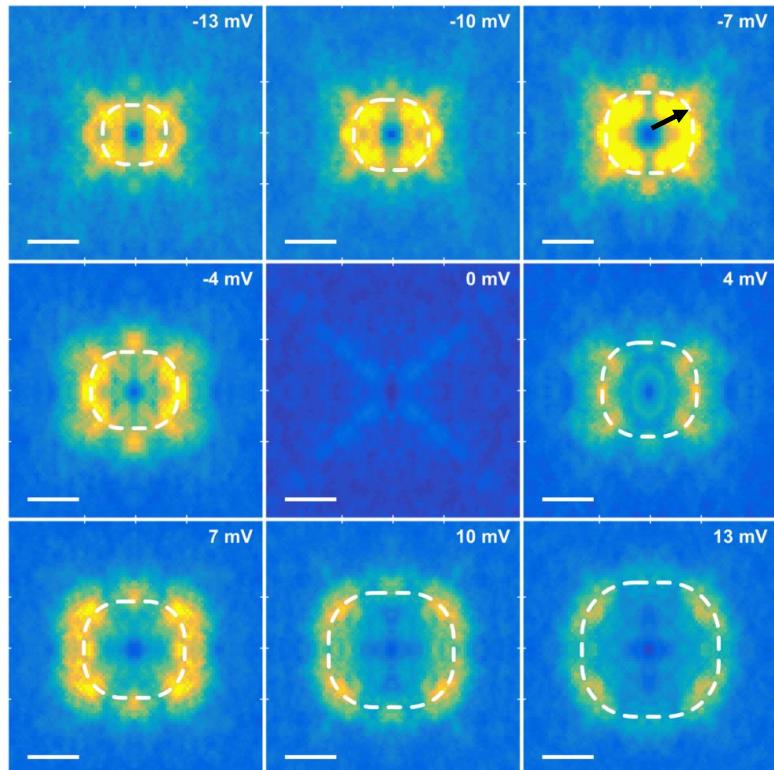


Fourier transform of the  
spectroscopy



## The CaKFe<sub>4</sub>As<sub>4</sub> system

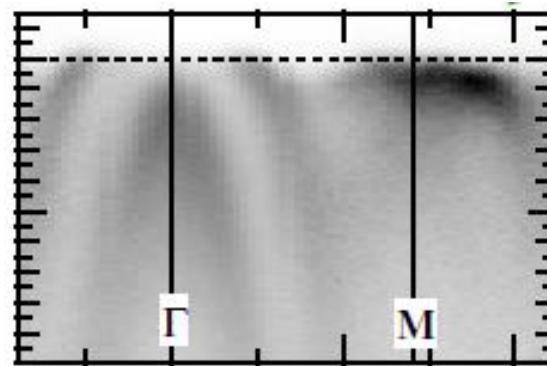
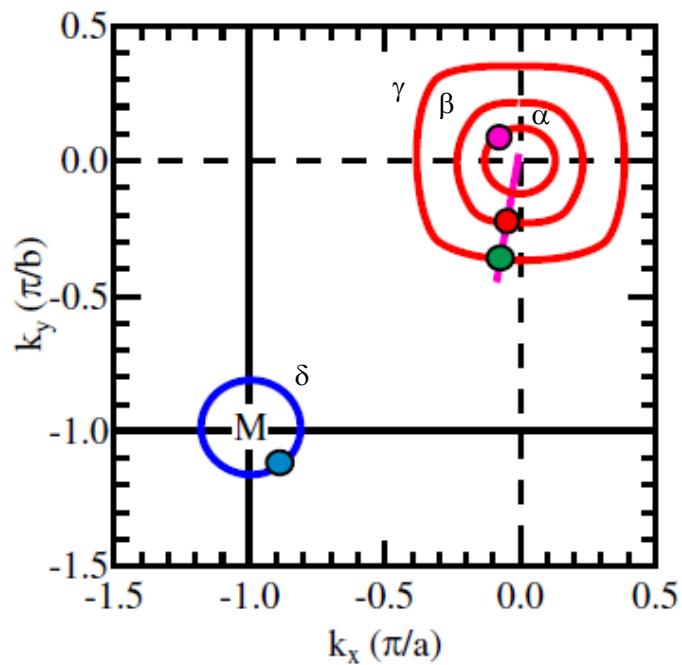
### Fourier transform of the spectroscopy



- One main scattering wavevector whose size increases with energy
- No scattering below about 4 mV

## The CaKFe<sub>4</sub>As<sub>4</sub> system

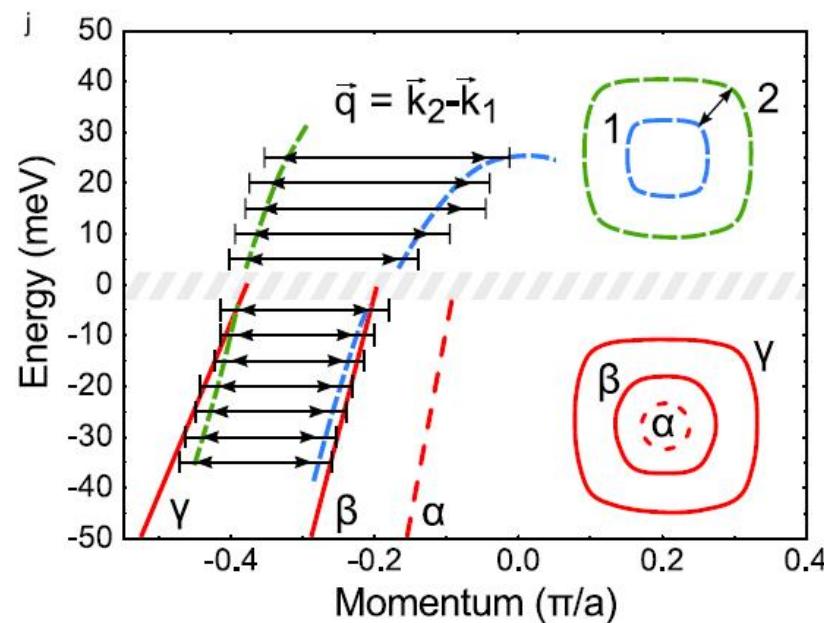
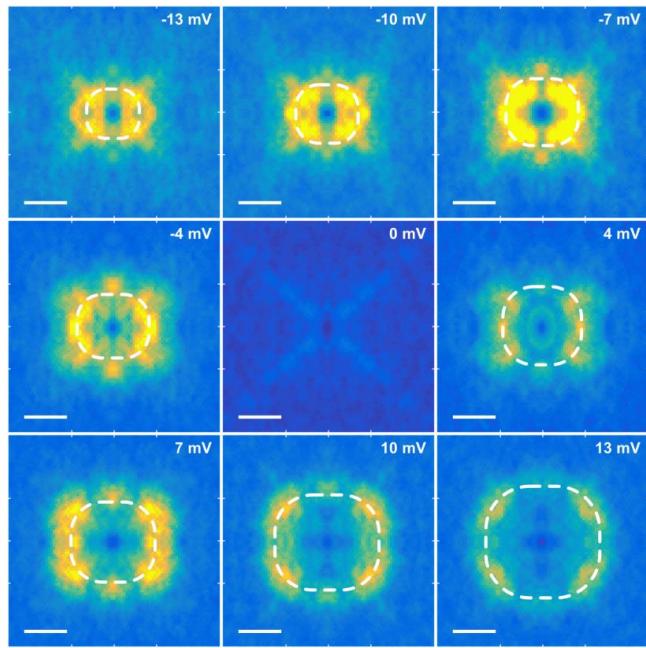
ARPES



