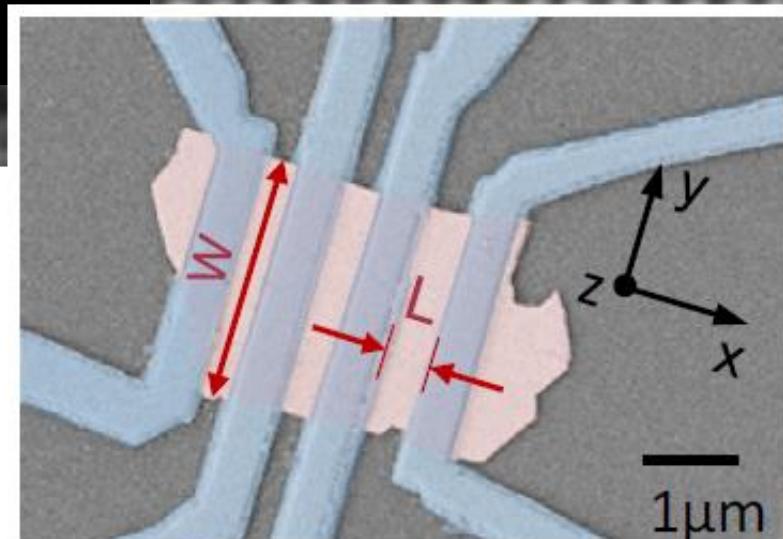
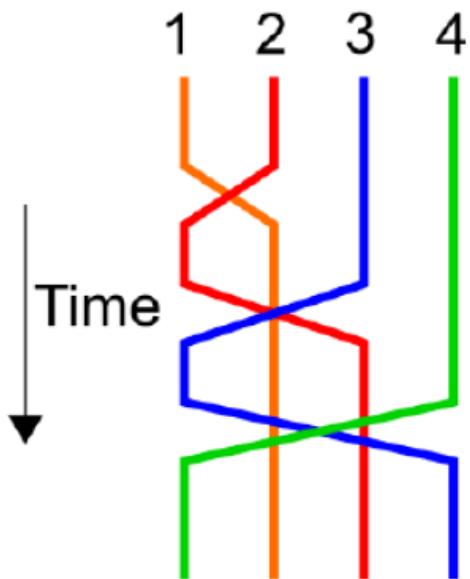


# Bismuth based topological Josephson junctions

Alexander Brinkman

University of Twente, the Netherlands



**MESA+**  
INSTITUTE FOR NANOTECHNOLOGY

# Acknowledgements

## University of Twente



Chuan  
Li



Jorrit  
de Boer



Bob  
de Ronde



Martijn  
Lankhorst

## University of Amsterdam

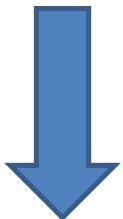
Yingkai Huang  
Anne de Visser  
Erik van Heumen  
Mark Golden

## Forschungszentrum Jülich

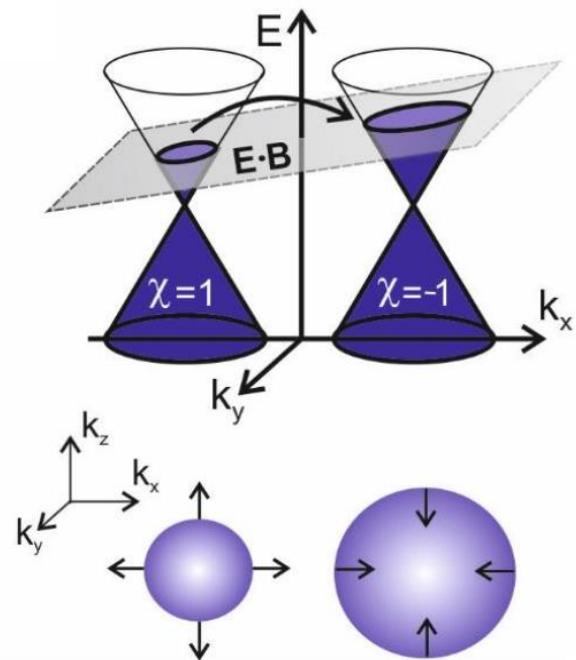
Peter Schüffelgen  
Daniel Rosenbach  
Thomas Schäpers  
Detlef Grüzmacher

# Message of this talk

3D helical **spin-momentum locking**  
of Dirac semimetal



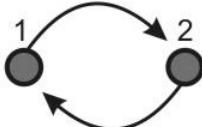
- Topological transport
- $4\pi$ -periodic Andreev bound states
- Zeeman induced  $\pi$ -junction



# Non-Abelian anyons

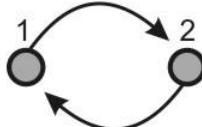
(a)

bosons



(b)

fermions



(c)

anyons

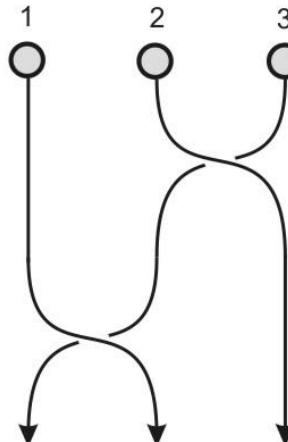


$$\Psi_{12} = \Psi_{21}$$

$$\Psi_{12} = -\Psi_{21}$$

$$\Psi_{12} = e^{i\phi} \Psi_{21}$$

non-abelian anyons

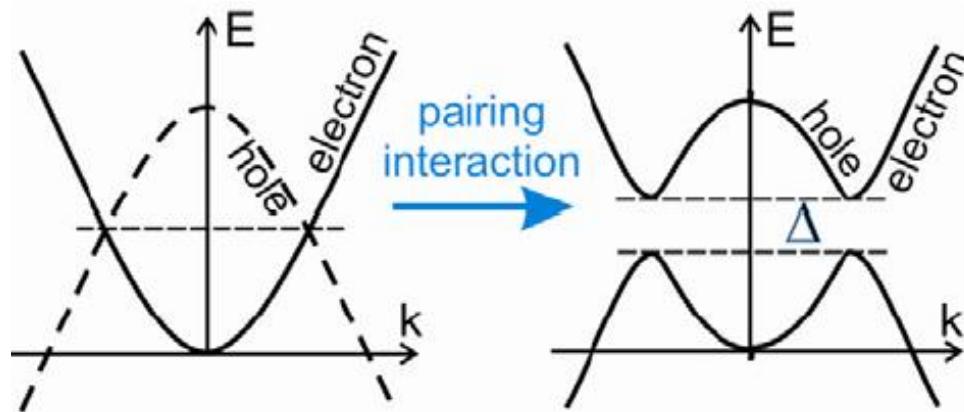


$$U_1 U_2 \neq U_2 U_1$$

*“Qubits stored in topology are less likely to be affected by their environment.”*

