

# The temperature dependence of superconducting single photon detectors is a vortex effect

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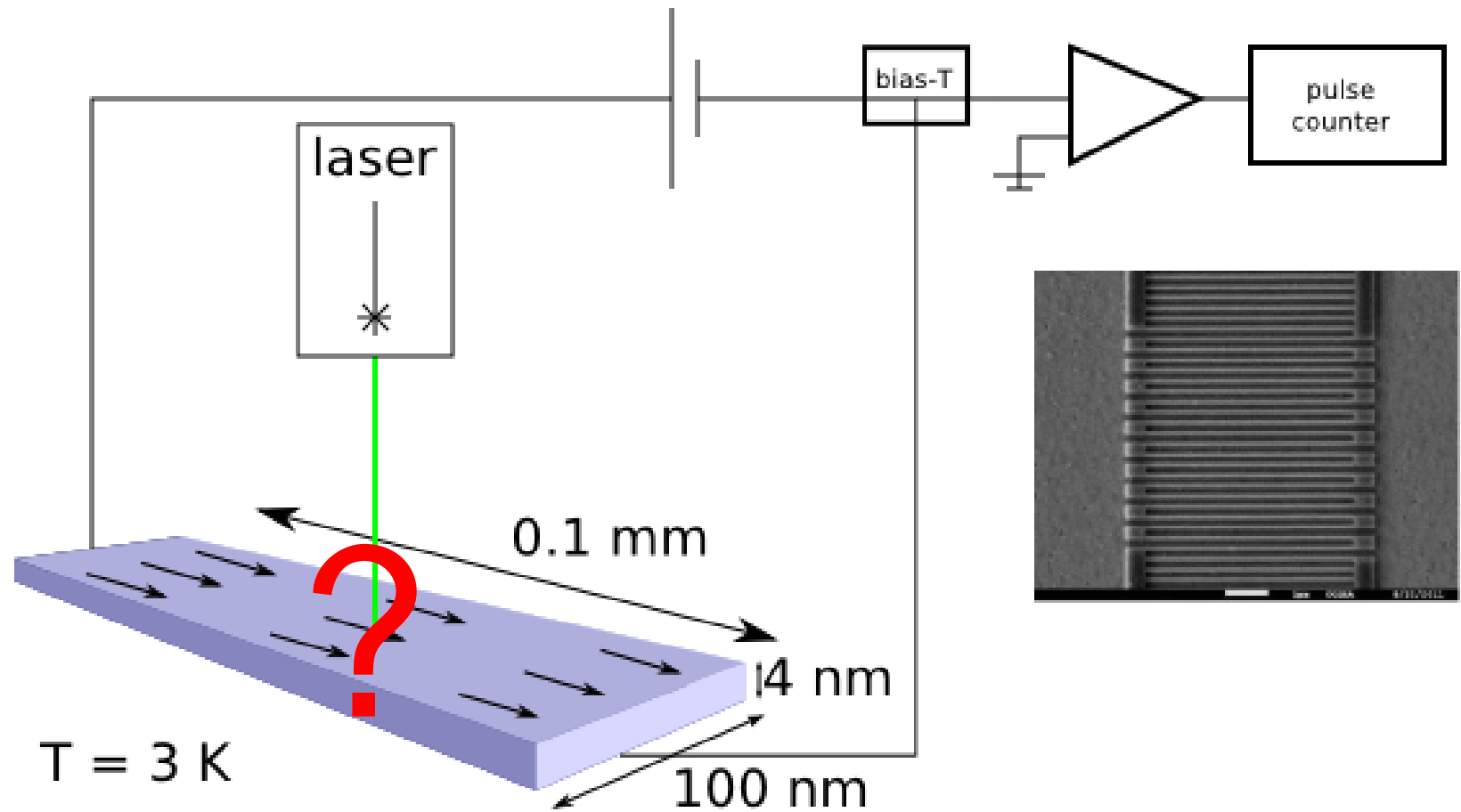
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Zili Zhou<sup>2</sup>, Alesandro Gaggero<sup>3</sup>, Francesco Mattioli<sup>3</sup>, Roberto Leoni<sup>3</sup>,  
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<sup>2</sup>) Cobra Research Institute, Eindhoven, the Netherlands

<sup>3</sup>) IFN, Rome, Italy

# Goal: investigate SSPD fundamentals



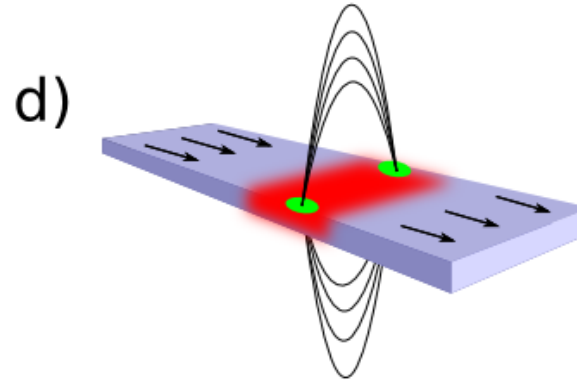
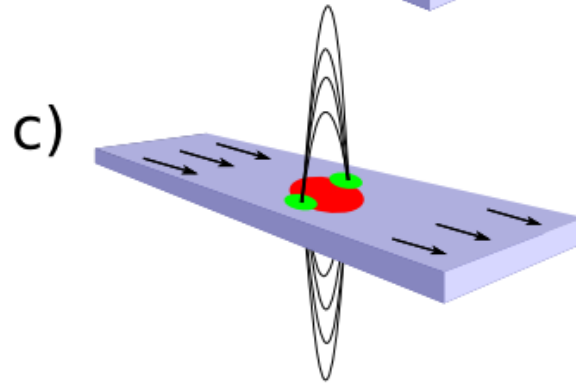
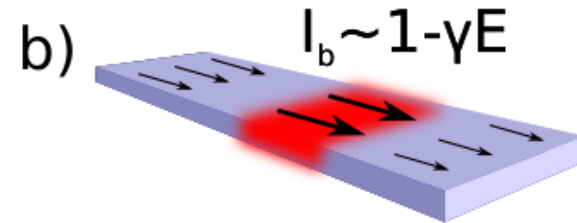
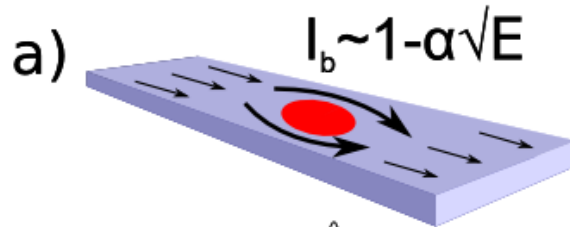
# Four models

Yes      ↓      Normal state?      ↓      No

No  
→

Vortices?

Yes  
→

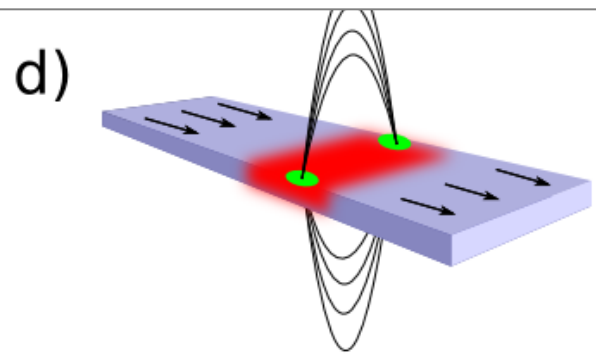
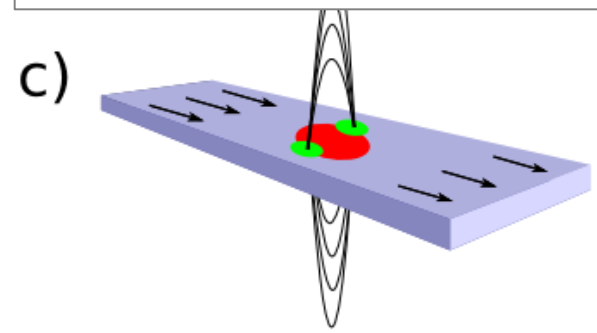


# Four models

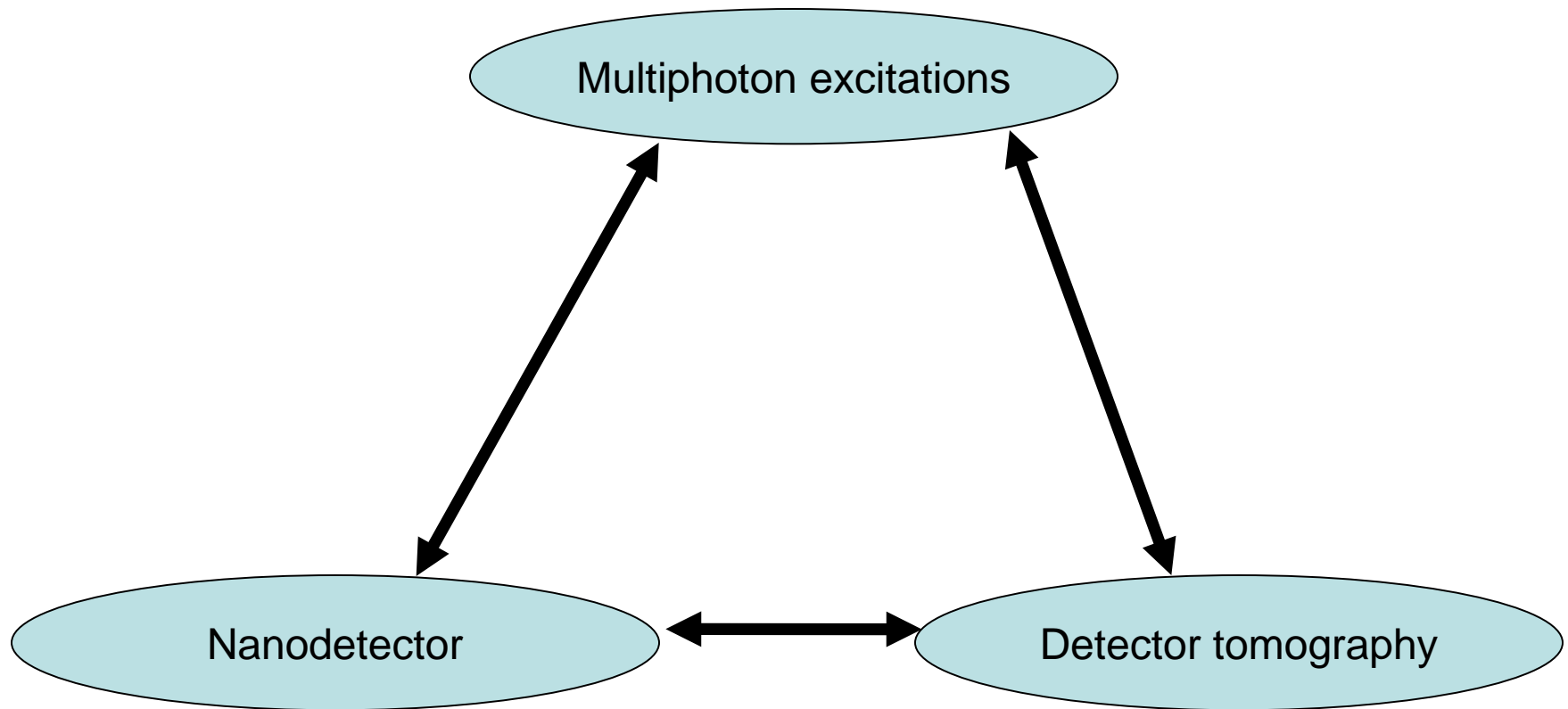
Yes ↓ Normal state? ↓ No

No →  
Vortices?  
Yes →

Energy-current relationship is key!