Mou et al., PRL 117 277001 (2016)

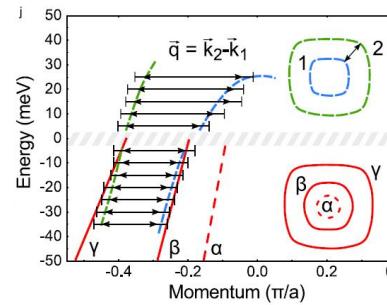
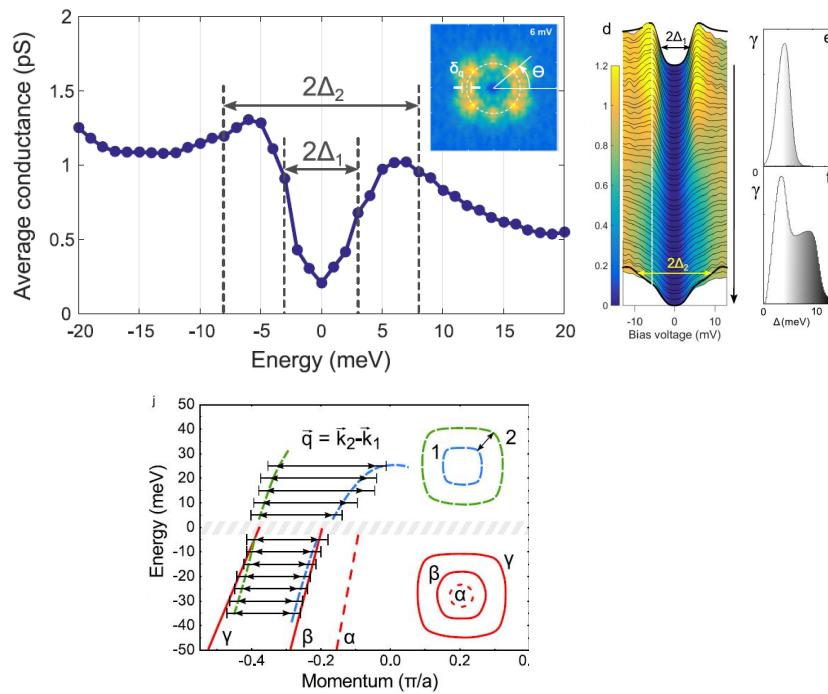
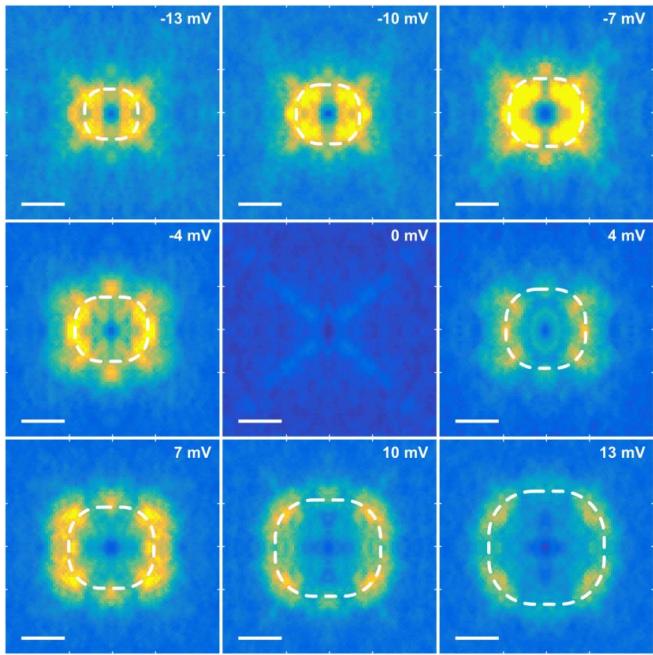
## The CaKFe<sub>4</sub>As<sub>4</sub> system

Electron dispersion relation from quasiparticle interference scattering  
Empty and full states.  
Innermost pocket closes a few mV above E<sub>F</sub>



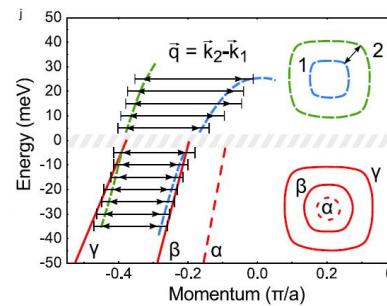
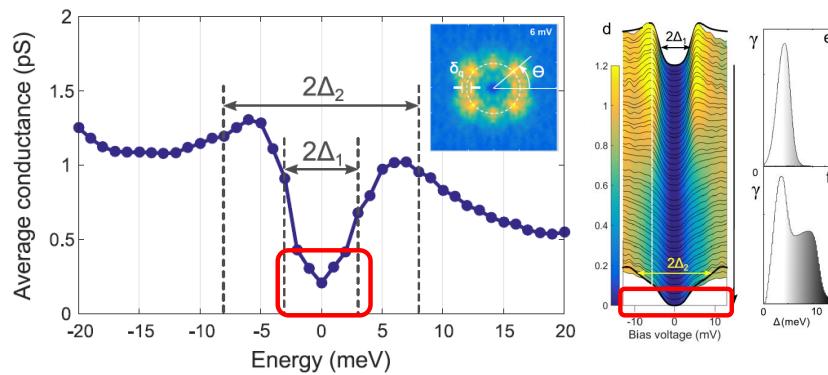
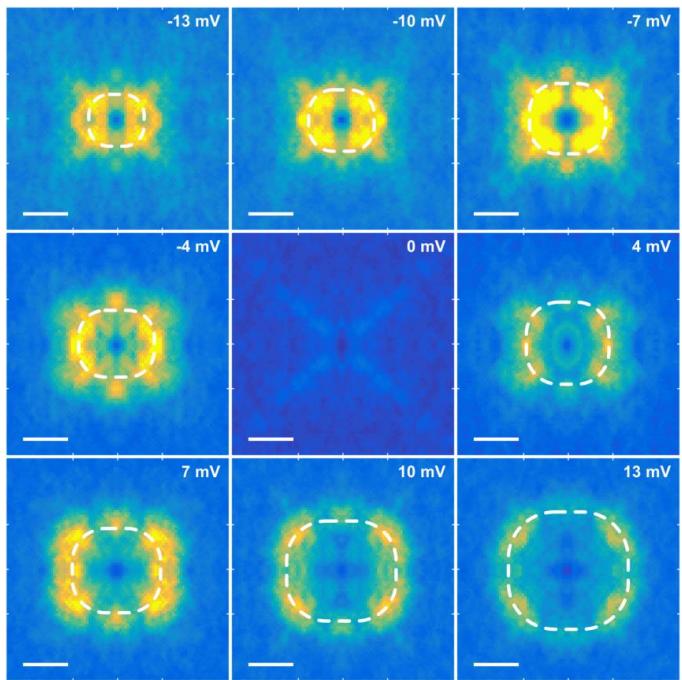
## The CaKFe<sub>4</sub>As<sub>4</sub> system