**Microsoft**

# Majoranas in superconductors?

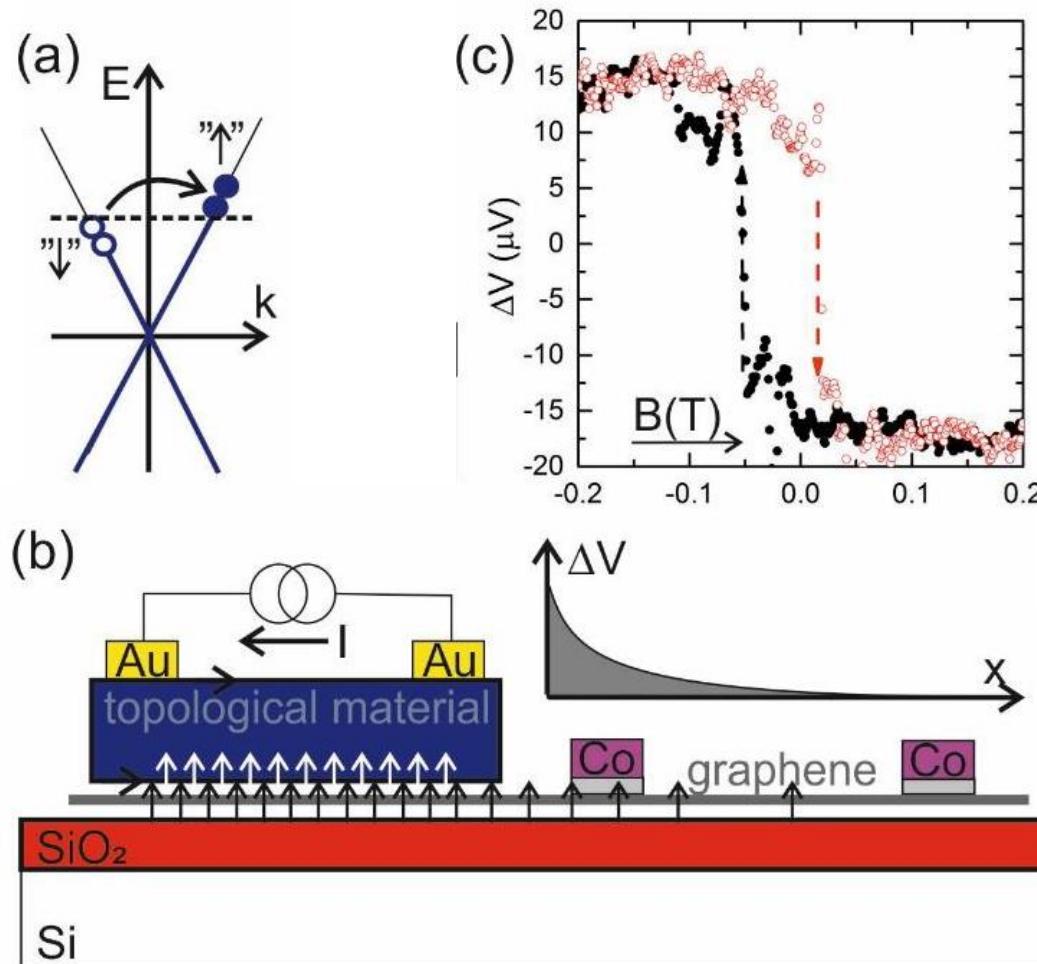


Bogoliubov quasiparticles:  $\gamma_\epsilon = \gamma_{-\epsilon}^\dagger$

Needed for Majorana:  $\gamma_\epsilon = \gamma_\epsilon^\dagger$

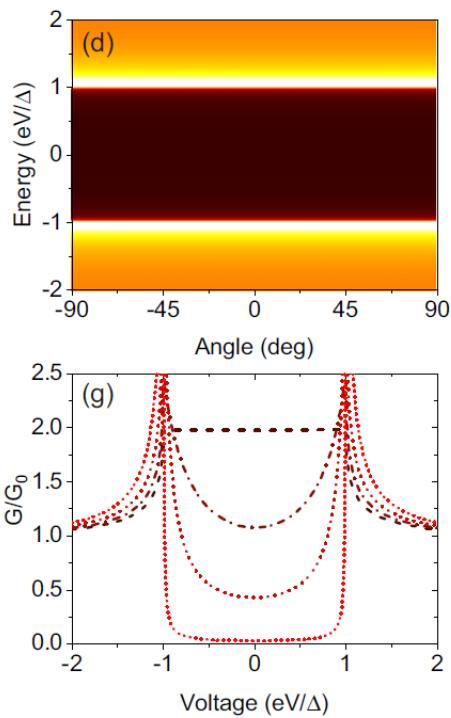
→ Zero energy state required

# Spin-momentum locking

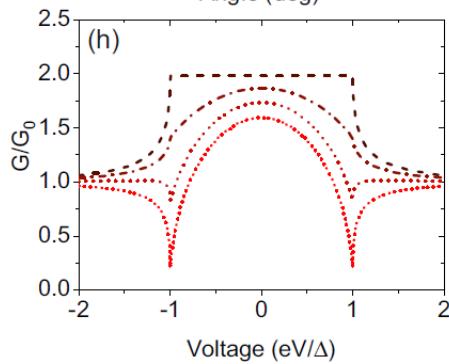
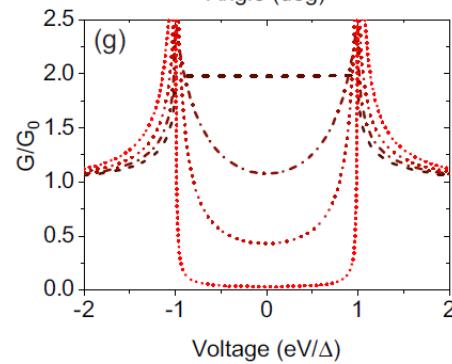
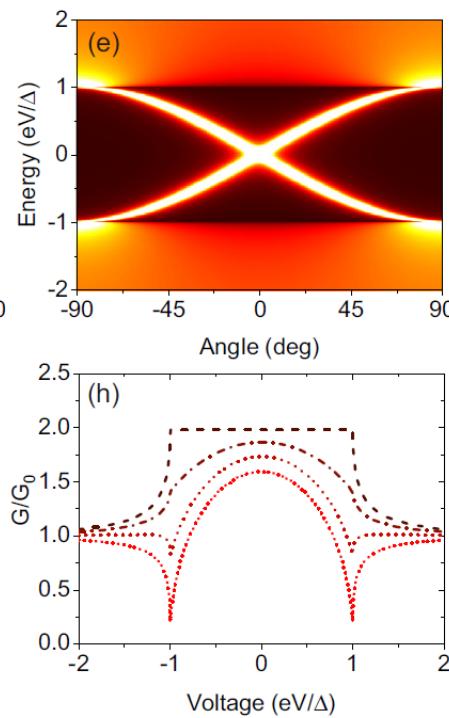