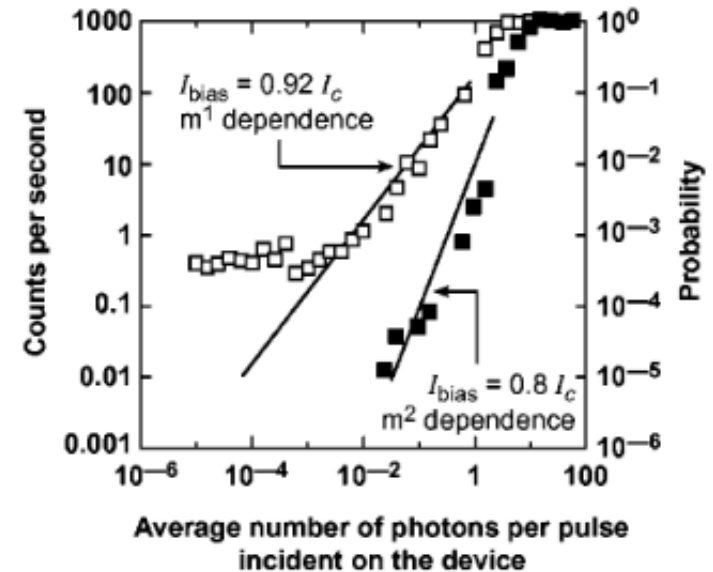
# Three inter-related techniques:



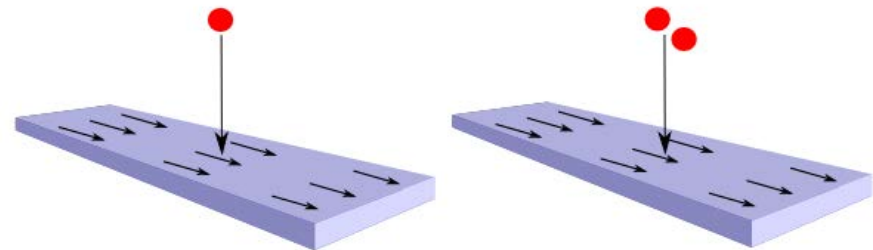
# Multiphoton excitations

Appl. Phys. Lett., Vol. 79, No. 6, 6 August 2001

- Observed in 2001 [1], but considered a curiosity
- Important experimental tool:
  - Enhanced dynamic range
  - Probe with multiple energies in a single experiment

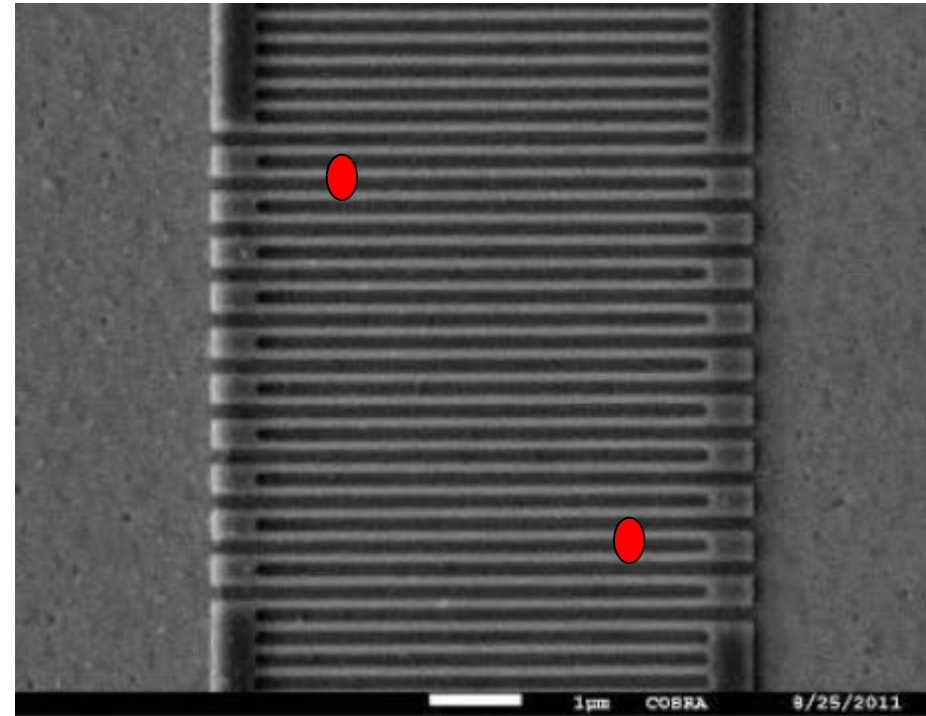


[1] Goltsman APL **79** (2001)



# How to study multiphoton excitations?

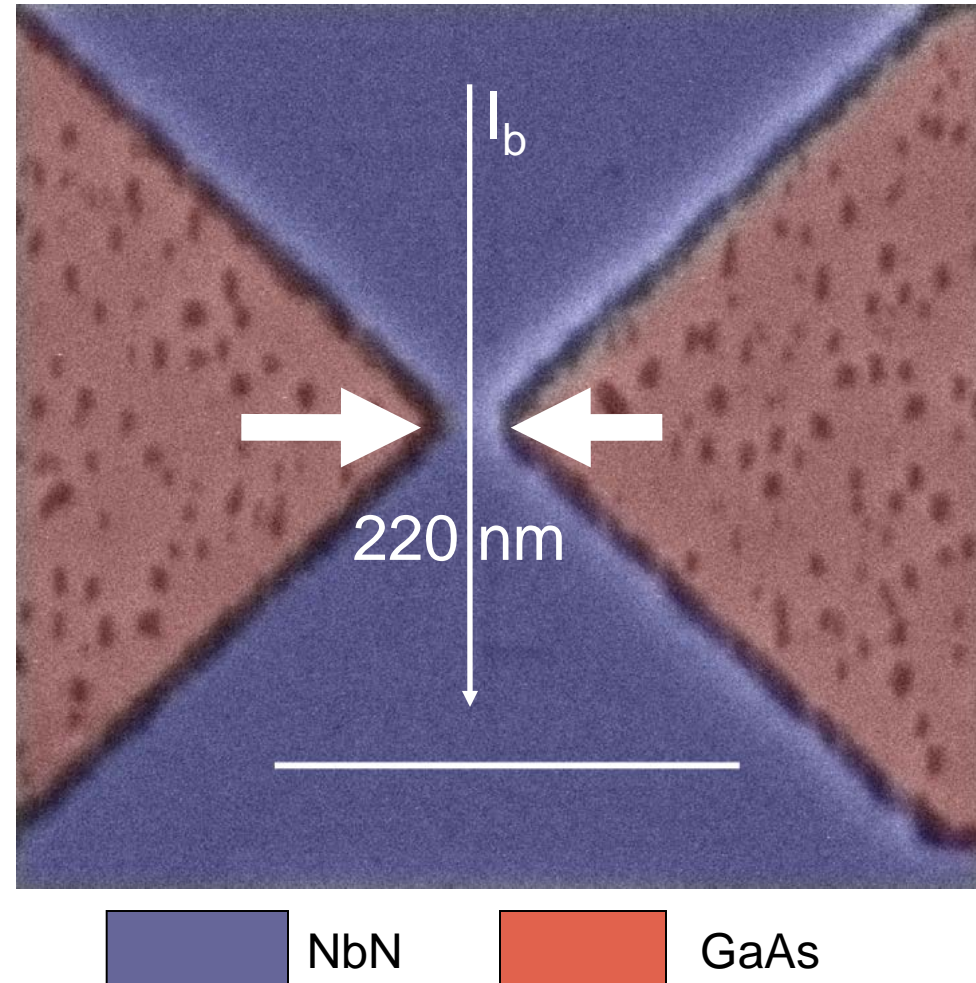
- Exist in meander, but suppressed due to geometry
- Furthermore: meander has:
  - Bends
  - ‘Constrictions’
- Fundamental study, so efficiency not an issue





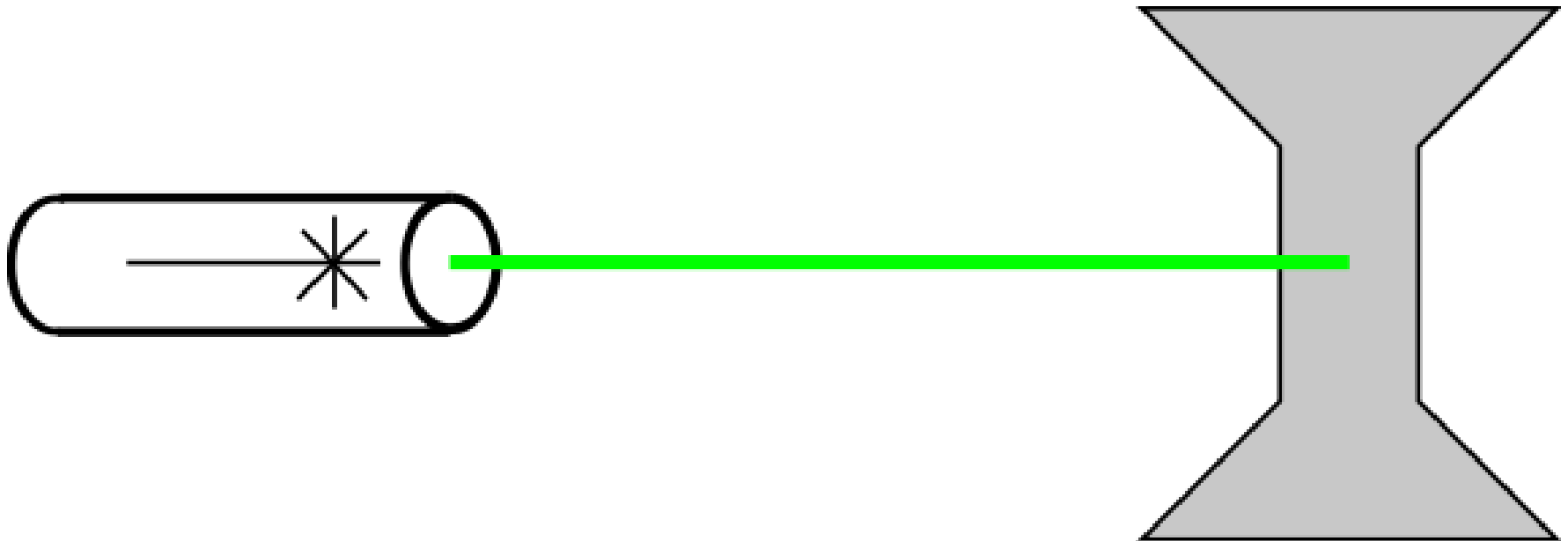
# Our sample: nanodetector

- One active point, 150, 220 nm wide NbN on GaAs (5 nm)
- Simple geometry
- Few fabrication errors
- Several multiphoton processes at once



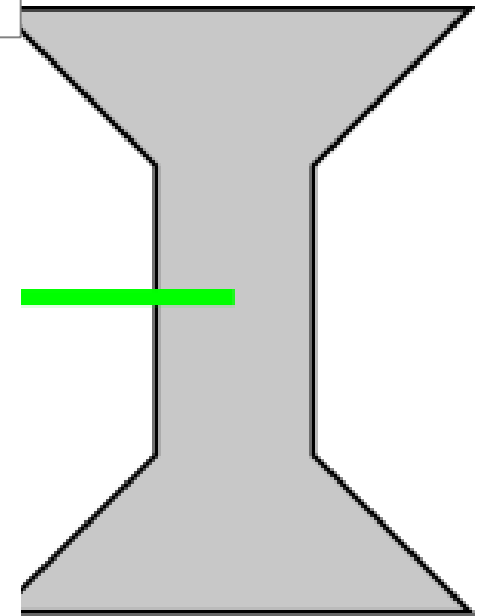
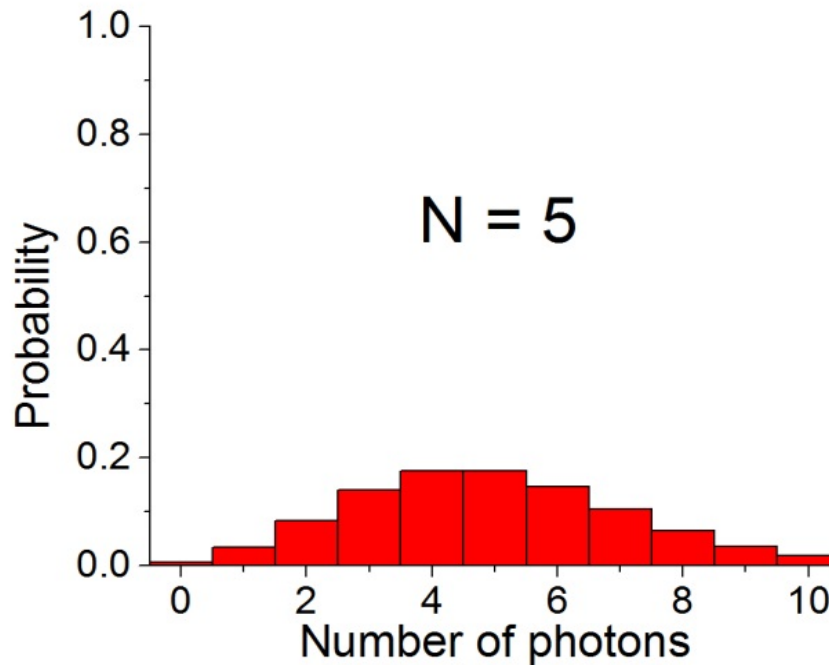
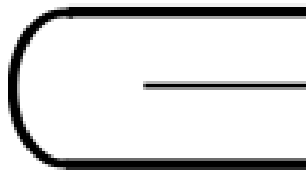