### Superconducting gap opening from quasiparticle interference scattering Isotropic gap opening in the hole sheets



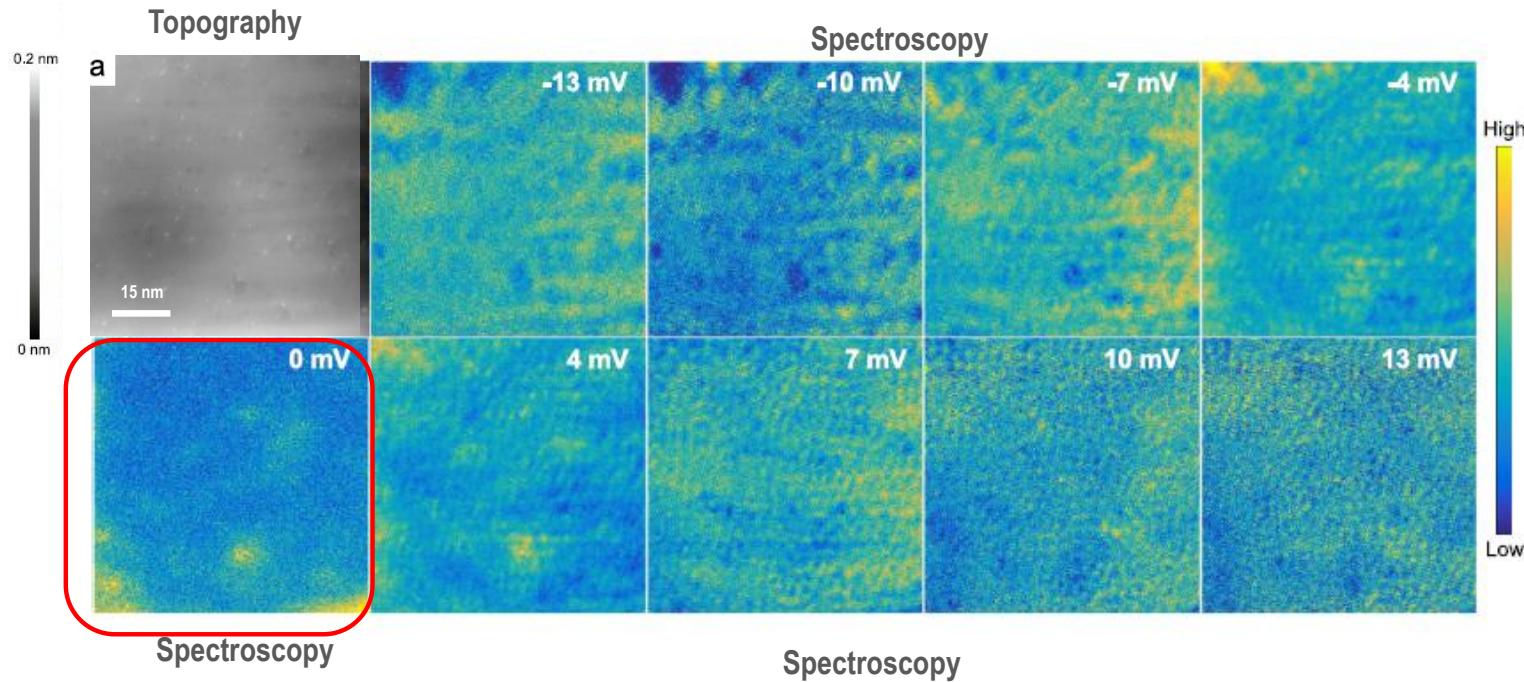
## The CaKFe<sub>4</sub>As<sub>4</sub> system

### Superconducting gap opening from quasiparticle interference scattering Zero bias DOS ??



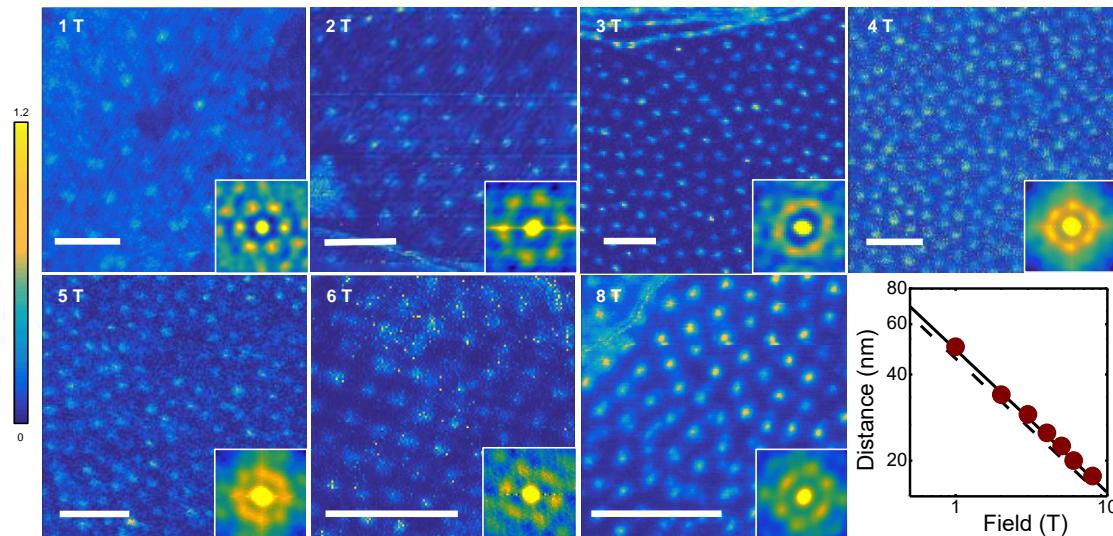
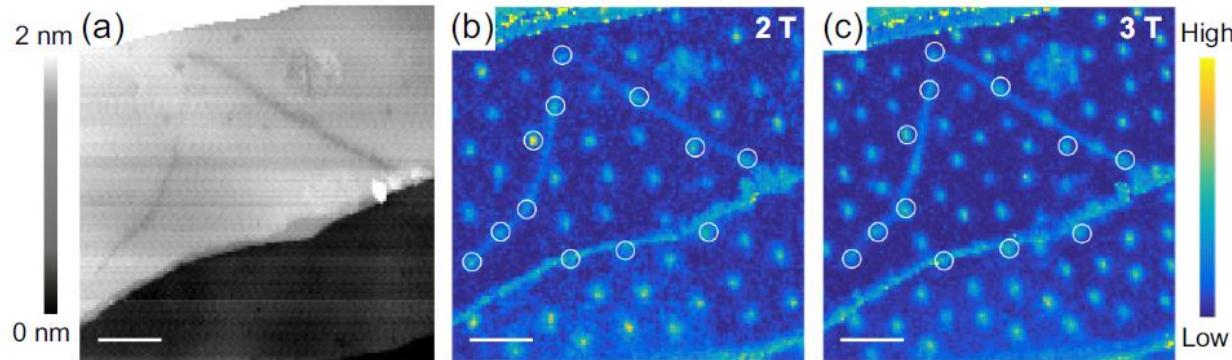
## The CaKFe<sub>4</sub>As<sub>4</sub> system