# Topology and superconductivity

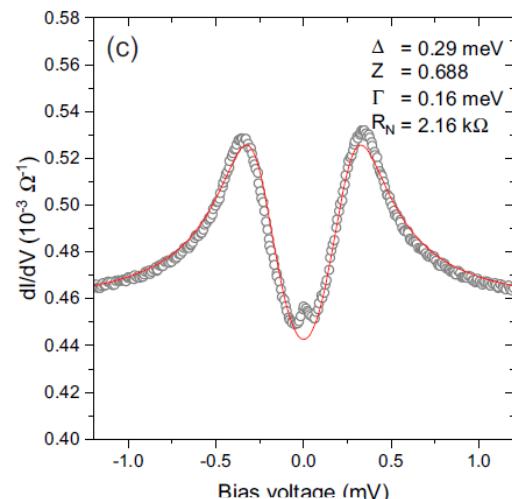
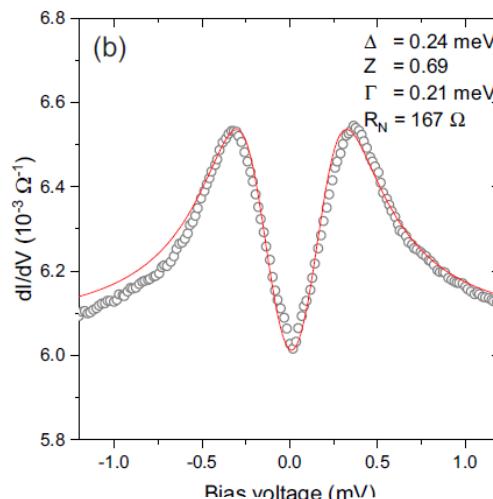
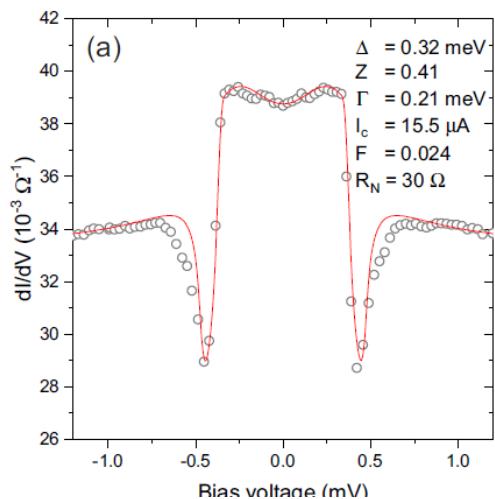
3D s-wave:



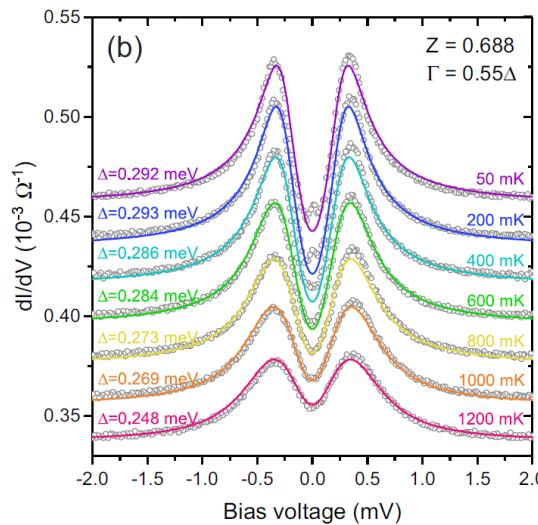
3D chiral  $p$ :



# PdTe<sub>2</sub> type II Dirac semimetal

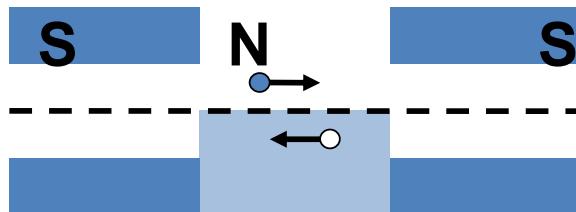


Dominant s-wave superconductor...

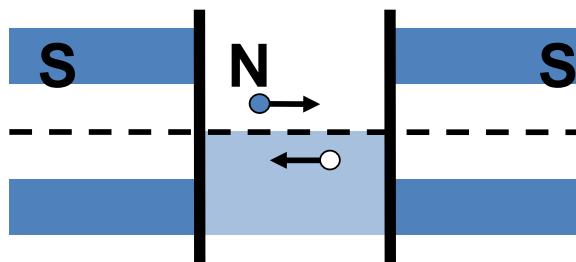


# Why topological Josephson junctions?

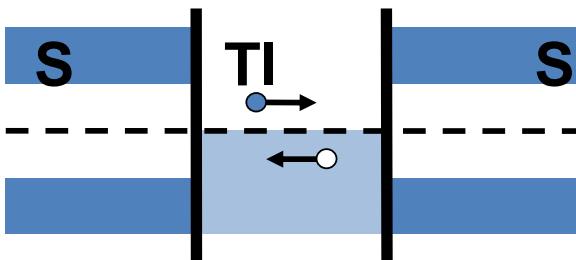
Standard SNS:



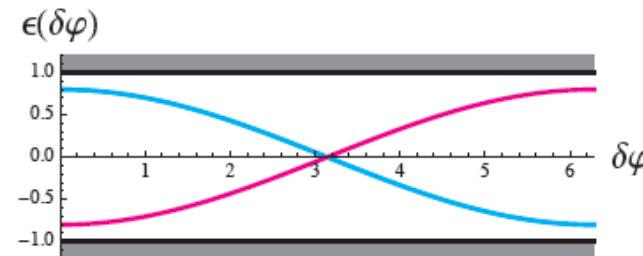
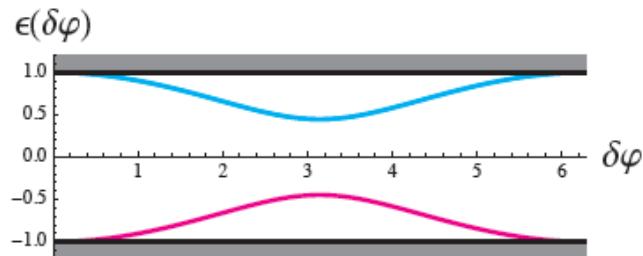
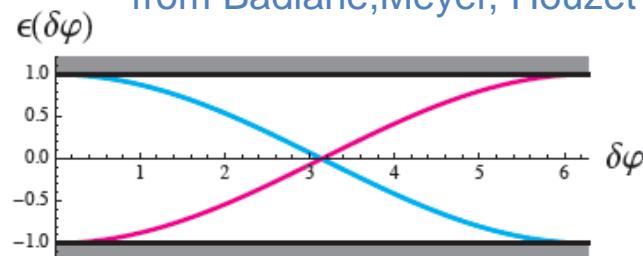
Including interface barrier:



With topological interlayer:

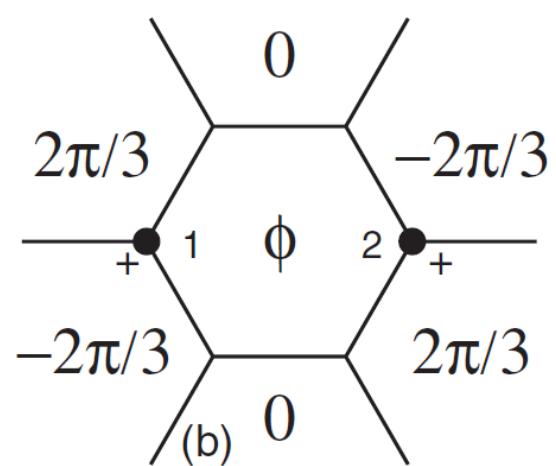
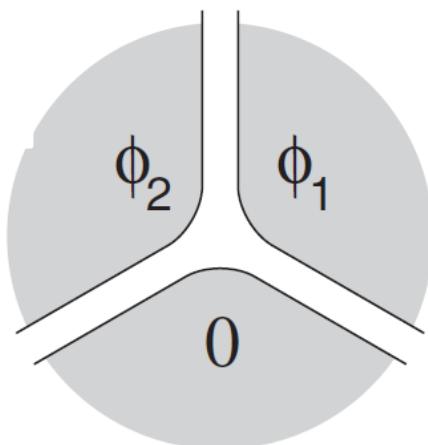
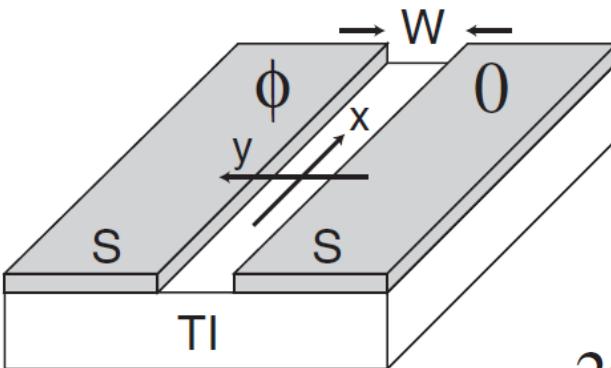
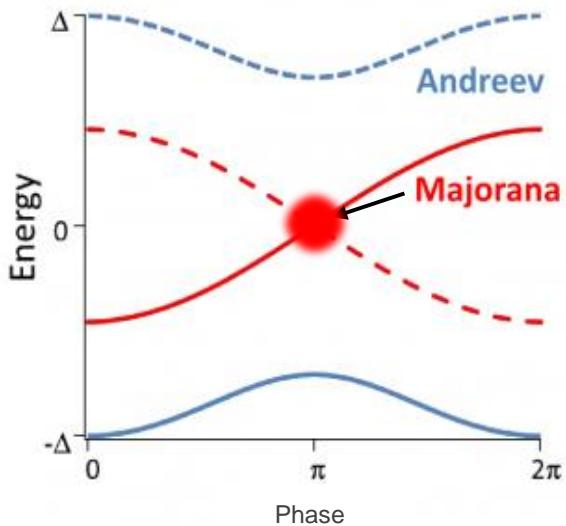


Calculations and figure  
from Badiane,Meyer, Houzet



# Majorana manipulation on a surface

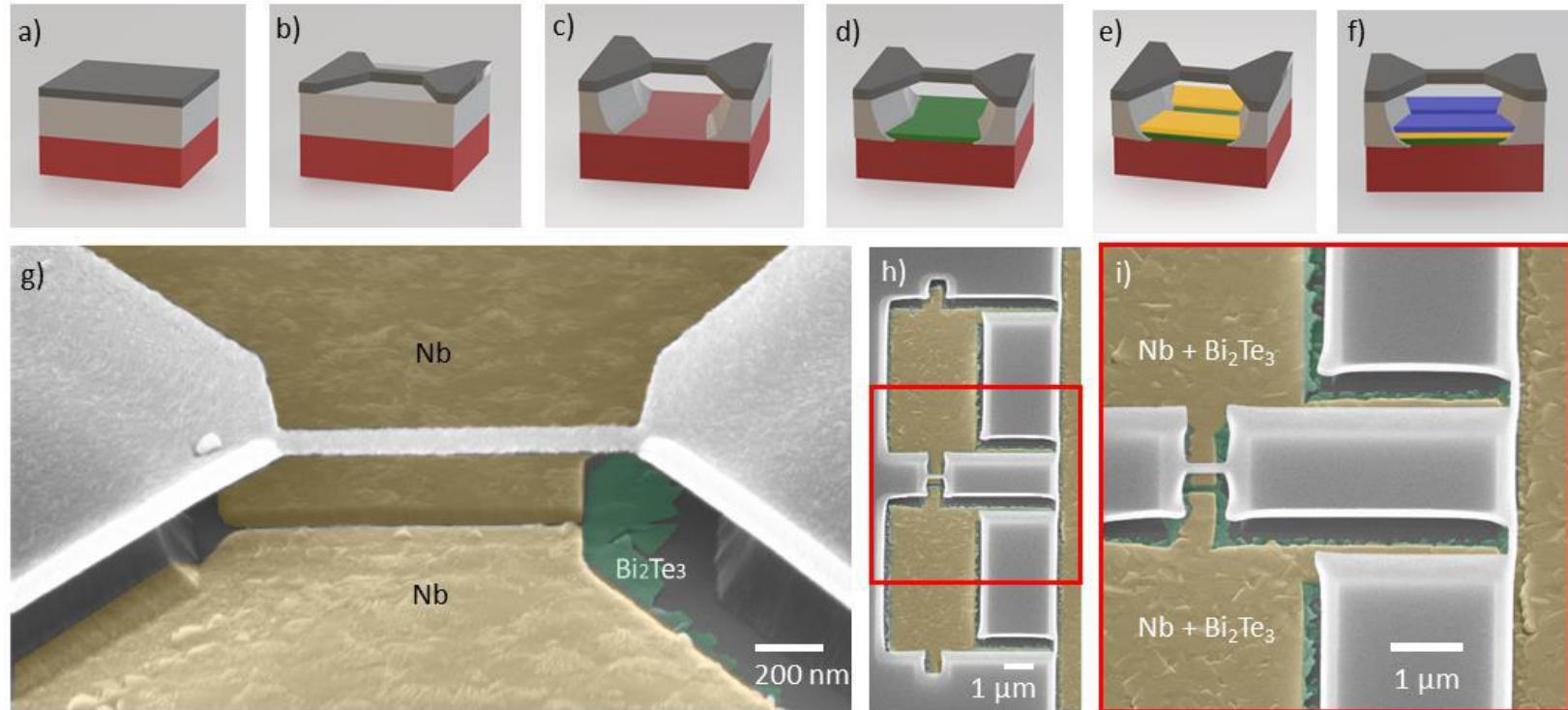
Fu & Kane, PRL **100**, 096407 (2008)