# How do you make multiphoton excitations?



# How do you make multiphoton excitations?

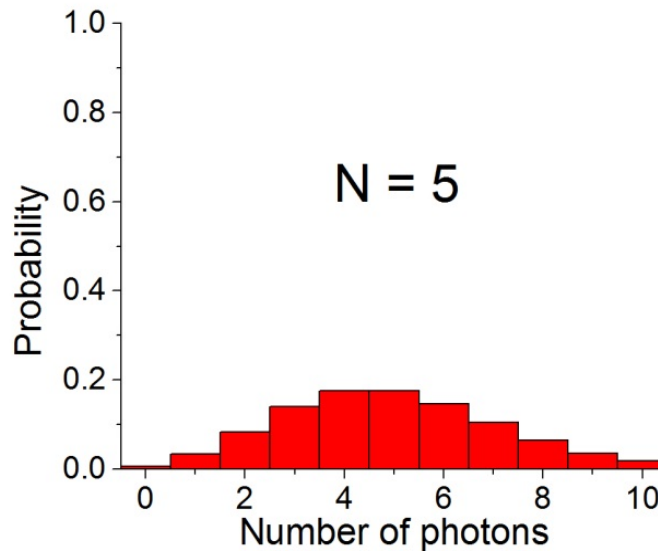
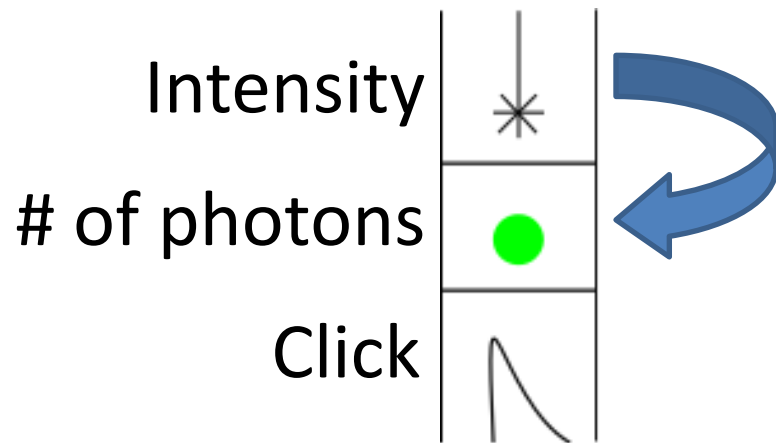
Bright laser pulses



# Quantum detector tomography

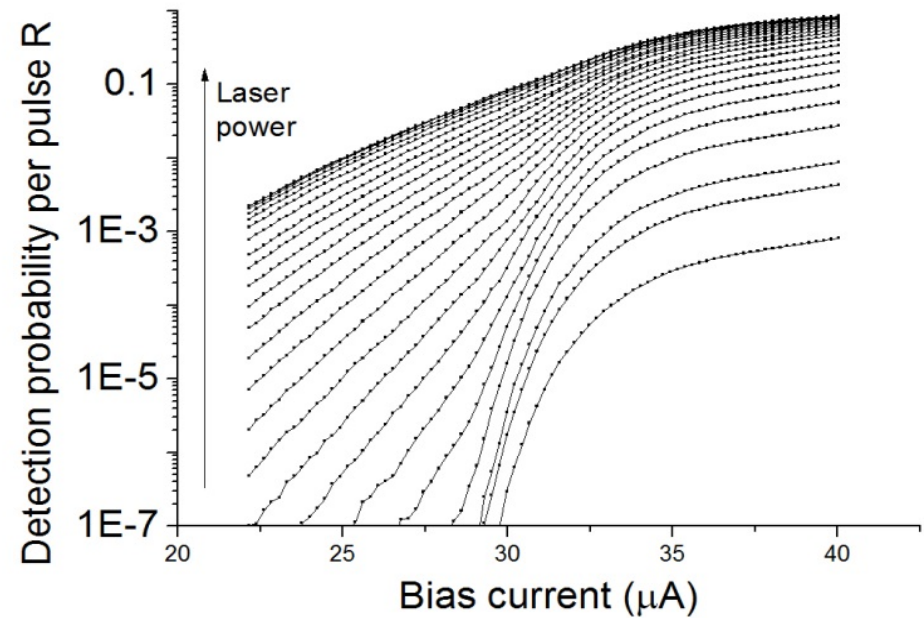
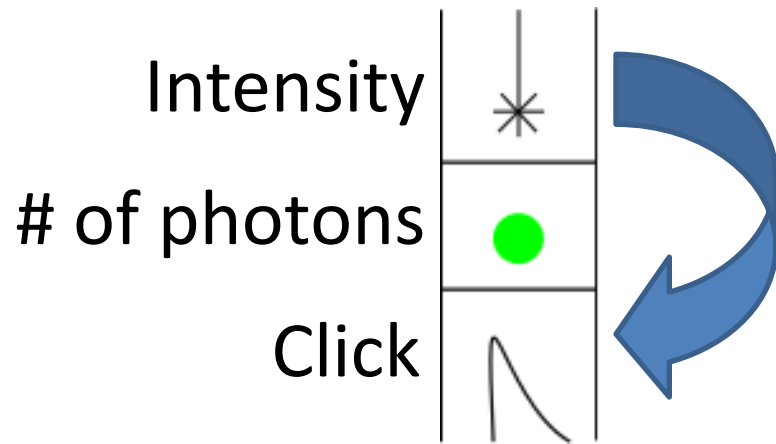
QDT is the bookkeeping of photon number probabilities,  
click probabilities and detection probabilities

Known from theory

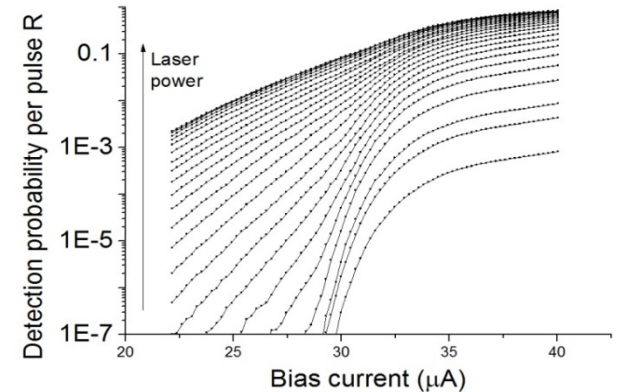
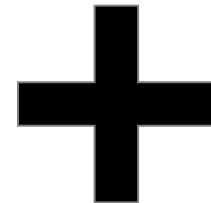
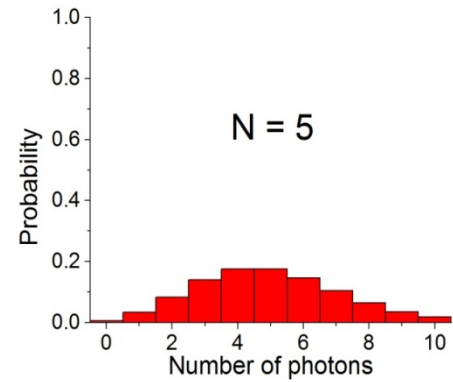
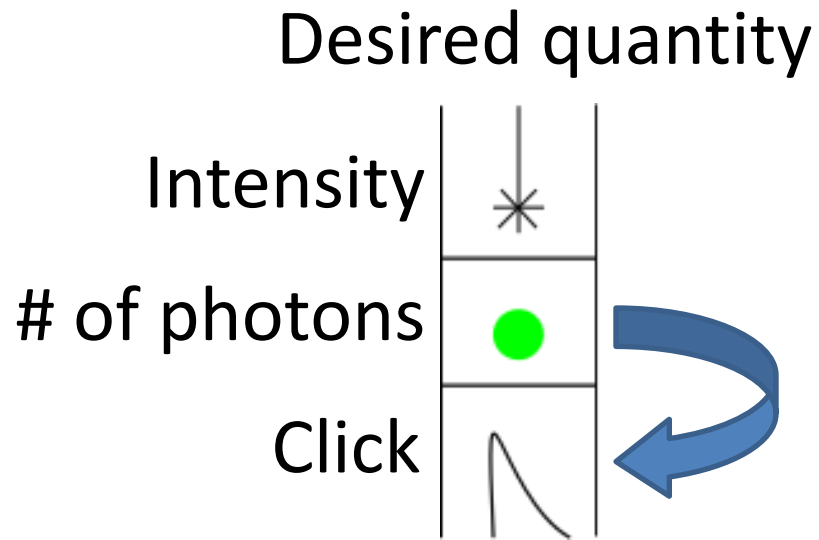


# Quantum detector tomography

Measured experimentally



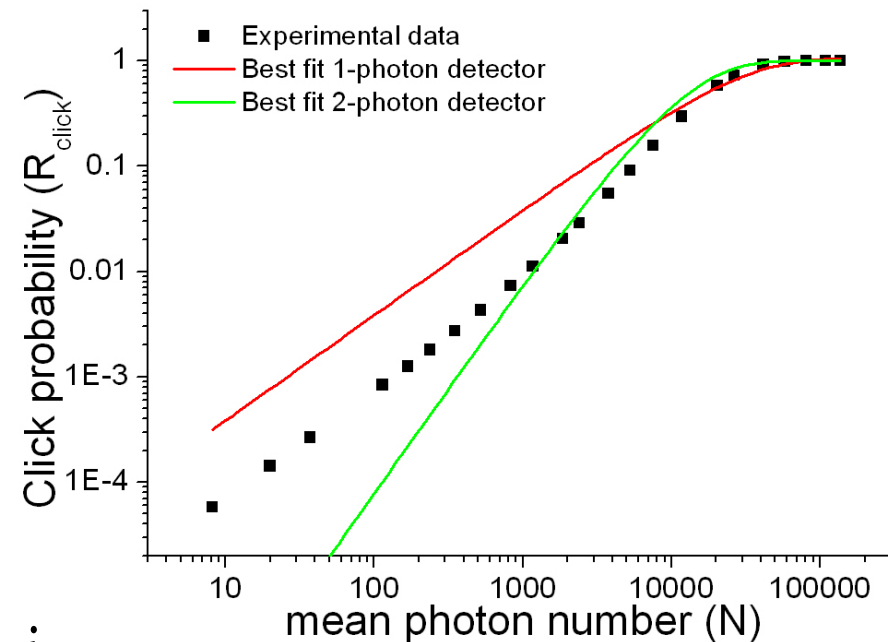
# Quantum detector tomography



# Detector Tomography

- Measure counts vs input intensity
- Response to  $i$  photons given by  $p_i$
- Treat linear efficiency separately, but as free parameter

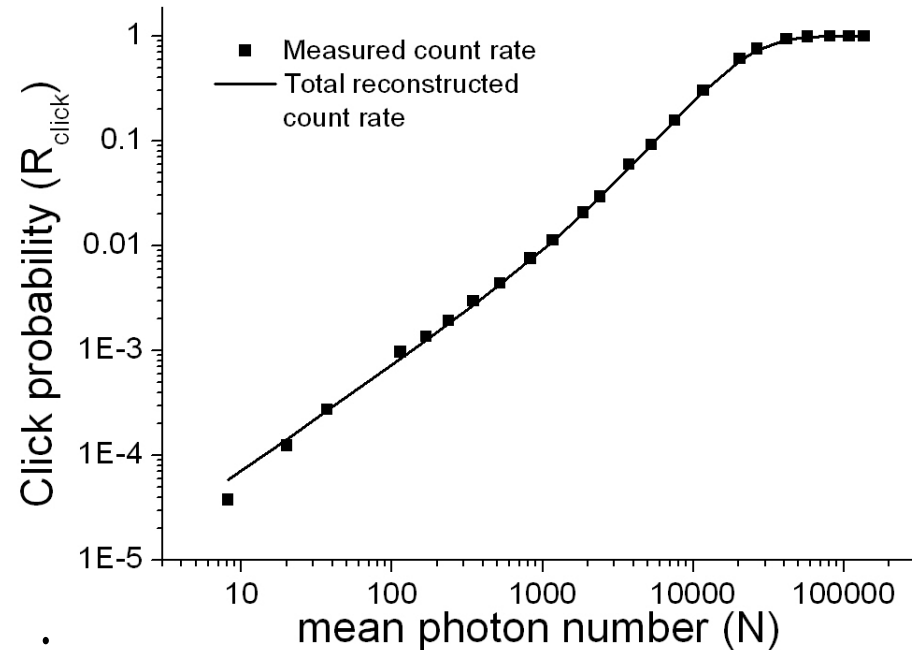
$$R(N) = e^{-\eta N} \sum_i p_i \frac{(\eta N)^i}{i!}$$