### Quasiparticle interference scattering at zero bias within the superconducting gap

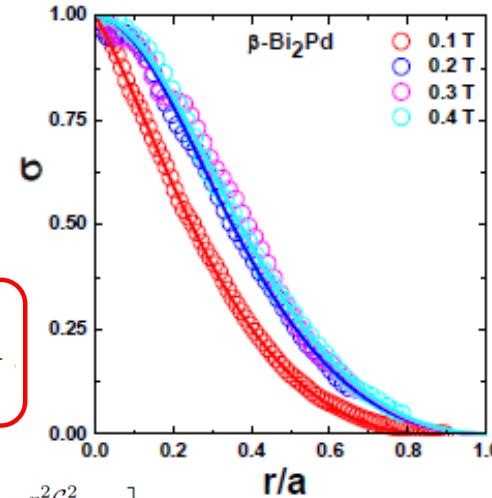
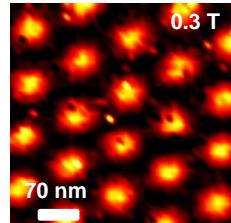
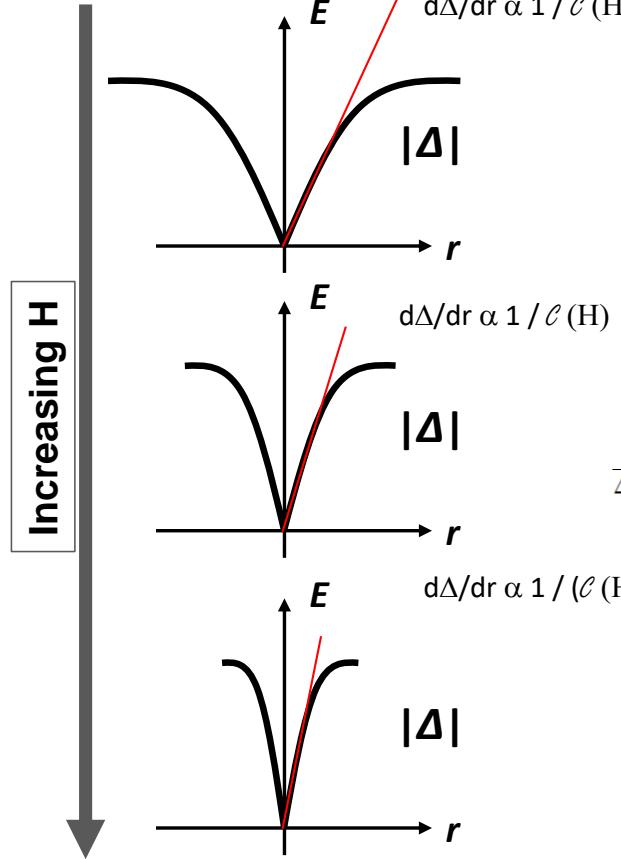


## The CaKFe<sub>4</sub>As<sub>4</sub> system

### Zero bias states and pinning



## Using de Gennes relation to understand $N(r)$ in a vortex core

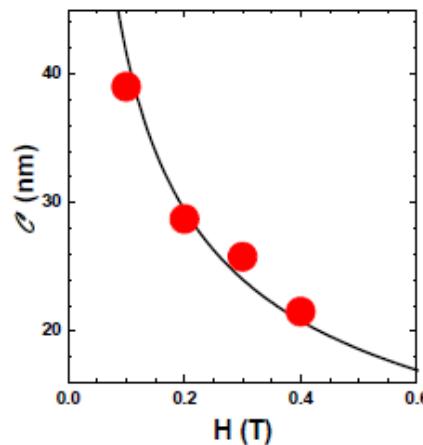


$$\frac{N(r)}{N_n} = 1 - \frac{|\Delta(r)|^2}{\Delta_0^2}$$

$$\frac{\Delta(r)}{\Delta_0(B, T)} = \frac{r}{\sqrt{r^2 + C^2}} \exp \left[ -\frac{r^2 C^2}{2a^2(C^2 + a^2)} \right]$$

$$\xi_H = \sqrt{\phi_0 / 2\pi H}$$

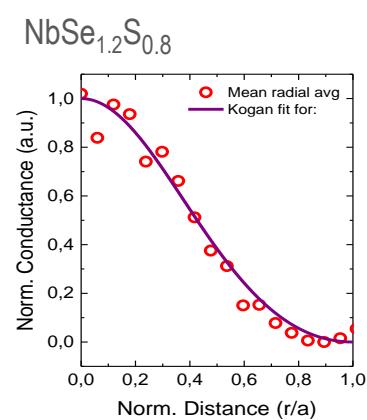
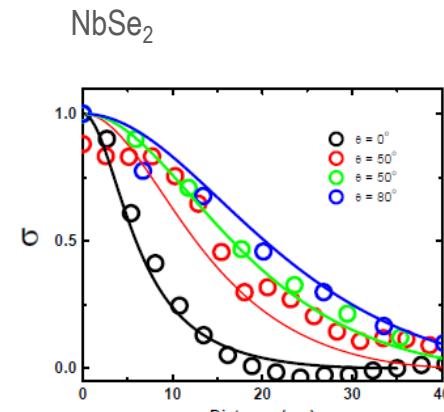
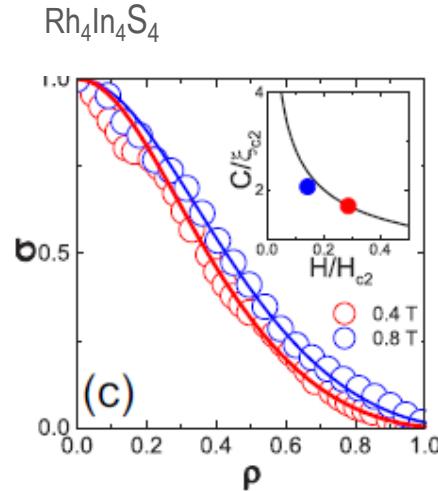
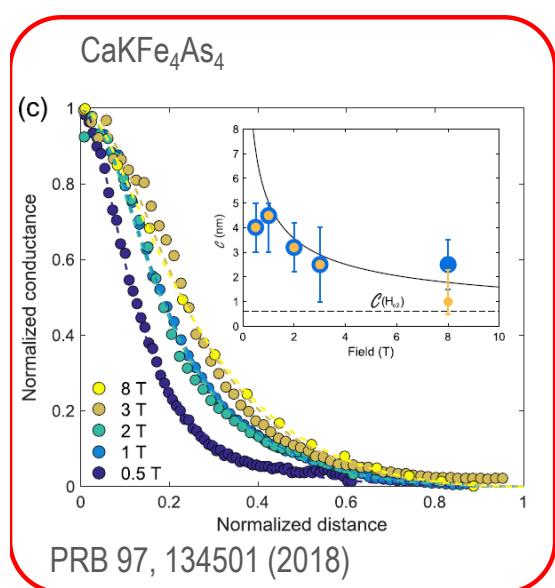
$$C = \eta a = \eta \sqrt{\frac{\phi_0}{\pi H}}$$