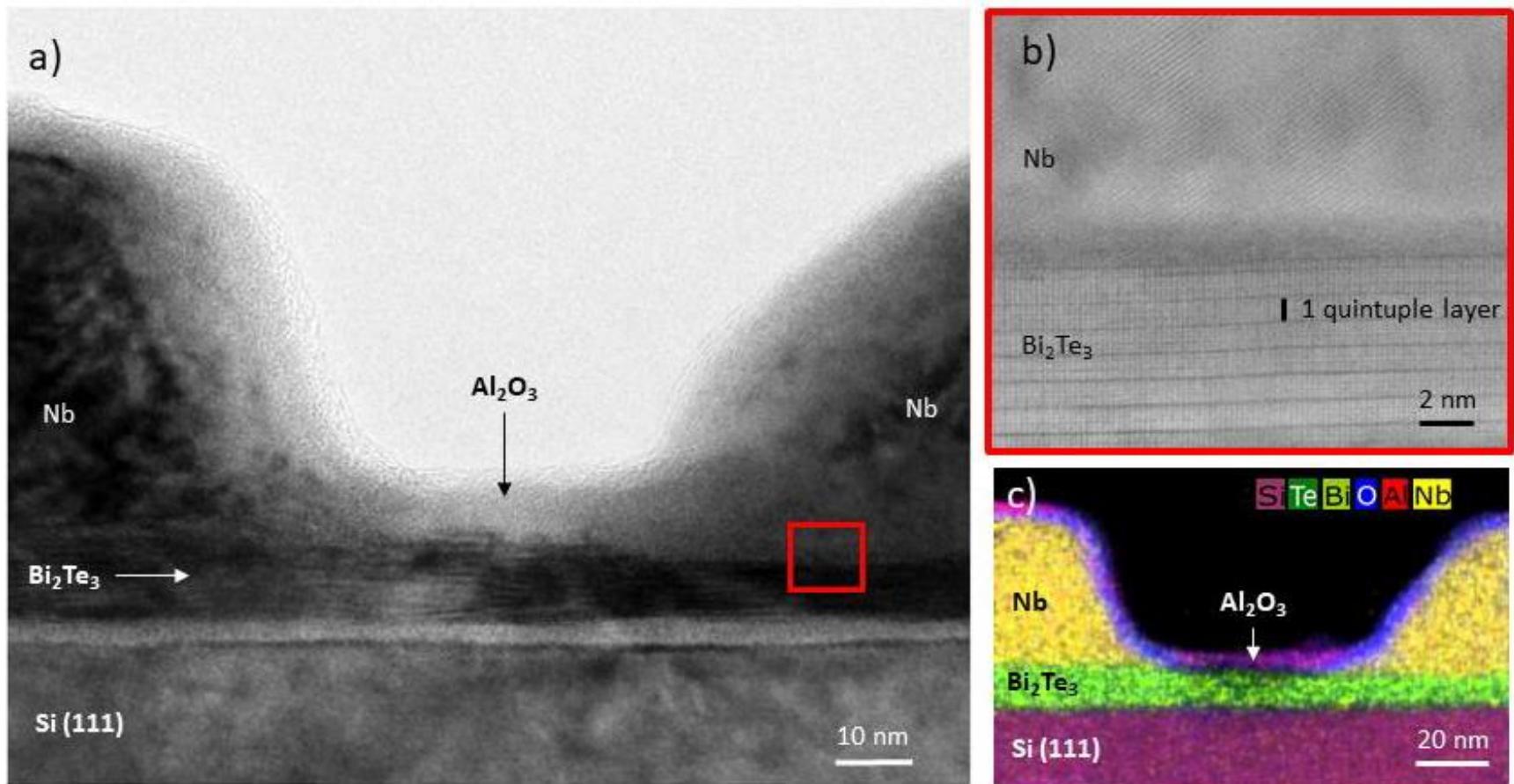
# Advanced device fabrication

In situ electrode deposition by shadow mask technology → enhanced I<sub>c</sub>R<sub>n</sub>

Selective area growth → towards 1D



# Advanced device fabrication

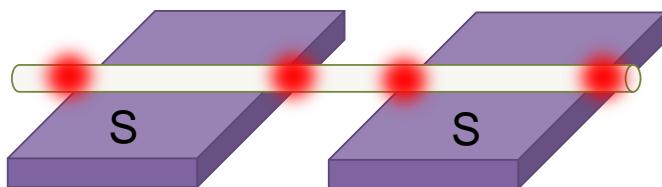


Nature Nanotechnology **14**, 821 (2019)

# Majoranas

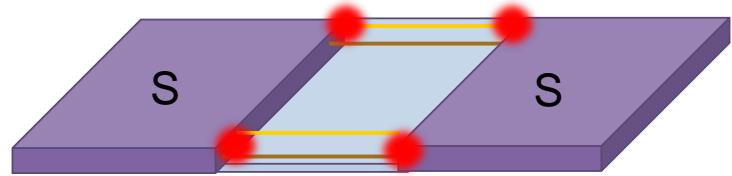
1D Nanowires (InAs, InSb...)