# Detector Tomography

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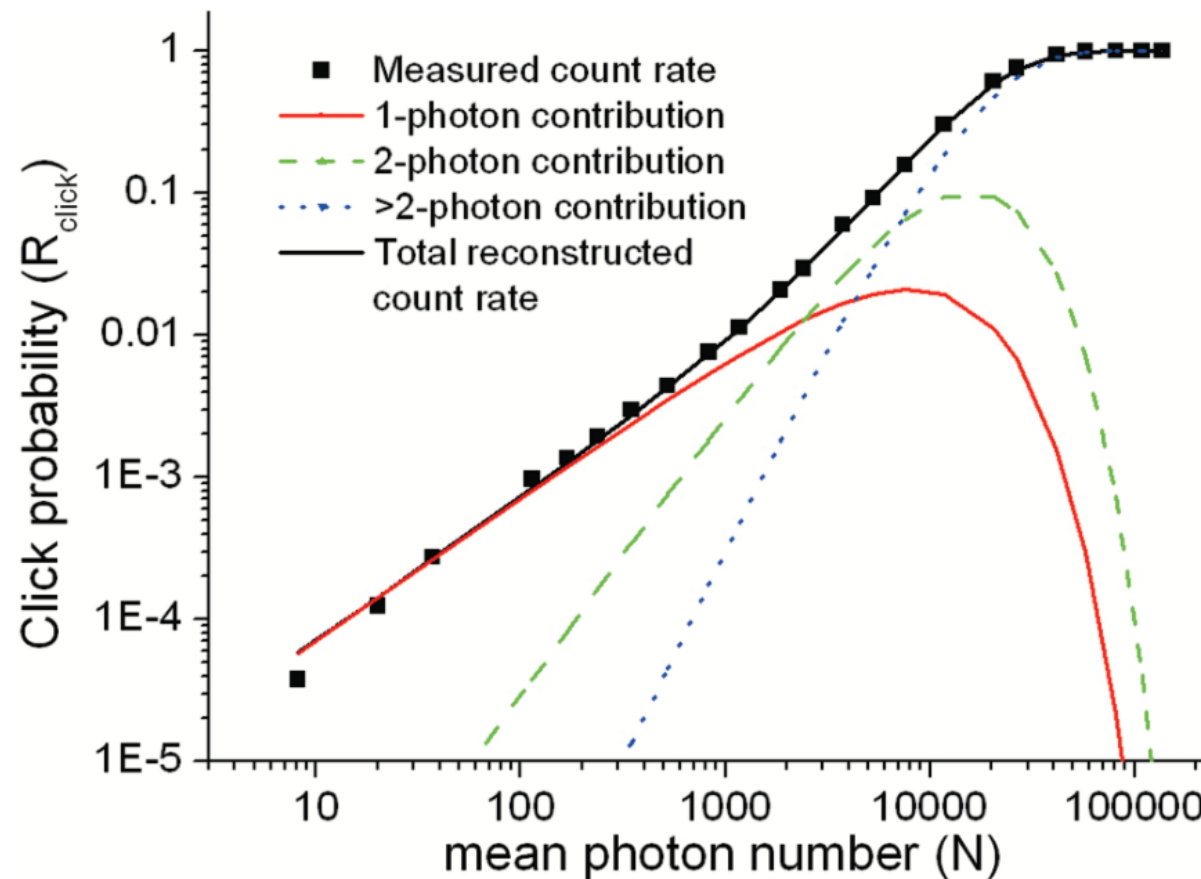
$$R(N) = e^{-\eta N} \sum_i p_i \frac{(\eta N)^i}{i!}$$





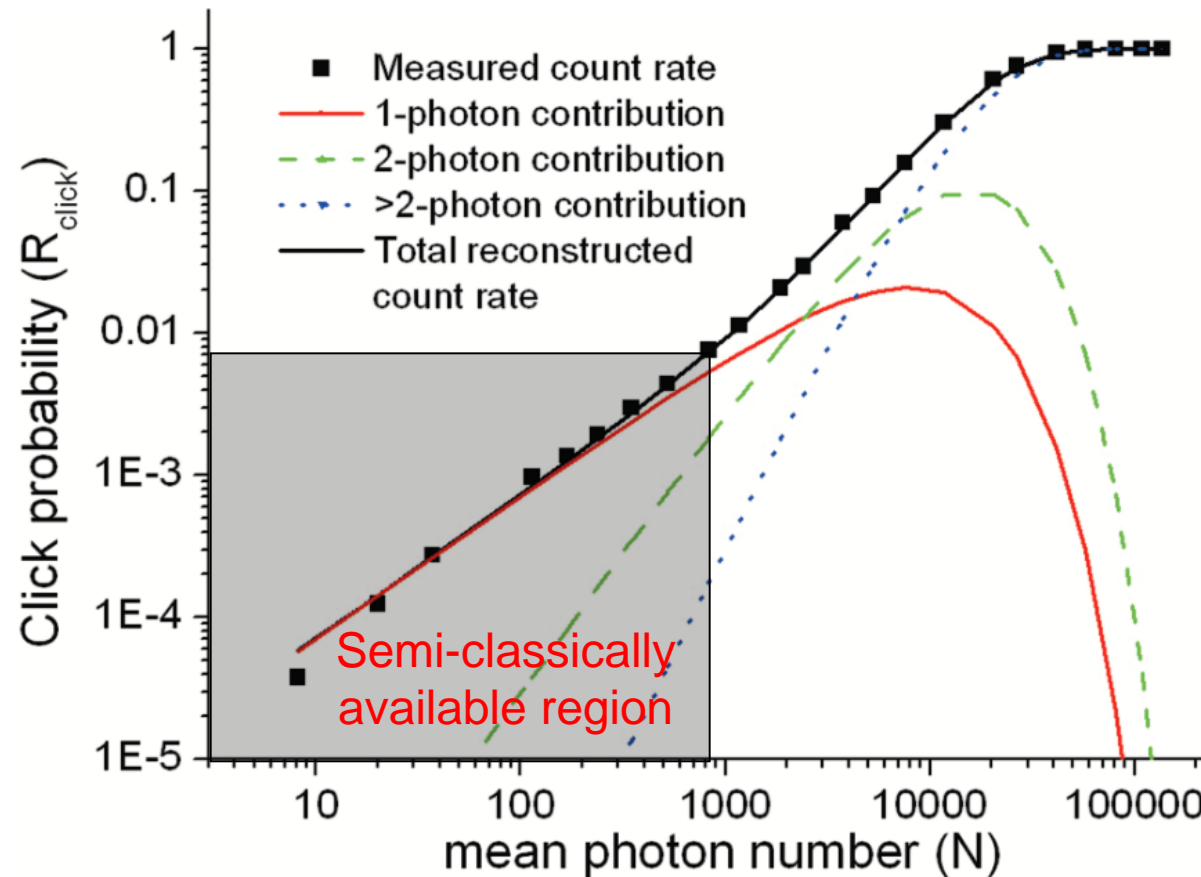
# Complete tomography

- 1, 2 photon processes present



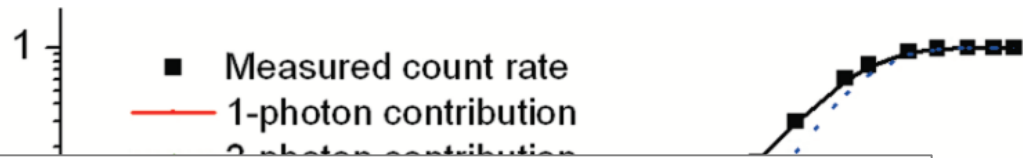
# Complete tomography

- 1, 2 photon processes present
- Usual method  $R = (\eta N)^i$  restricted to  $\eta N \ll 1$ , lowest  $i$



# Complete tomography

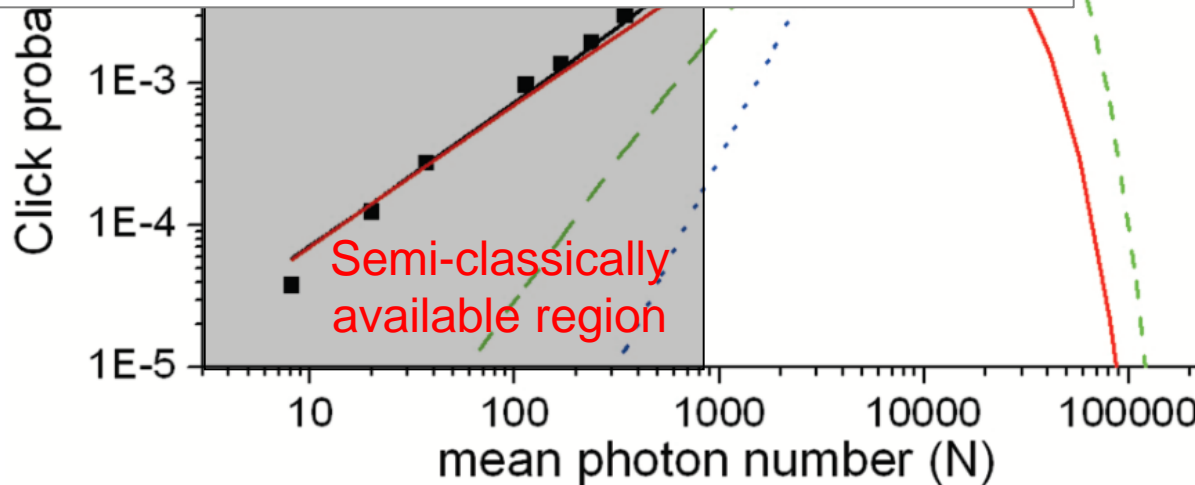
- 1, 2 photon processes



QDT can quantify multiphoton processes

- We are open to sharing this technique

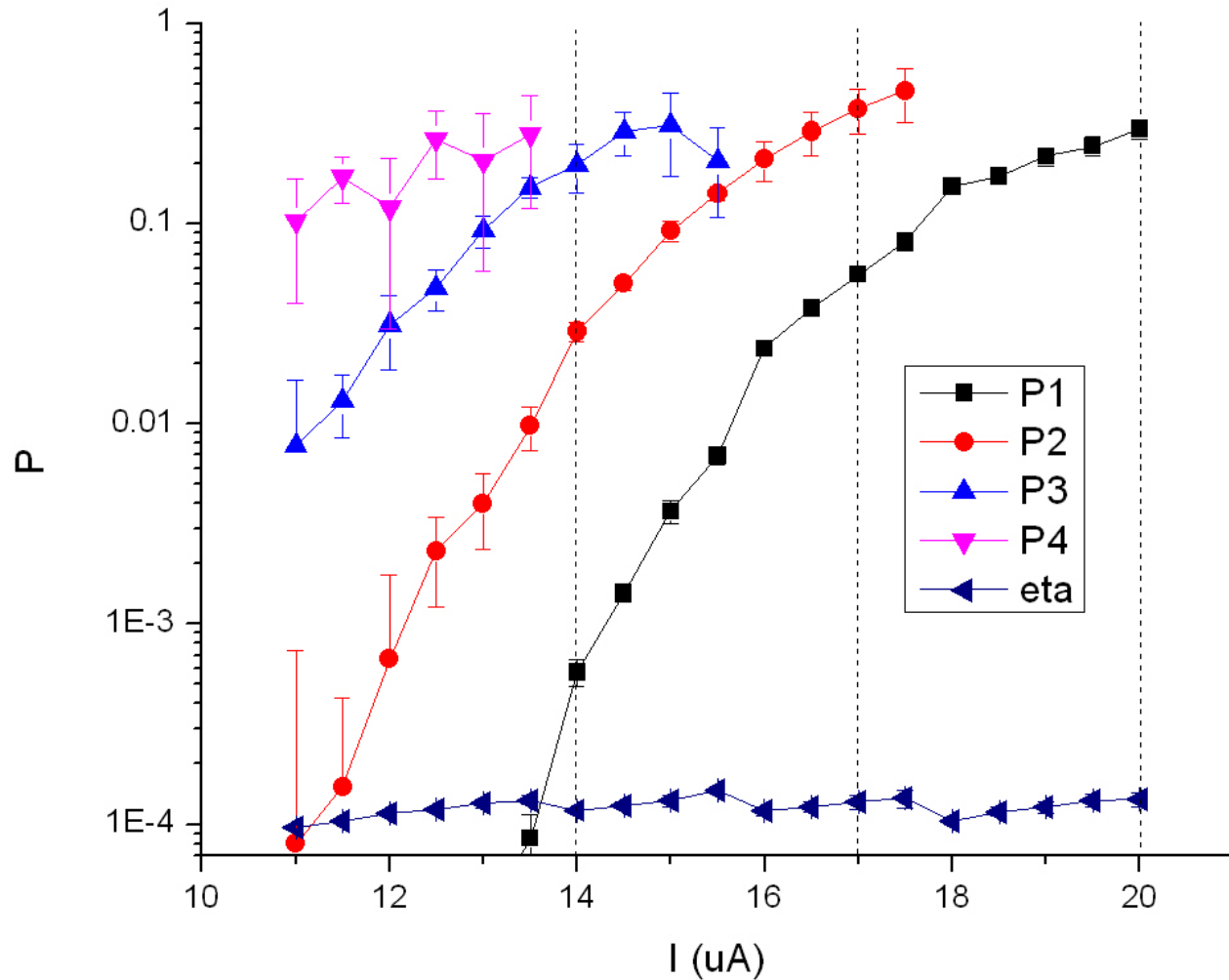
$R = (\eta N)^i$   
restricted to  
 $\eta N \ll 1$ ,  
lowest  $i$