Field dependence of the vortex core size  
 V.G. Kogan and N.V. Zhelezina,  
 Phys. Rev. B, 71, 134505 (2005).

$$C = \eta a = \eta \sqrt{\frac{\phi_0}{\pi H}} = \eta \sqrt{2} \xi_H$$

## Applied to different superconductors



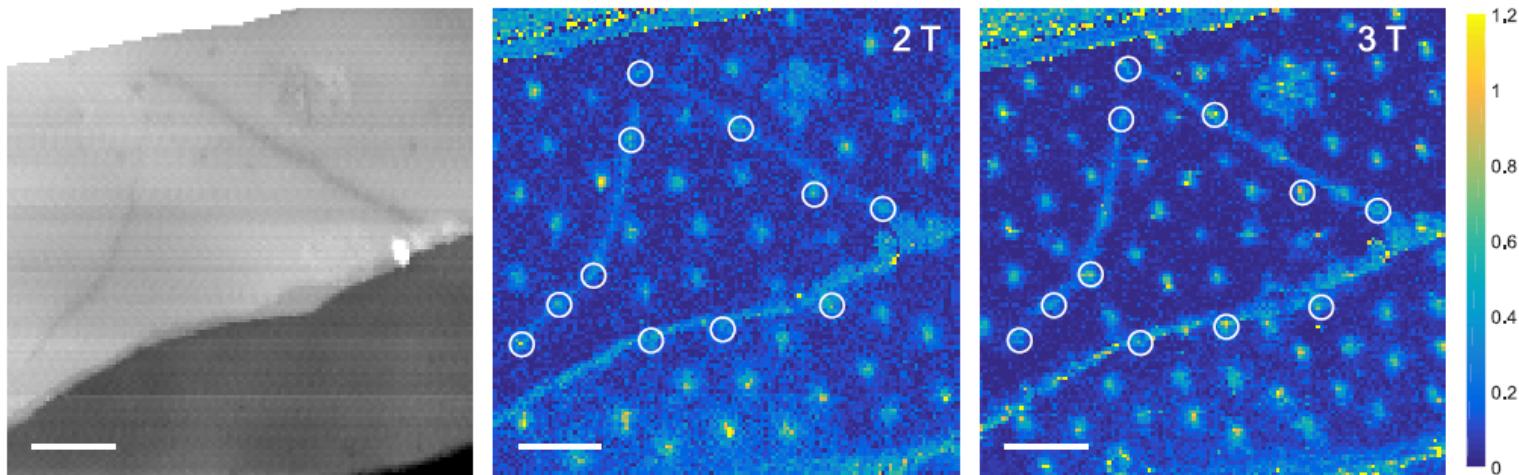
$$\frac{N(\mathbf{r})}{N_n} = 1 - \frac{|\Delta(\mathbf{r})|^2}{\Delta_0^2}$$

$$\frac{\Delta(r)}{\Delta_0(B, T)} = \frac{r}{\sqrt{r^2 + C^2}} \exp \left[ -\frac{r^2 C^2}{2a^2(C^2 + a^2)} \right]$$

Field dependence of the vortex core size probed by STM  
 A. Fente et al, Phys. Rev. B, 94, 014517 (2016)

## The CaKFe<sub>4</sub>As<sub>4</sub> system

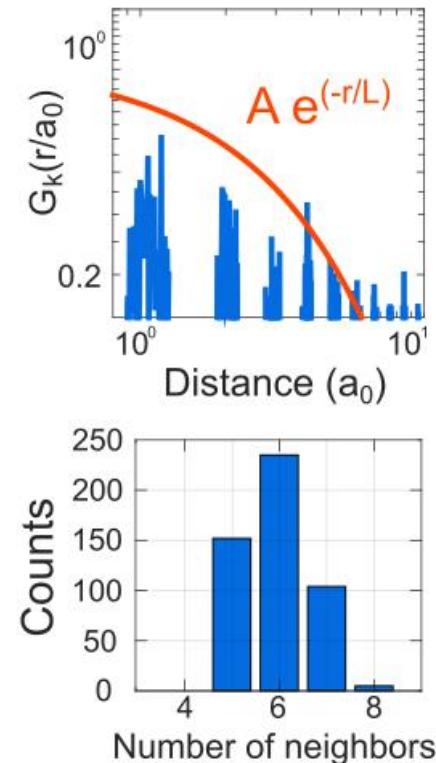
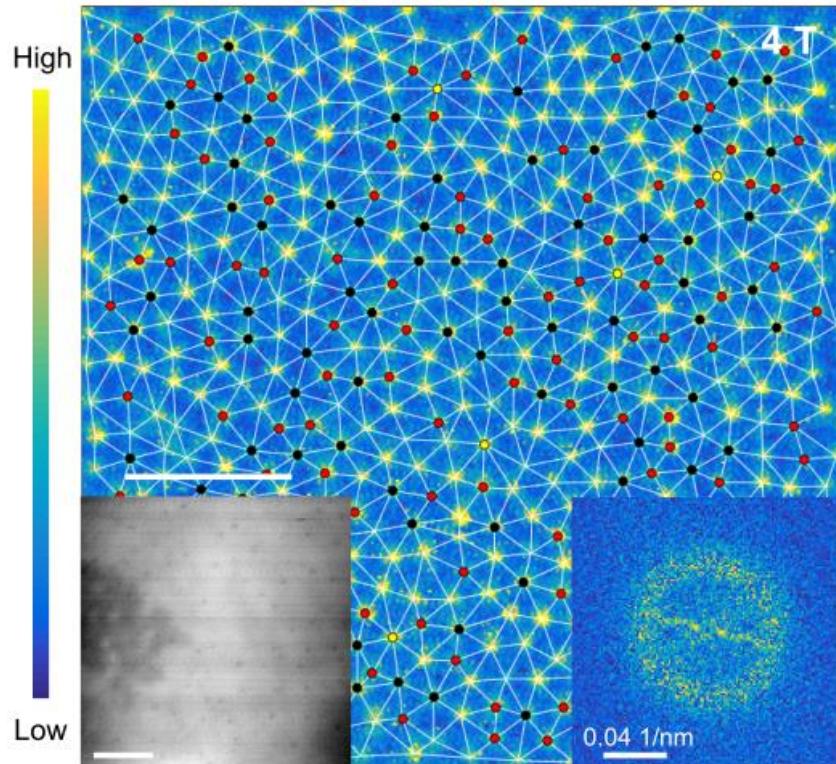
### Pair breaking and vortex pinning in CaKFe<sub>4</sub>As<sub>4</sub>