Strong evidence (Delft, Lund, Copenhagen)



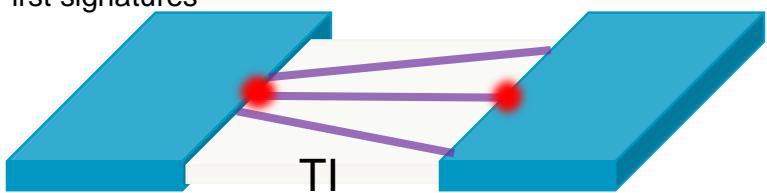
1D edges of 2D QSH (HgTe)

Strong evidence (Würzburg)



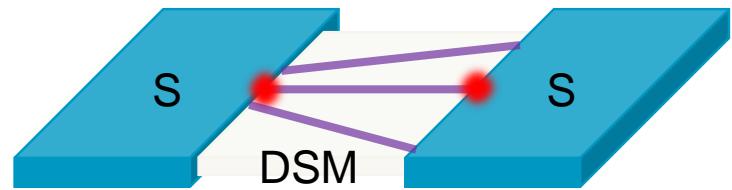
2D surface states of 3D topological insulators  
( $\text{Bi}_2\text{Se}_3$ ,  $\text{Bi}_2\text{Te}_3$ ,  $\text{Bi}_1\text{Sb}_1\text{Te}_1\text{Se}_2$ )

First signatures

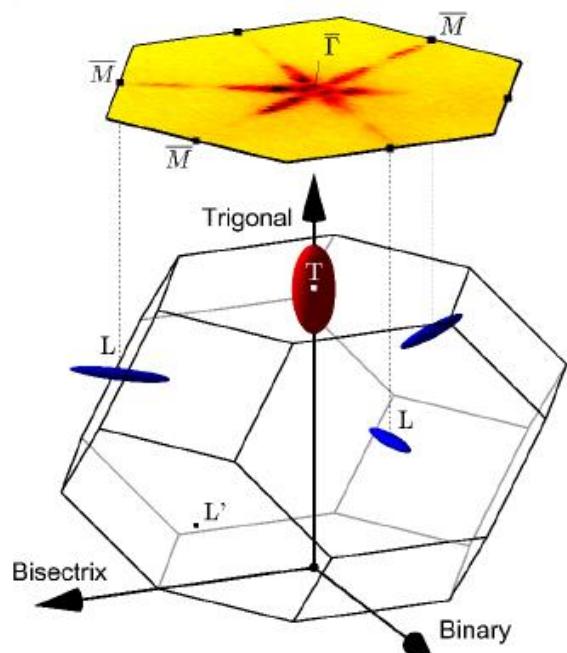


3D bulk states (Dirac/Weyl semimetals)