# Now repeat this many times

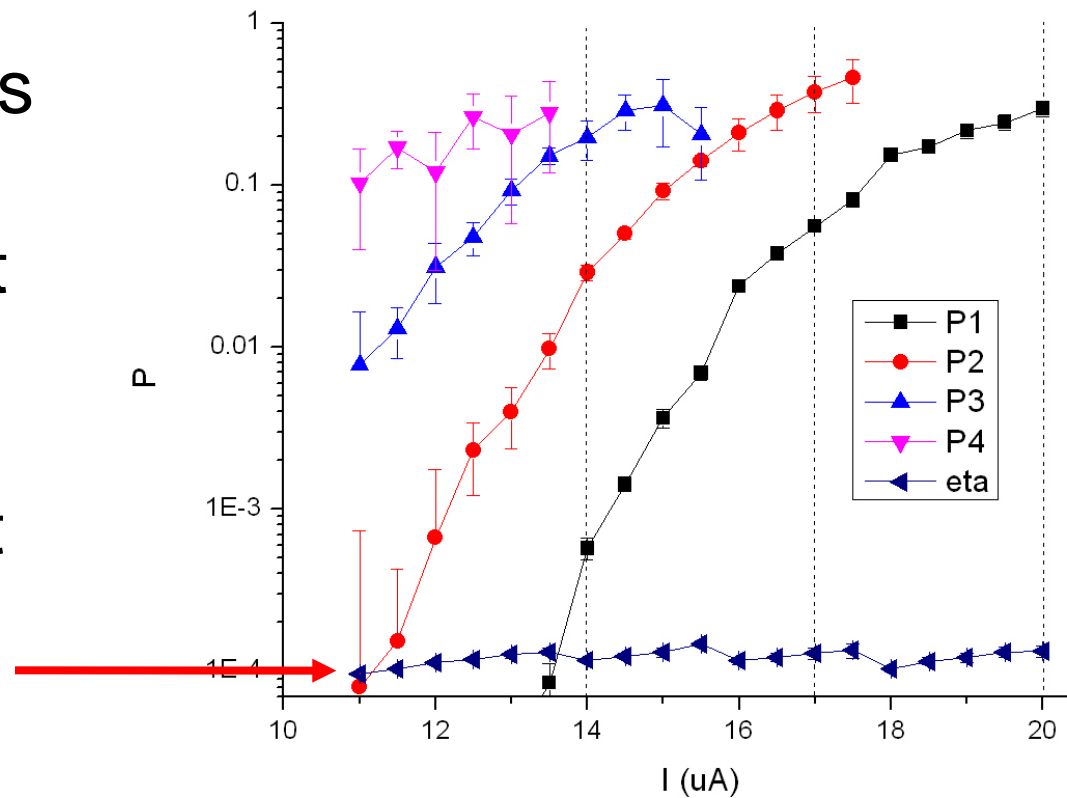
- For each current, vary the input power
- From the power dependence, reconstruct which photon processes are present

# Result from tomography



# Result from tomography

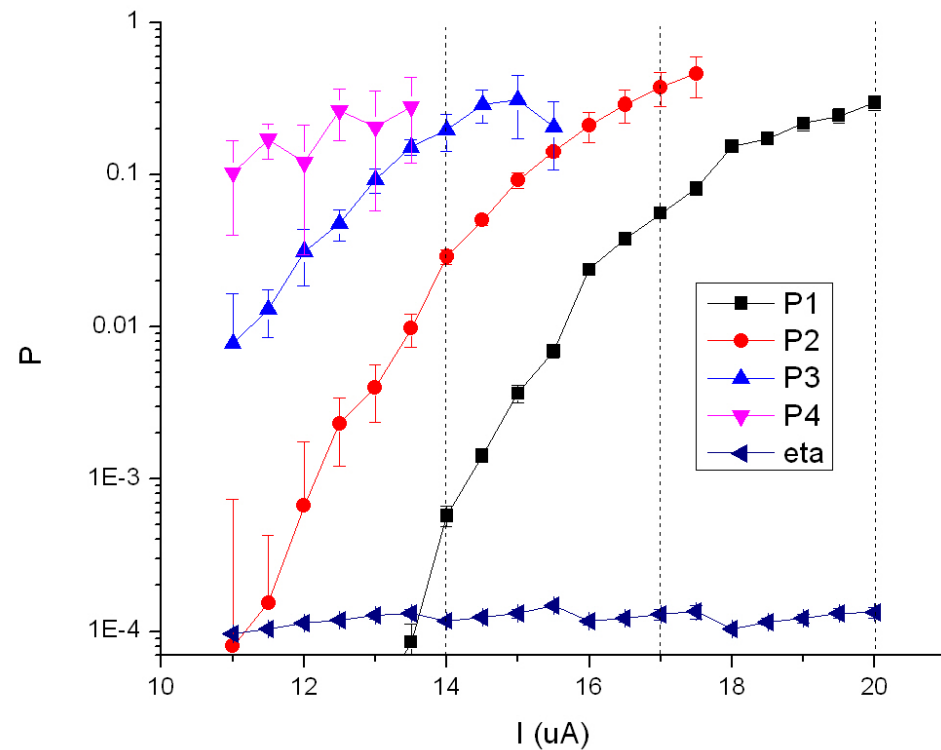
- We find: linear efficiency is independent of bias current
- This is a result, not an assumption (agnostic)
- Number consistent with overlap x absorption



Renema *et al*, Optics Express **20**, 2806-2813 (2012)

# Result from tomography

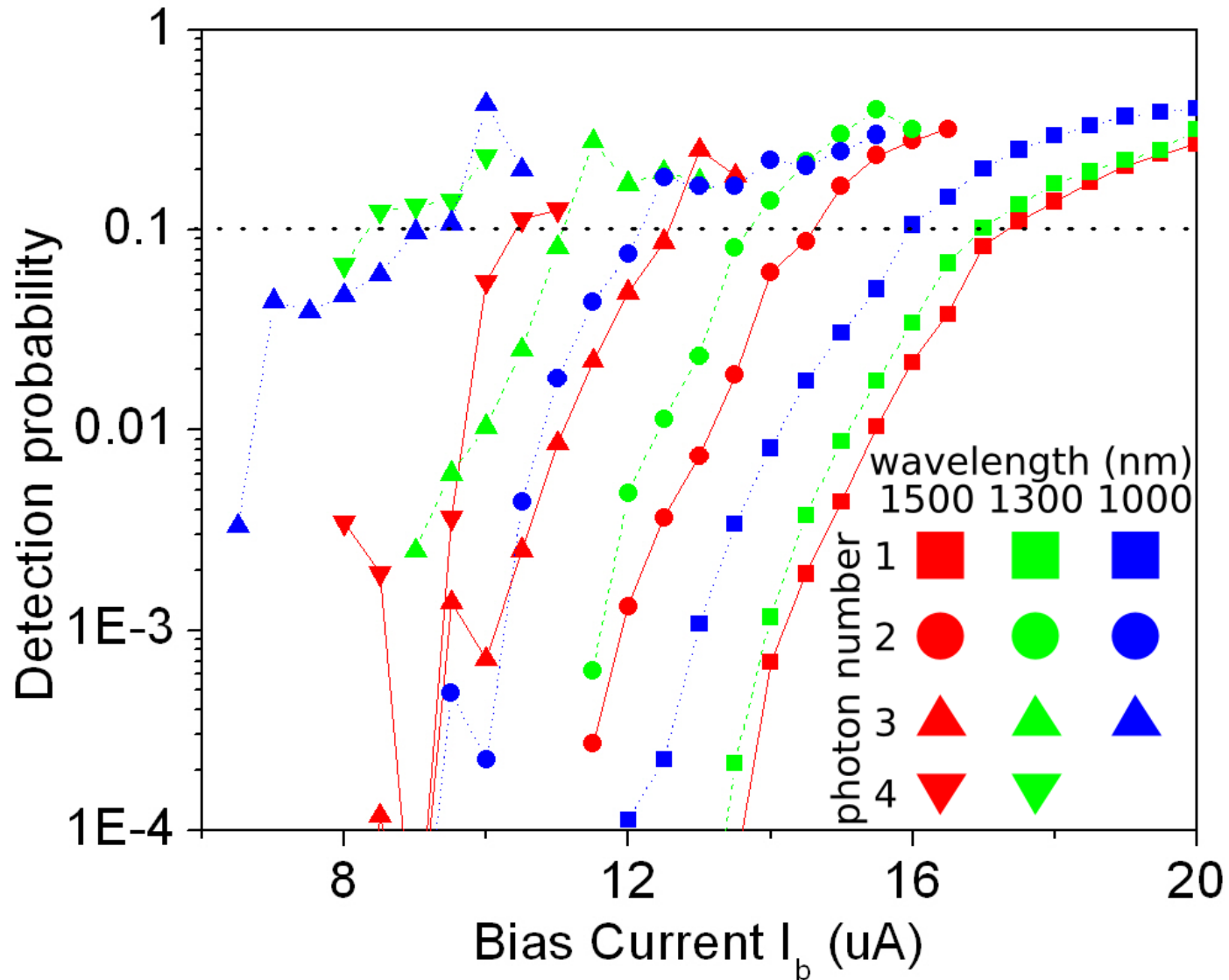
- $P_i$  internal response of the detector
- Independent of absorption, independent of incoupling
- There is more than linear efficiency



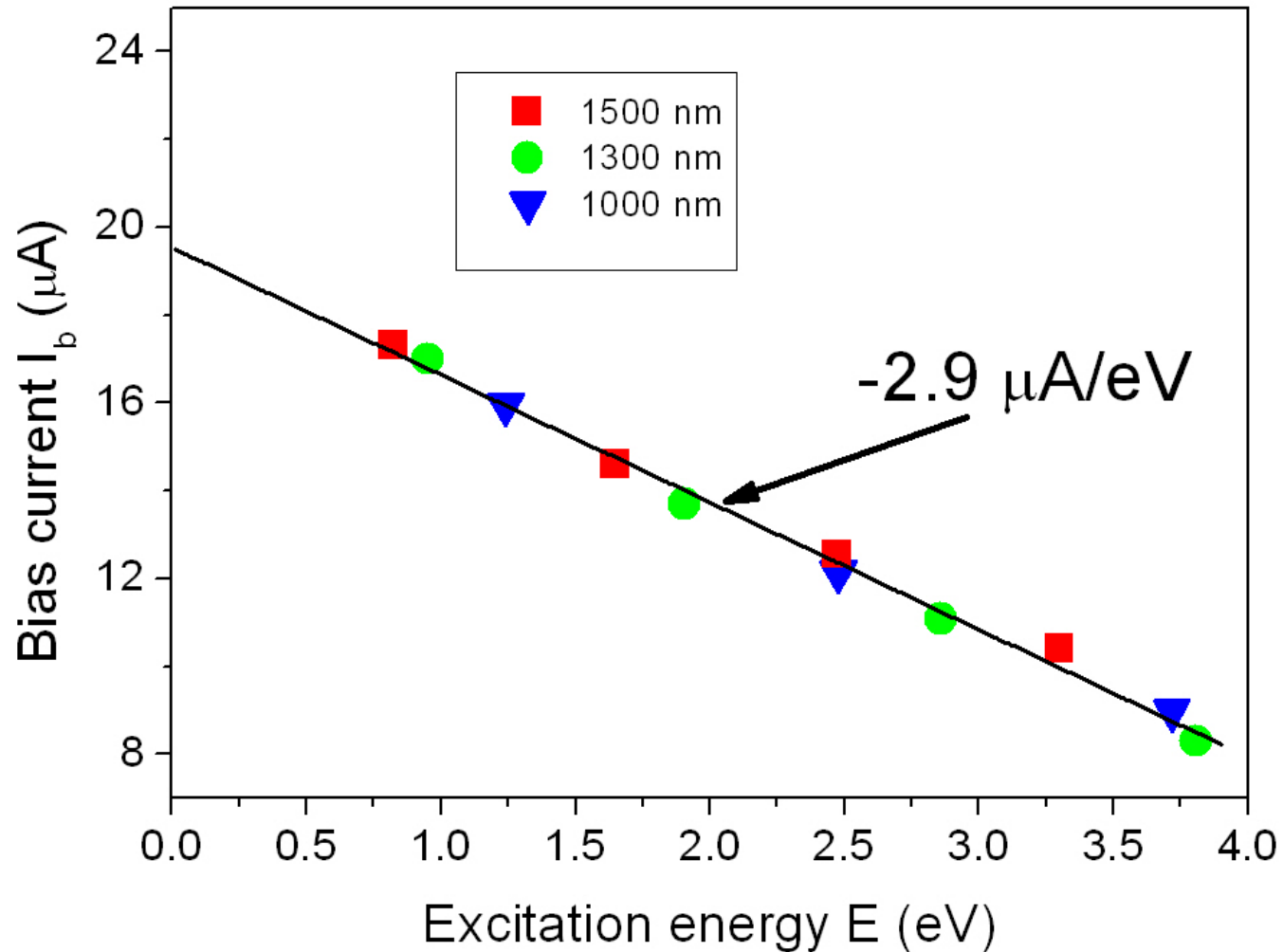
Renema *et al*, Optics Express **20**, 2806-2813 (2012)