Core size  $\ell$   $\leftrightarrow$  size of pair breaking features  
Single vortex pinning

Field dependence of the vortex core size probed by STM  
A. Fente et al, Phys. Rev. B, 94, 014517 (2016)

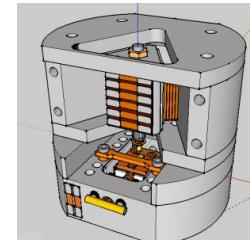
## Disordered vortex lattice



## Real space imaging of the superconducting vortex lattice: Recent results and prospects



- Real space imaging and spectroscopy of the vortex lattice
- Superconducting CaKFe<sub>4</sub>As<sub>4</sub>, bandstructure, vortex core, vortex lattice and pinning
- Vortex lattice structure: interactions vs disorder
- Observation of vortex creep on cooling
- Josephson effect at atomic scale

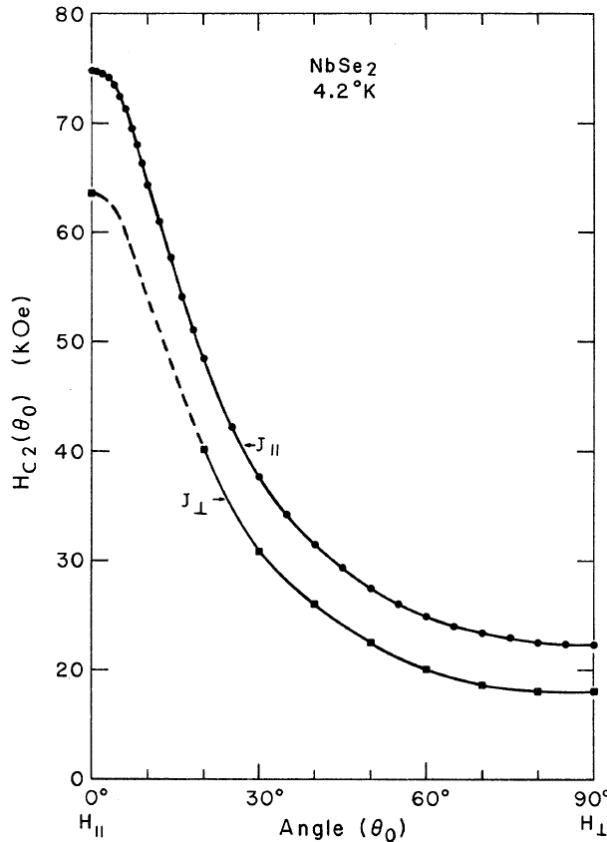