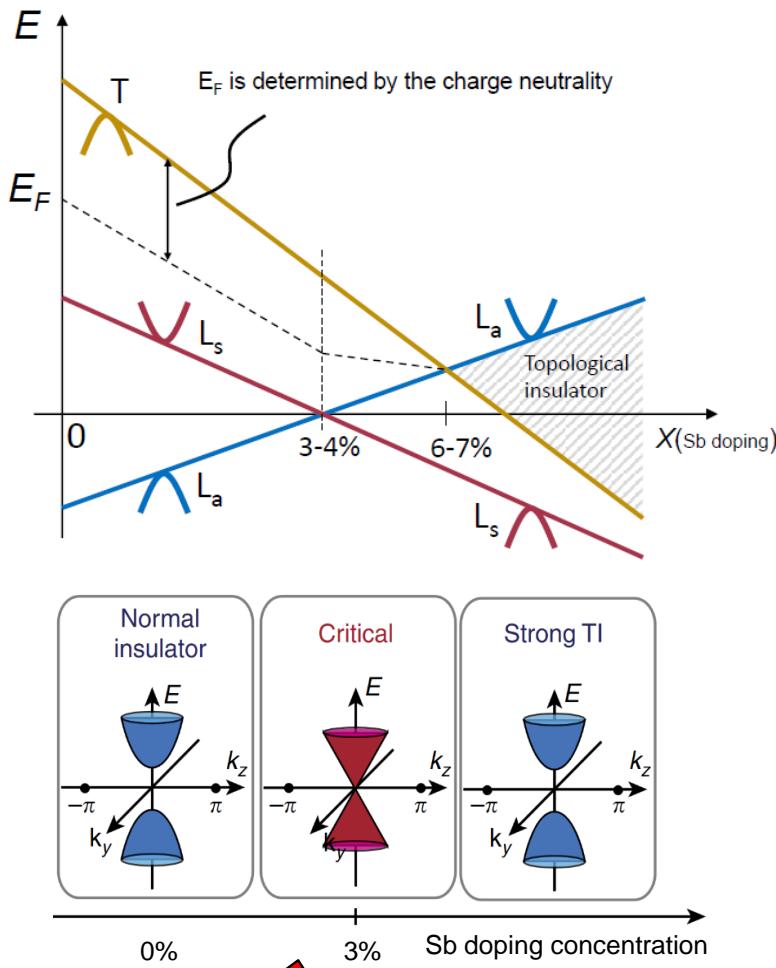
This talk



# 3D Dirac semimetal $\text{Bi}_{1-x}\text{Sb}_x$

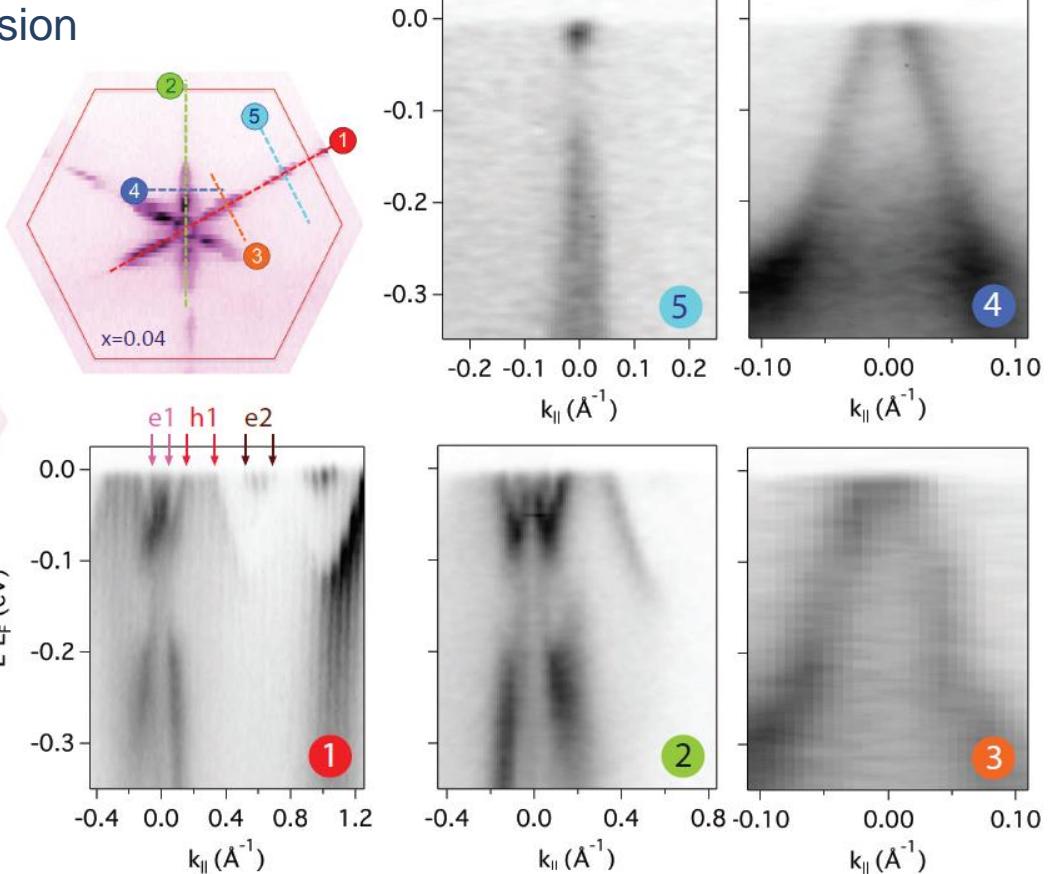
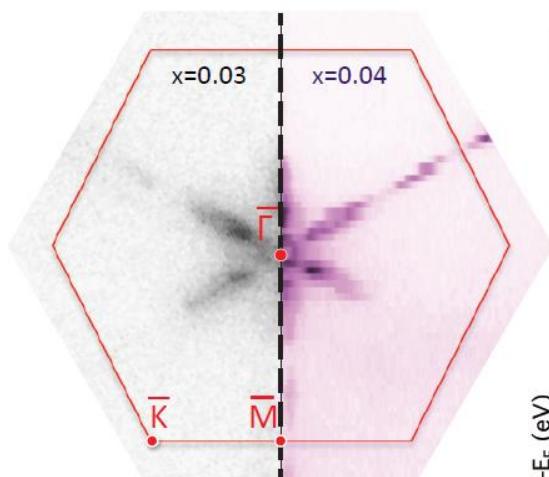


“3D graphene”



# Surface (111) of $\text{Bi}_{1-x}\text{Sb}_x$

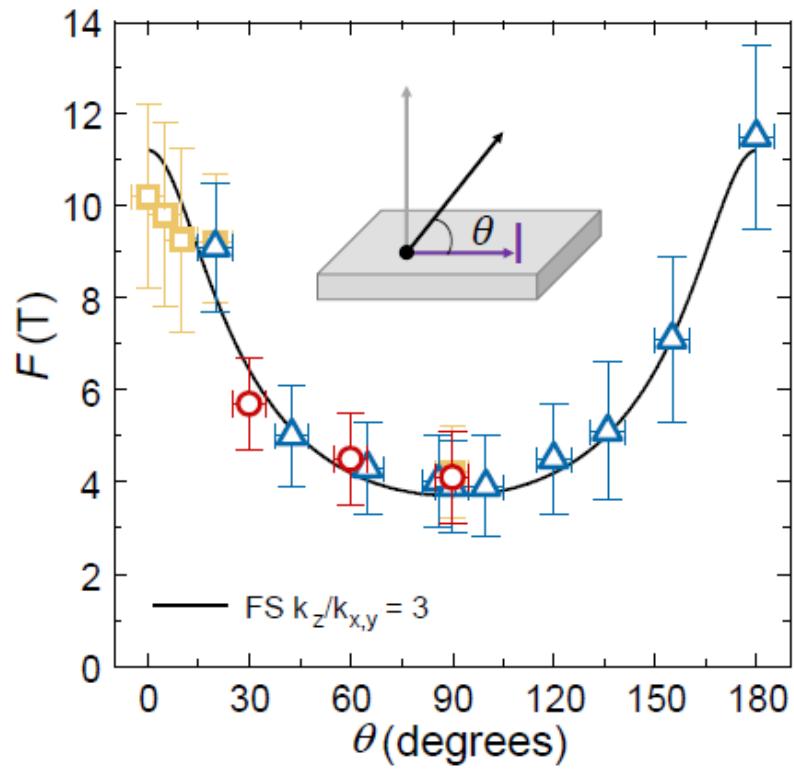
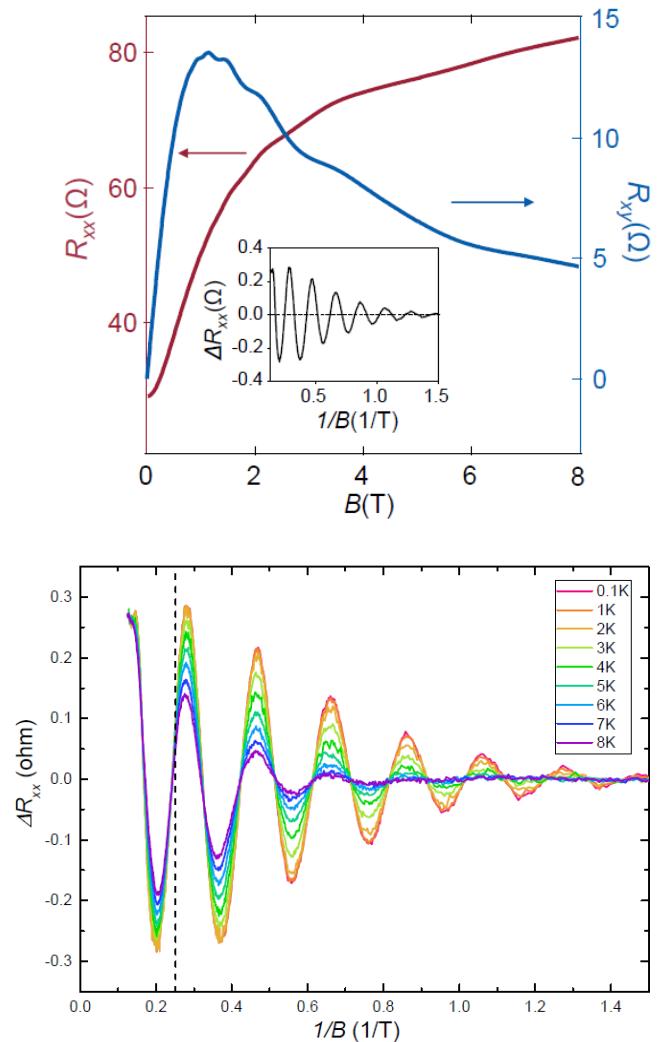
Angle resolved photoemission spectroscopy (ARPES):