# Multiple wavelengths



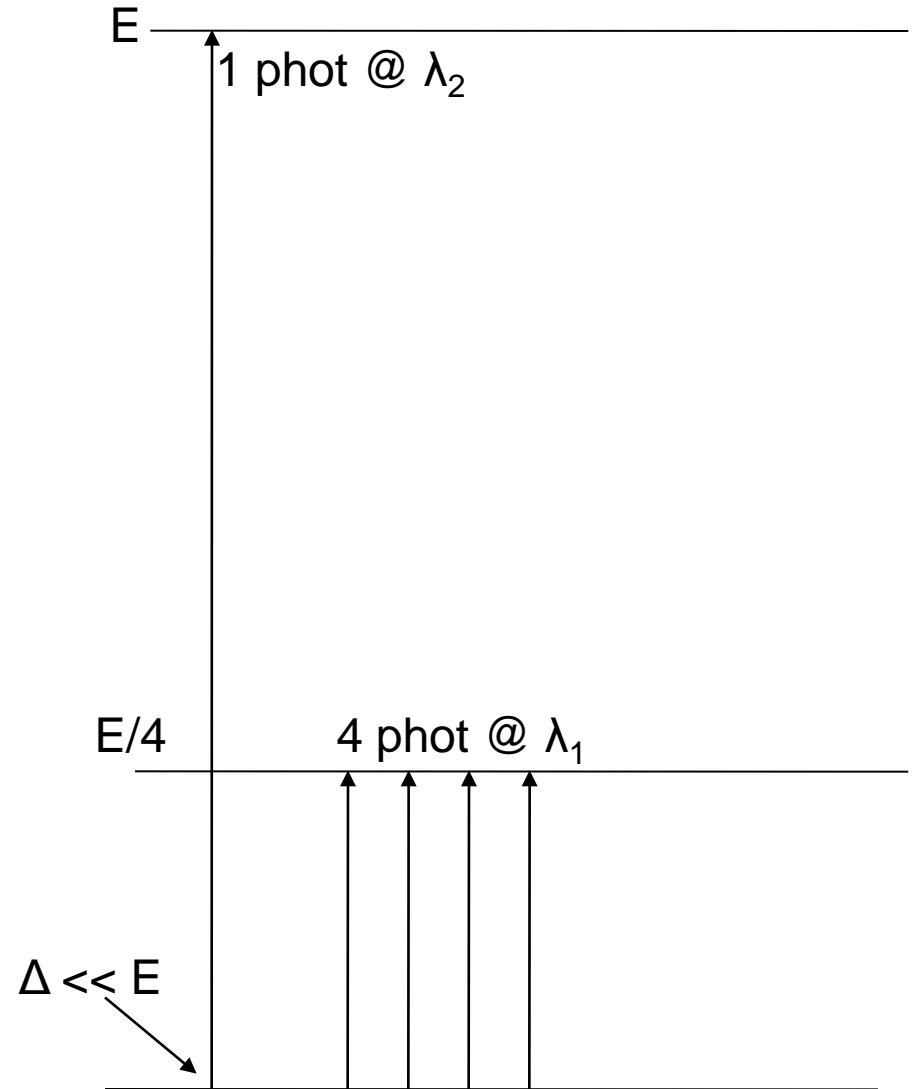
# Interchange energy/current



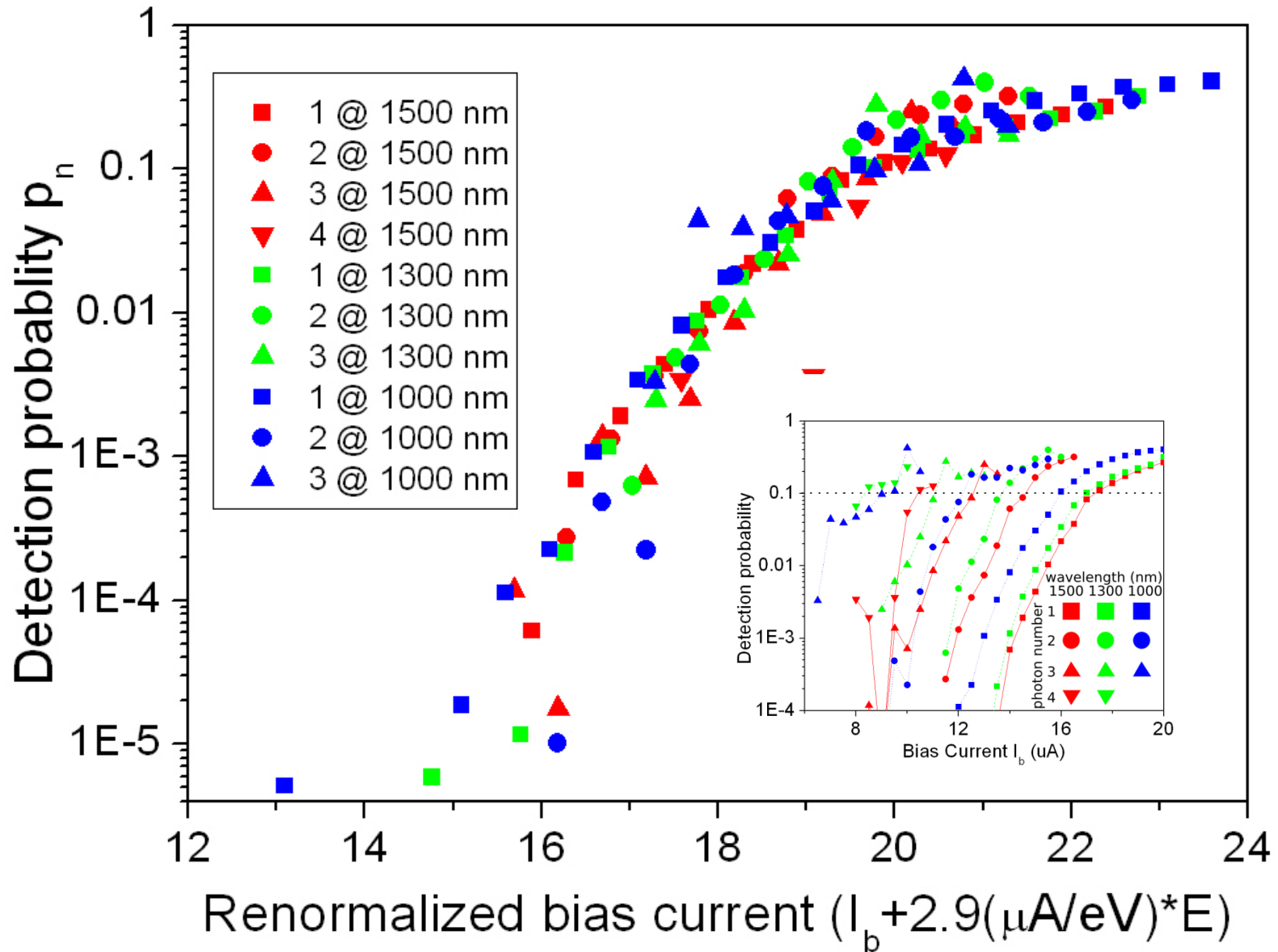
# QP conversion is linear

- No dependence on initial number of photons, only energy
- Detector is an energy detector

Renema *et al*, Phys Rev B **87**, 174526  
(2013)

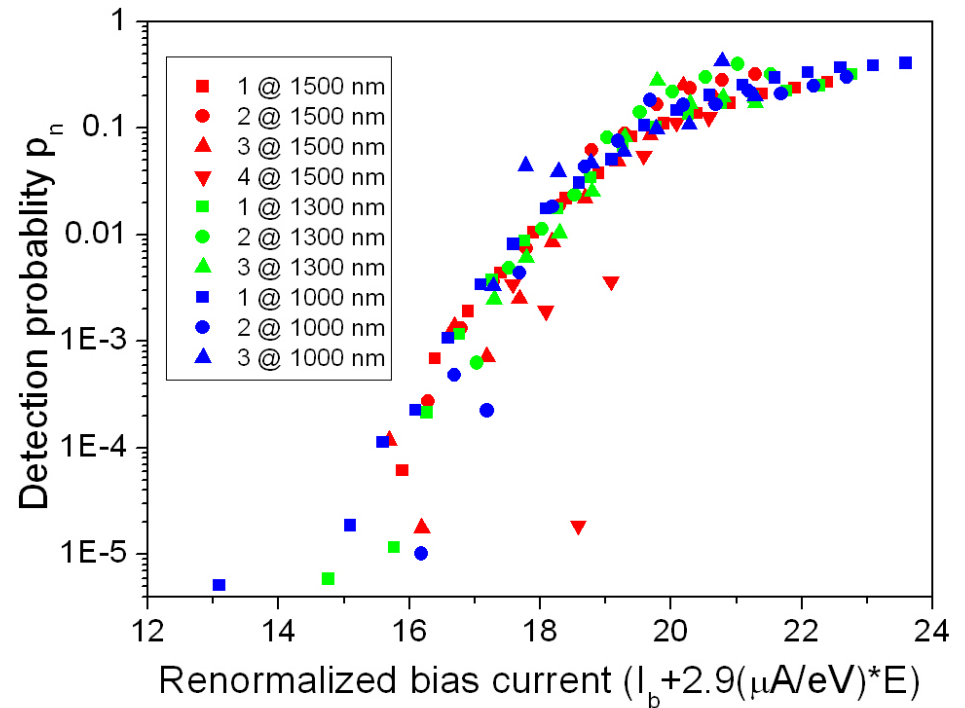


# Universal curve



# Universal curve

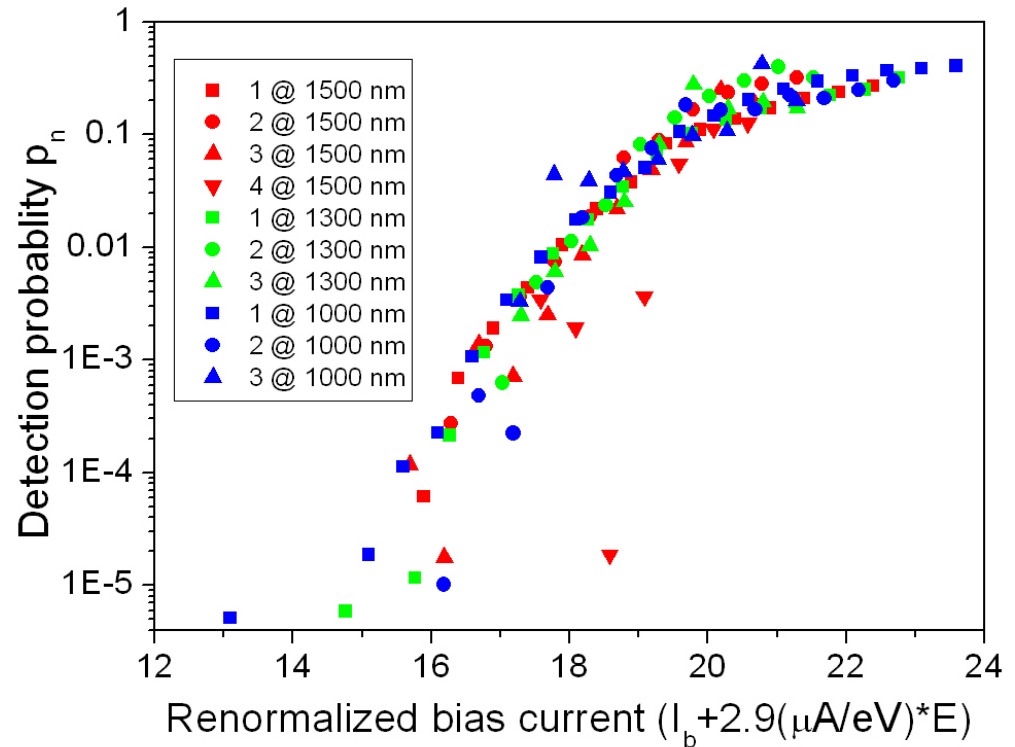
- $R(I, \lambda, N) = R(I + \gamma E)$  with  $E = N \cdot hc/\lambda$
- Goes beyond measuring edge of the plateau region



Renema *et al*, Phys Rev B **87**, 174526  
(2013)

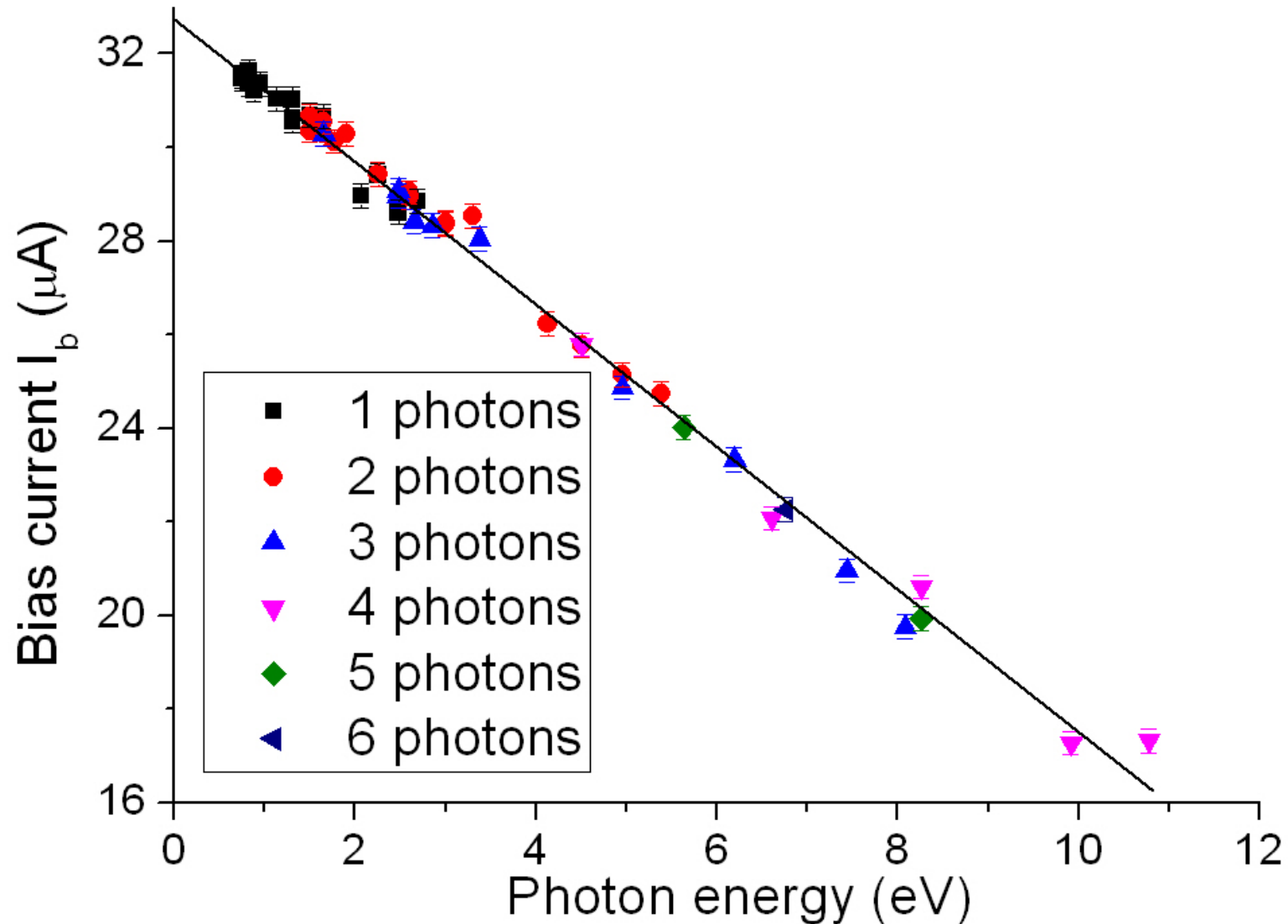
# Universal curve

- Fluctuation-assisted scales in the same way as plateau response
- Results compatible with theory (both Engel & Vodolazov)