## Upper critical field of the anisotropic superconductor 2H-NbSe<sub>2</sub>

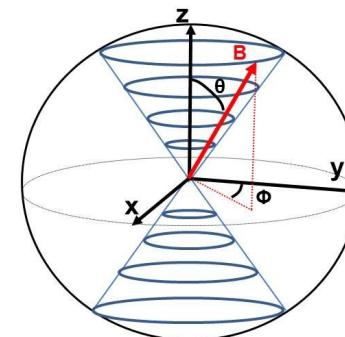
$T_c = 7.2 \text{ K}$

$$H_{C2\perp} = 4.2T$$

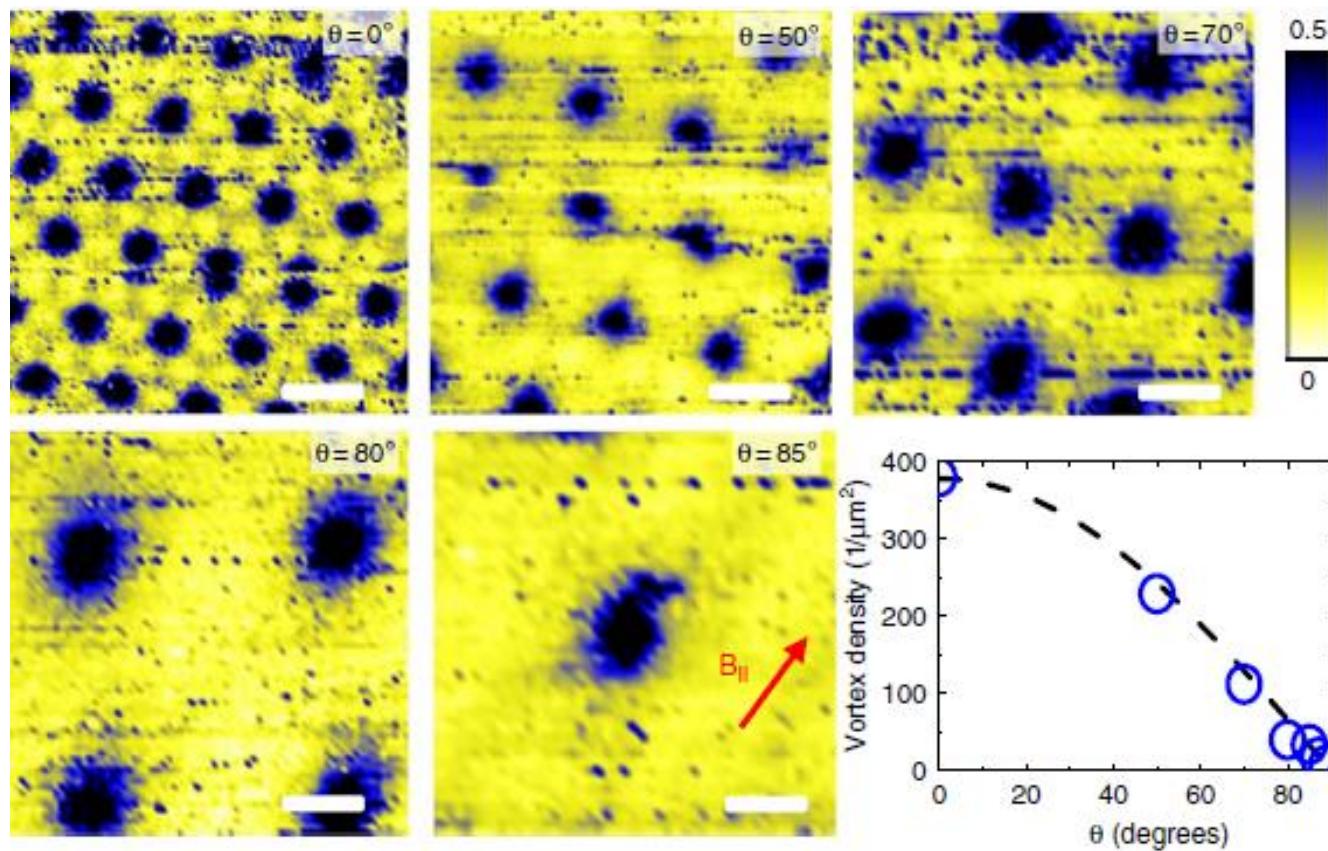
$$H_{C2\parallel} = 12.7T$$



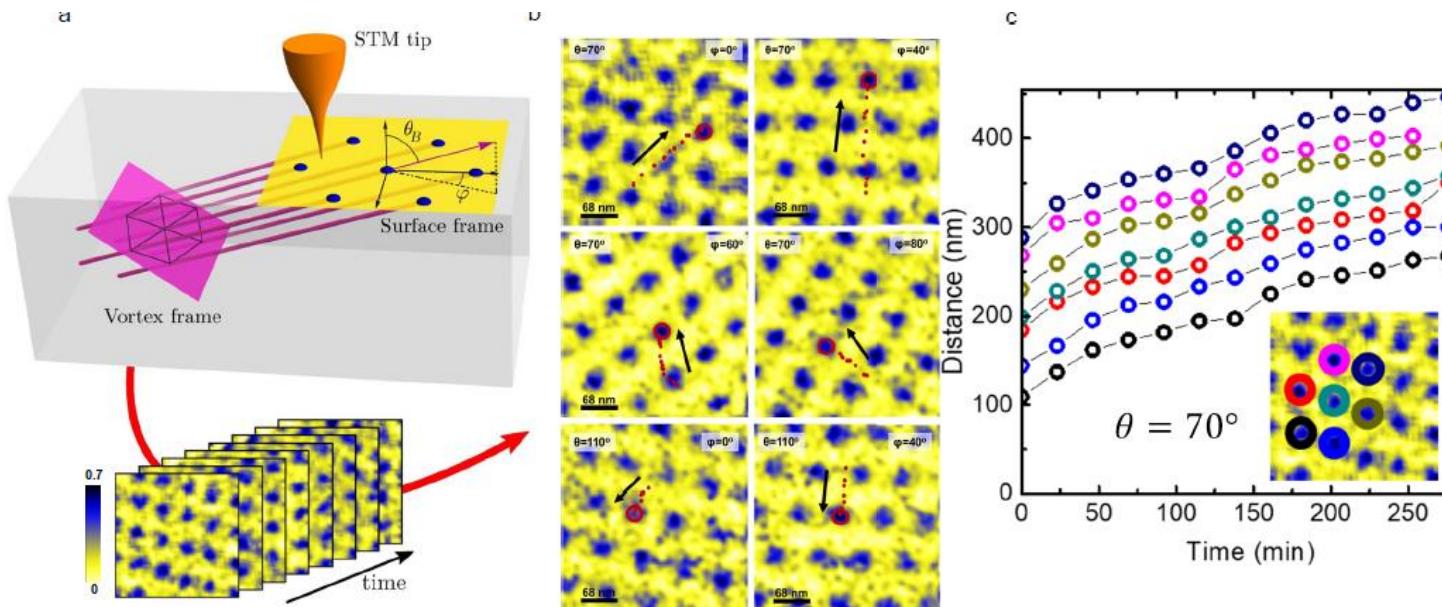
$$\Gamma = \frac{H_{C2\parallel}}{H_{C2\perp}} = \frac{\xi_{\parallel}}{\xi_{\perp}} \approx 3$$



## 2H-NbSe<sub>2</sub>: Vortex lattice in tilted magnetic fields

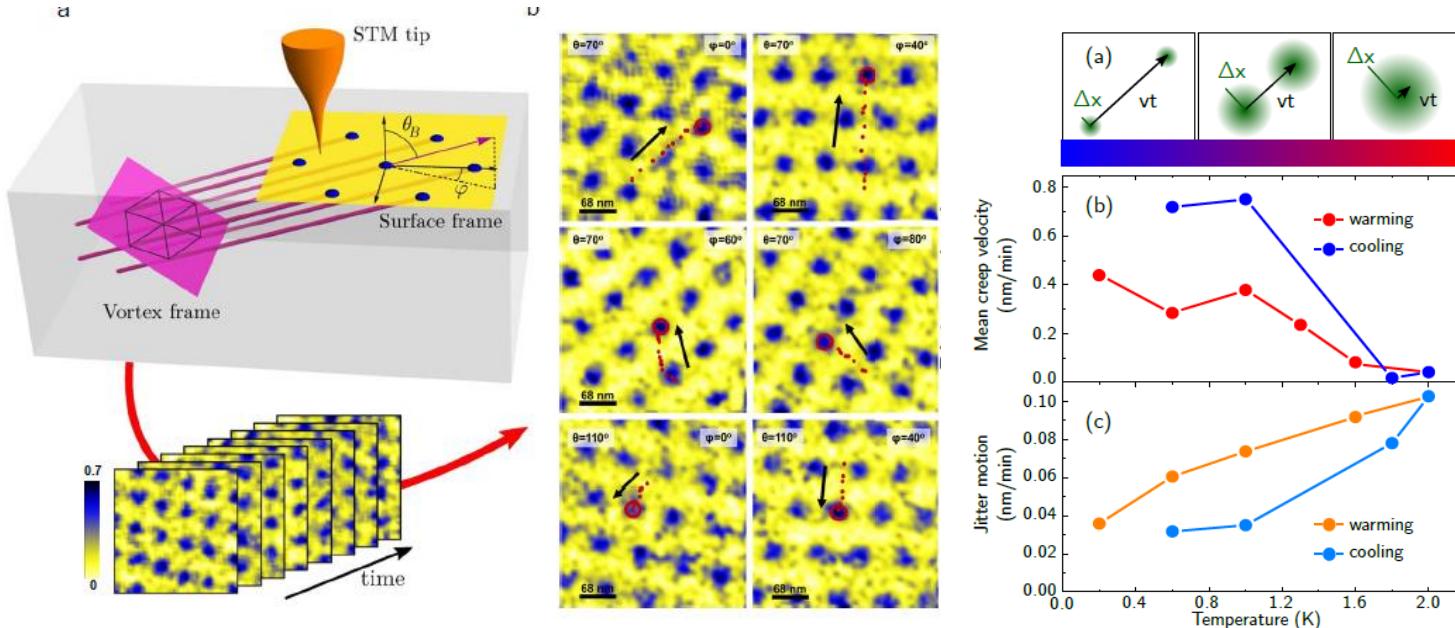


## 2H-NbSe<sub>2</sub>: Vortex lattices spontaneously move sometimes in tilted magnetic fields



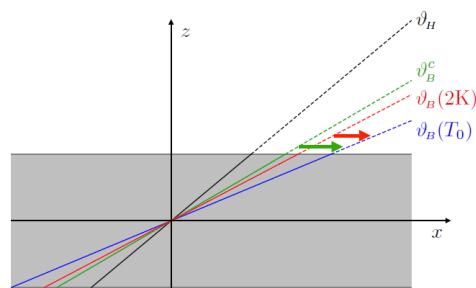
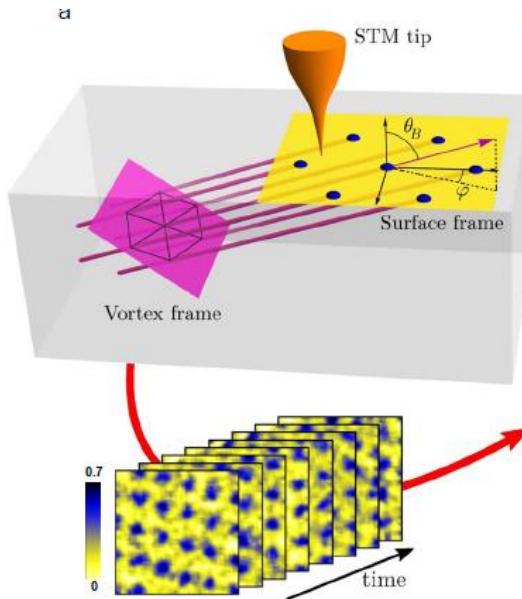
Slow vortex creep  
Washboard potential