Renema *et al*, Phys Rev B **87**, 174526  
(2013)

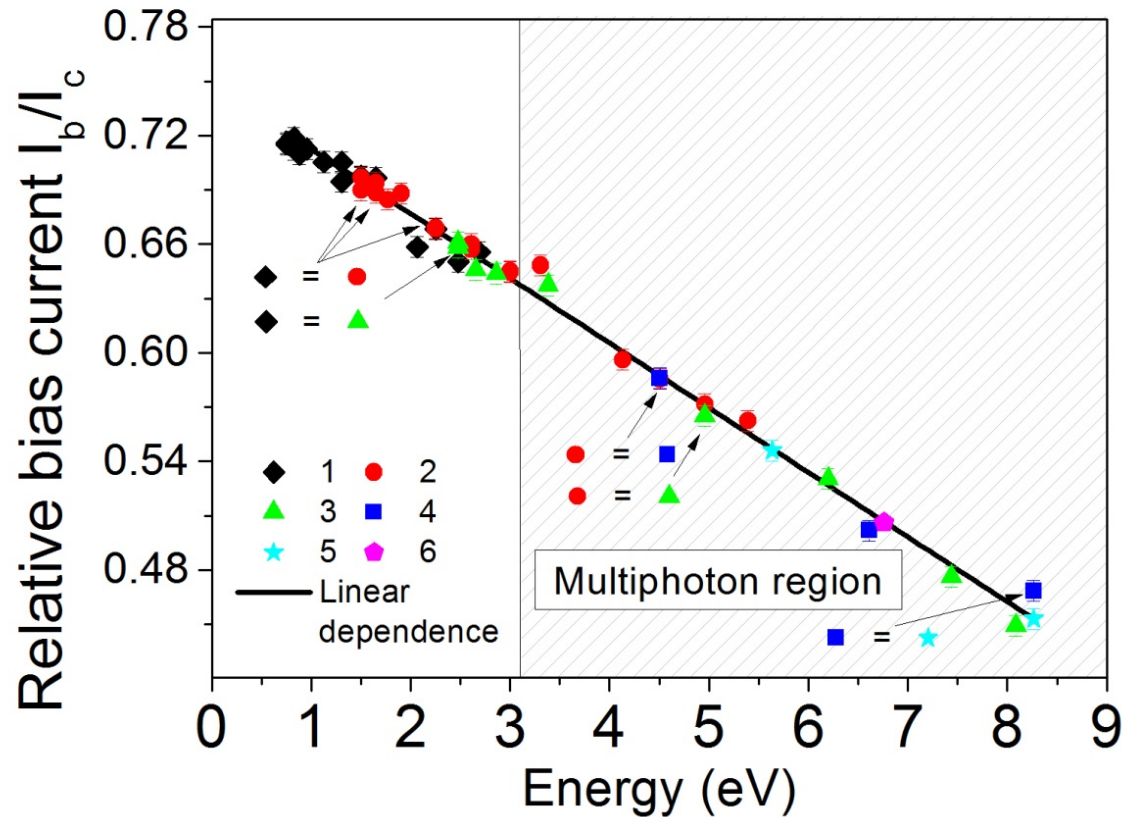
# Result on 220 nm detector



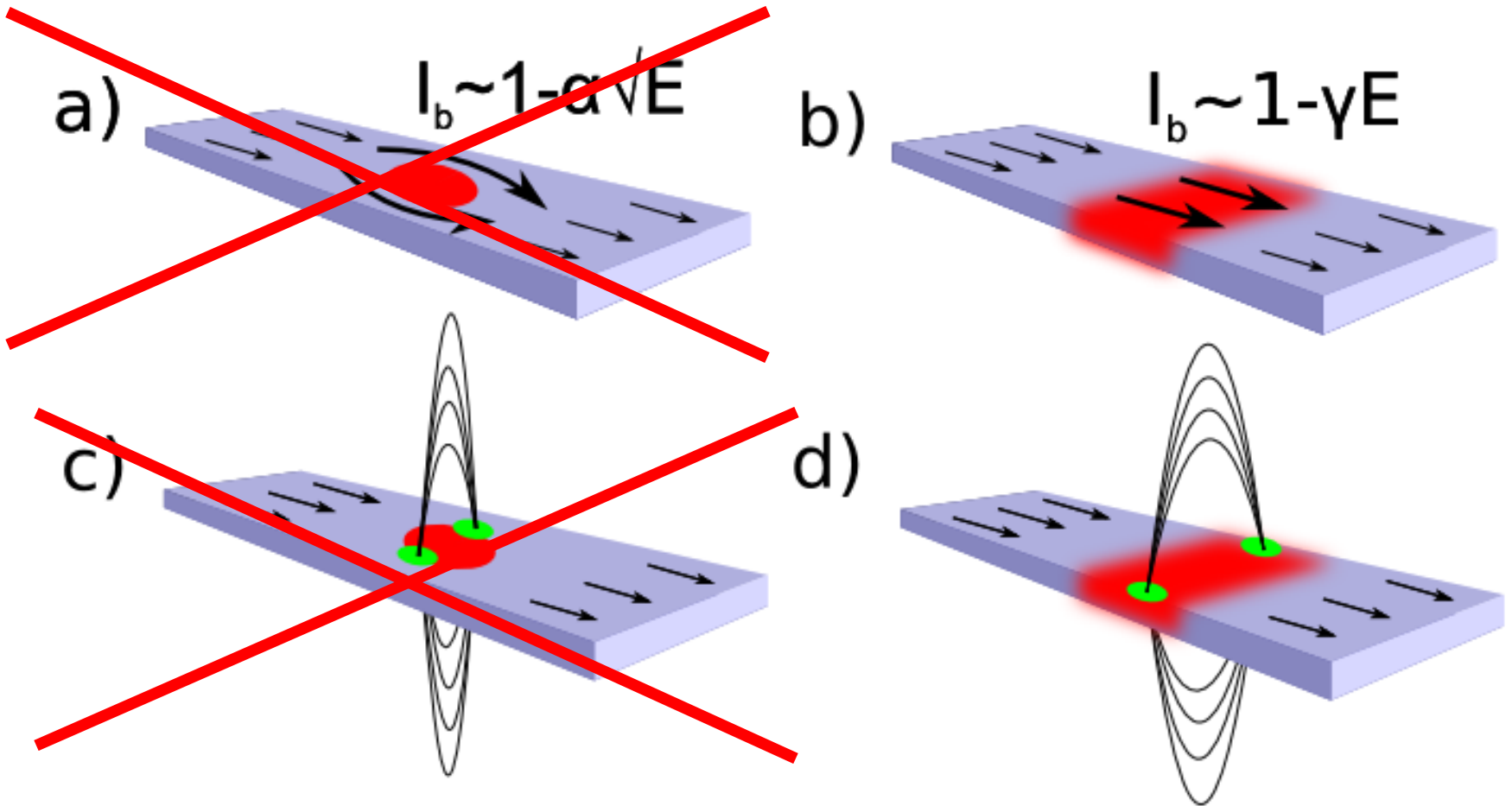


# Extreme dynamic range

- Find  $I_b = I_0 - \gamma E$
- 10.8 eV (X-UV):  
 $\lambda_{\text{eff}} = 115 \text{ nm}$
- Photon regimes overlap -> no stitching errors