## Creeping lattices when cooling

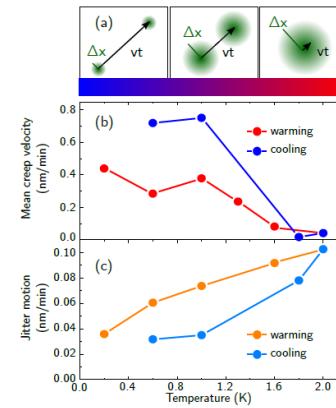


Increasing temperature stops creep  
Decreasing temperature again -> creep restarts

## Creeping lattices when cooling



$$\sin(\theta_B - \theta_H) = \frac{H_{c1}}{H} \frac{(1 - \varepsilon^2) \sin \theta_B \cos \theta_B}{(\varepsilon^2 \sin^2 \theta_B + \cos^2 \theta_B)^{1/2}}$$



R. Willa  
 (Wednesday, 10h30, PC3-1)

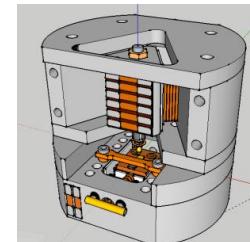
**Increasing temperature stops creep  
 Decreasing temperature again -> creep restarts**

Thermal creep emerging from cooling a tilted vortex lattice in uniaxial superconductor  
 R. Willa et al., arxiv 1906.06294

## Real space imaging of the superconducting vortex lattice: Recent results and prospects



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***Thank you very much for  
your attention***

V. Barrena, F. Martín Vega

E. Herrera, J. Benito

B. Wu

J.G. Rodrigo, S. Vieira

I. Guillamón



C. Munuera, F. Mompean and M. García-Hernández

ICMM-CSIC

W. Meier, S. Bud'ko and P.C. Canfield

V.G. Kogan

Ames Lab

D. Aoki, G. Knebel, JP Brison and J Flouquet

Grenoble

Y. Anahory, Weizmann, Hebrew University

Y. Fasano, Bariloche

M. Milosevic, Antwerp



S. Park, A. Martín, A. Levy-Yeyati

J. Herrera William  
Bogotá



A. I. Buzdin  
Bordeaux

K. Kadowaki  
Tsukuba





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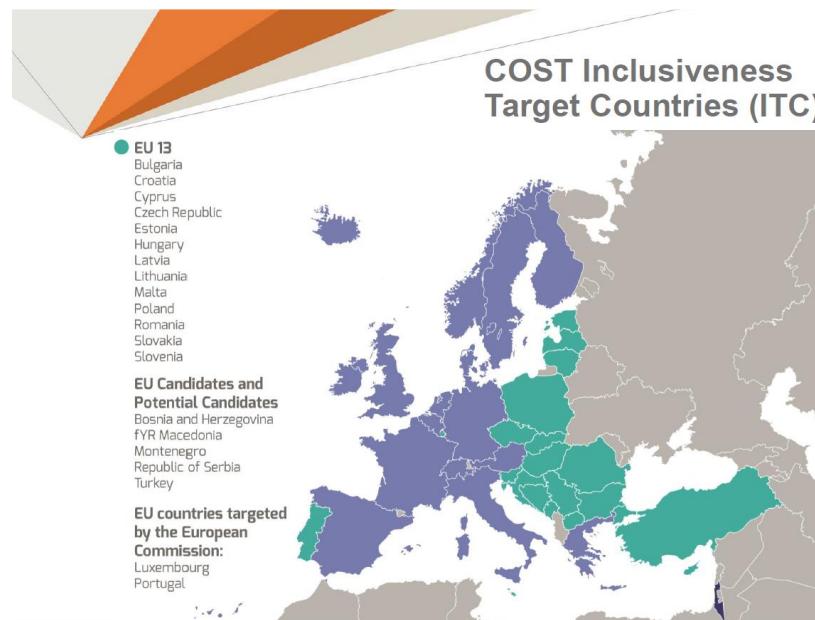
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Dimitri Roditchev and Brigitte Leridon, March 2018



### School on quantum materials and workshop on vortex behavior in unconventional superconductors

7-12 October, 2018



Probing Coherent  
Superconducting Hybrids at  
the Nanoscale, 17-20  
February 2019

13th November 2018    Meetings

Program and schedule The final schedule and abstract book are now available. Important dates February 17th: arrival in Eilat for dinner (around...



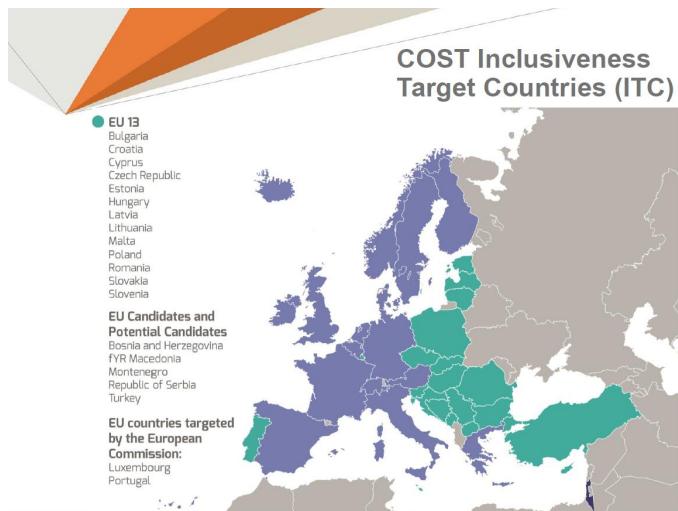
Yony Anahory and Beena Kalisky, February 2019



Joint workshop between  
MOLSPIN and  
NANOCOHBRI –  
Superconductivity meets  
Molecular Spins – 20-22  
March 2019

CA COST Action CA16218

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- *Open network*
- *Meetings (2/year)*
- *Short term stays*

[https://www.cost.eu/actions/CA16218/#tabs\\_1](https://www.cost.eu/actions/CA16218/#tabs_1)  
Name:management-committee

[www.nanocohy bri.eu](http://www.nanocohy bri.eu)



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*Short term stays*

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