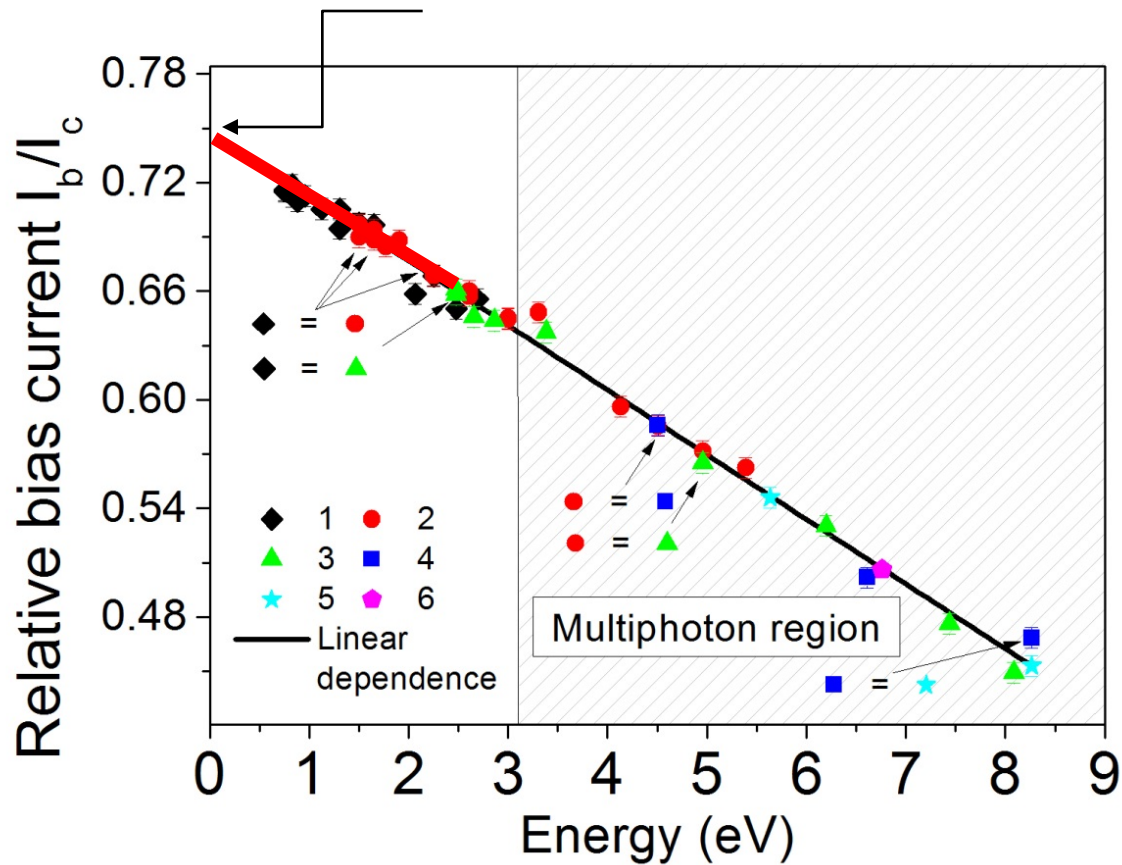


# First conclusion

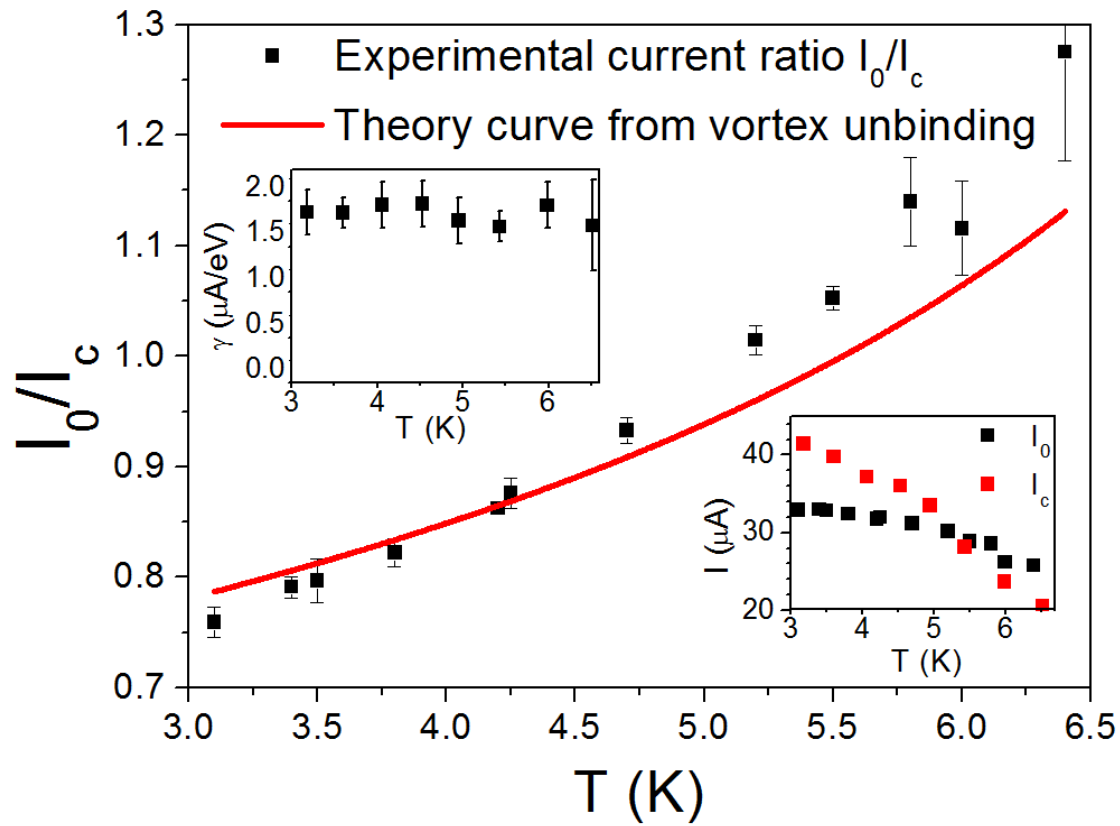


# Moreover

$$I_0 \neq I_c !!!$$

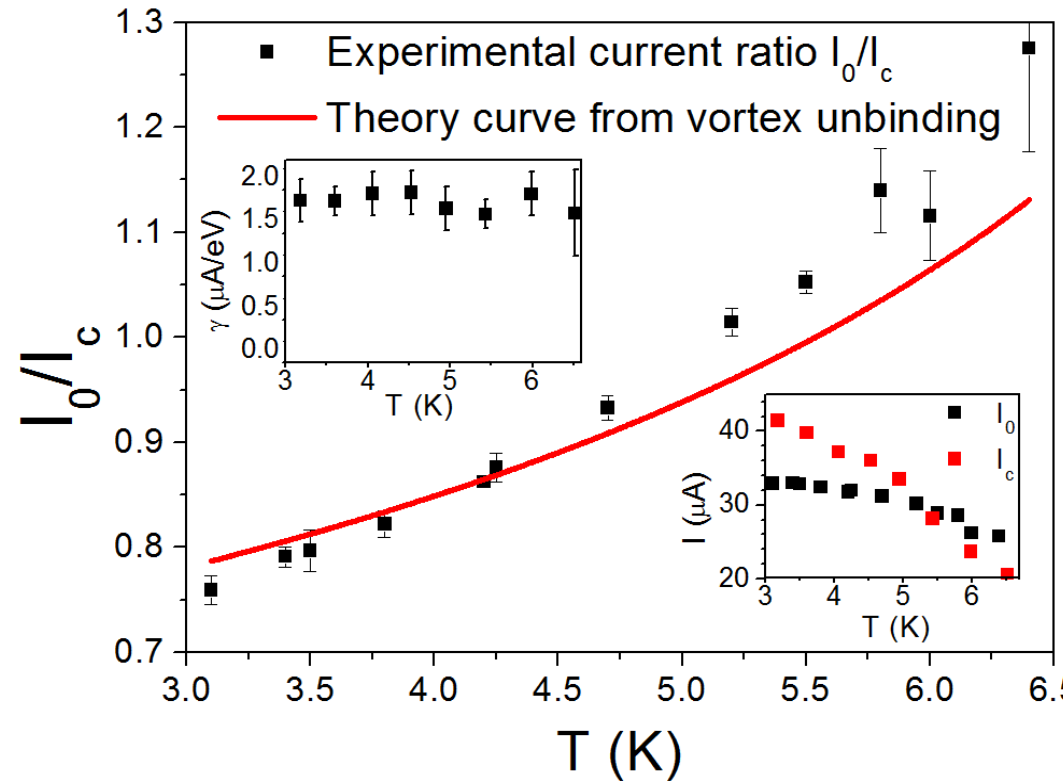


# Temperature dependence of $I_0$

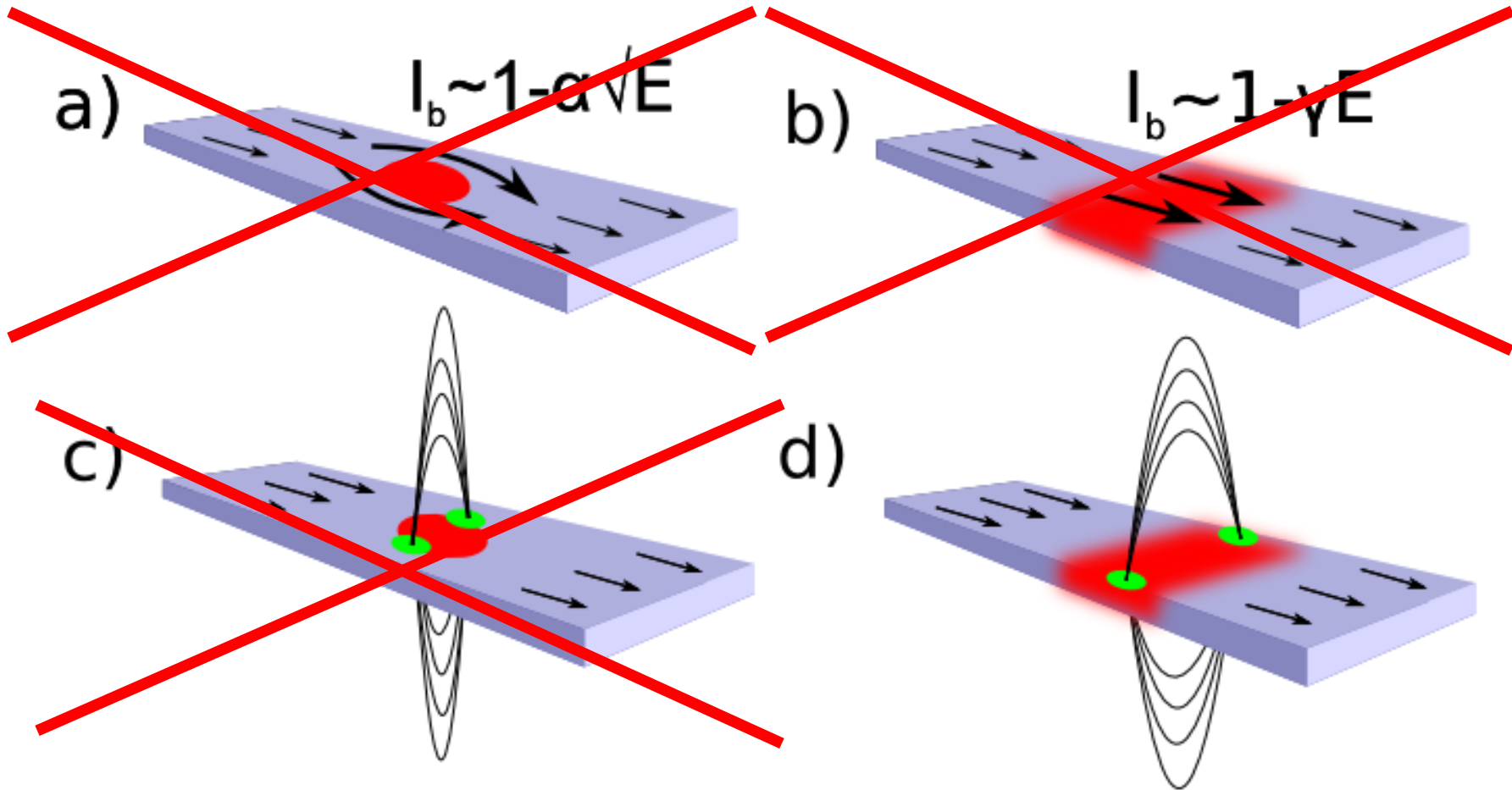


# Temperature dependence of $I_0$

- $I_b = I_0 - \gamma E$
- Only  $I_0$  temperature dependent
- Find cases  $I_0 > I_c$  and  $I_0 < I_c$
- Ratio  $I_0 / I_c$  follows vortex entry energy prediction



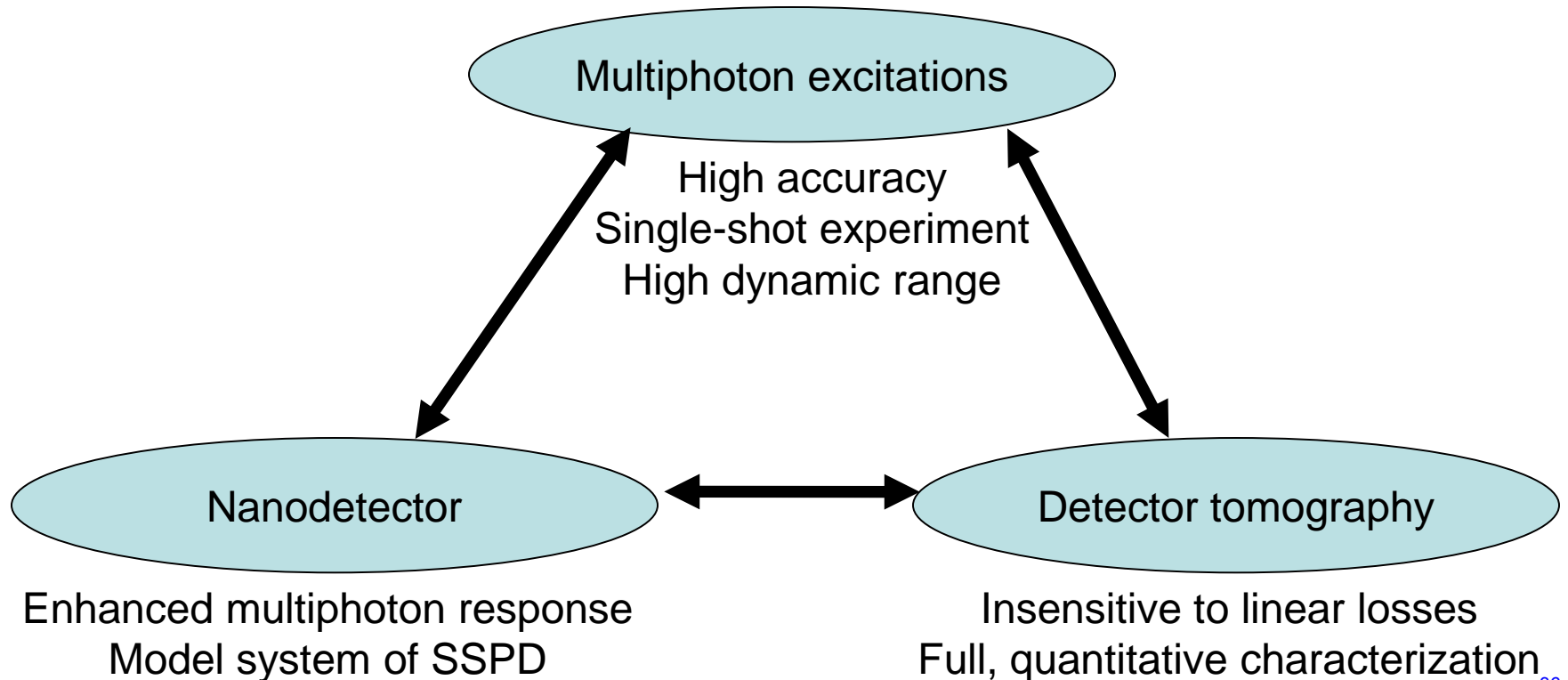
# Second conclusion



Renema *et al*, Phys Rev Letters **112**, 117604 (2014)

# Conclusions

- There is more in the detector than linear efficiency
- Quantum tomography useful for inner workings of detector
  - Linear energy-current relation up to X-UV
- Temperature dependence fixed by vortex behaviour





# Acknowledgements

Supervisor: M.P. van Exter

Co-workers: Q. Wang,  
M.J.A. de Dood

Sample Fabrication TUE / IFN: R. Gaudio, A. Fiore, G. Frucci, Z. Zhou, A. Gaggero, F. Mattioli, R. Leoni

Mathematics: R. Gill

Numerical simulations: A. Engel  
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Renema *et al*, Optics Express **20**, 2806-2813 (2012)  
JJ Renema *et al*, Phys Rev A **86**, 062113 (2012)  
Renema *et al*, Phys Rev B **87**, 174526 (2013)  
Renema *et al*, Phys Rev Letters **112**, 117604 (2014)