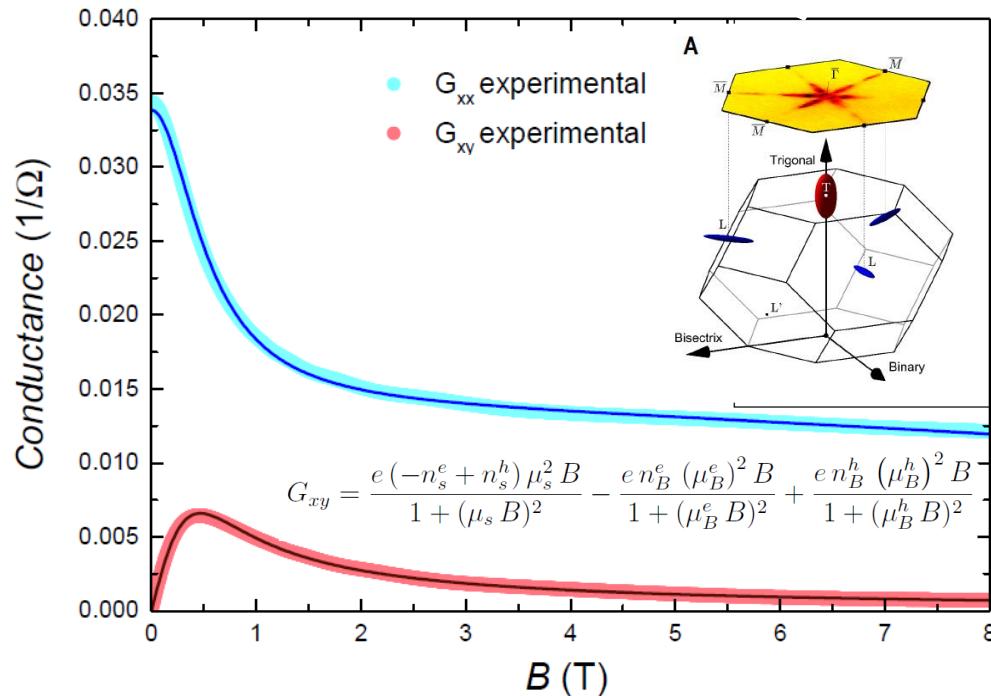
electrons:  $1-2 \times 10^{12} \text{ cm}^{-2}$

holes:  $1-2 \times 10^{12} \text{ cm}^{-2}$

# Quantum oscillations of bulk holes



# Multiband magnetotransport

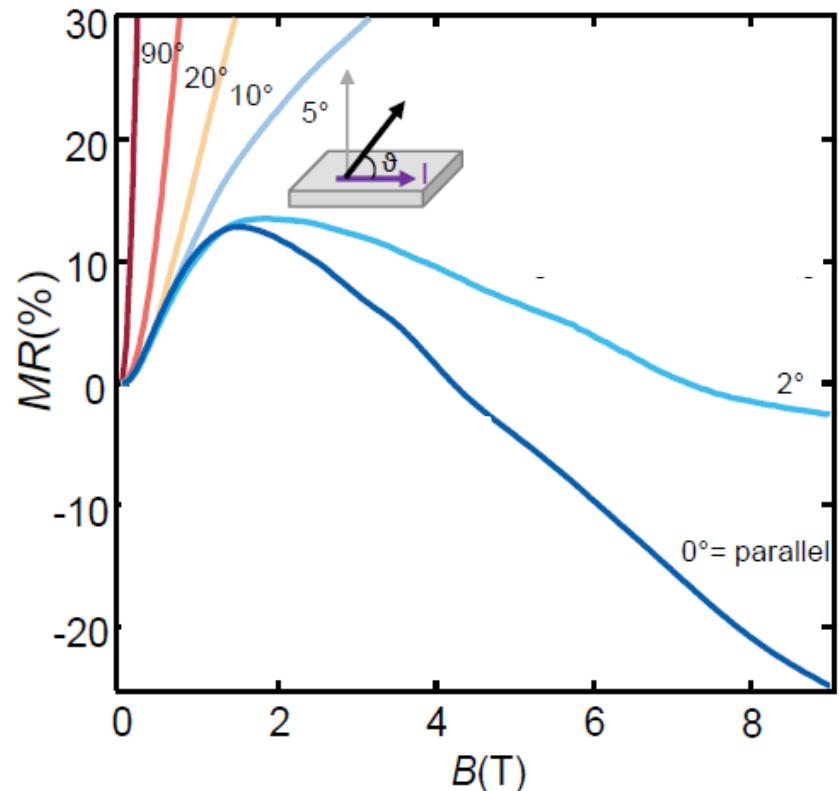
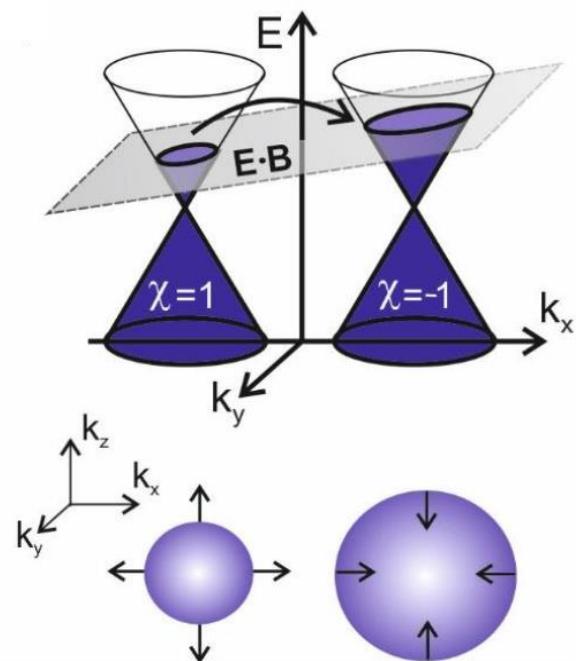


Electrons in the bulk Dirac cone:  $0.2 \times 10^{17} \text{ cm}^{-3}$

$$E_F = \hbar (\pi^2 n_e v_1 v_2^2)^{1/3} = 16 \text{ meV}$$

# Chiral anomaly in $\text{Bi}_{1-x}\text{Sb}_x$

Polarization of chirality:



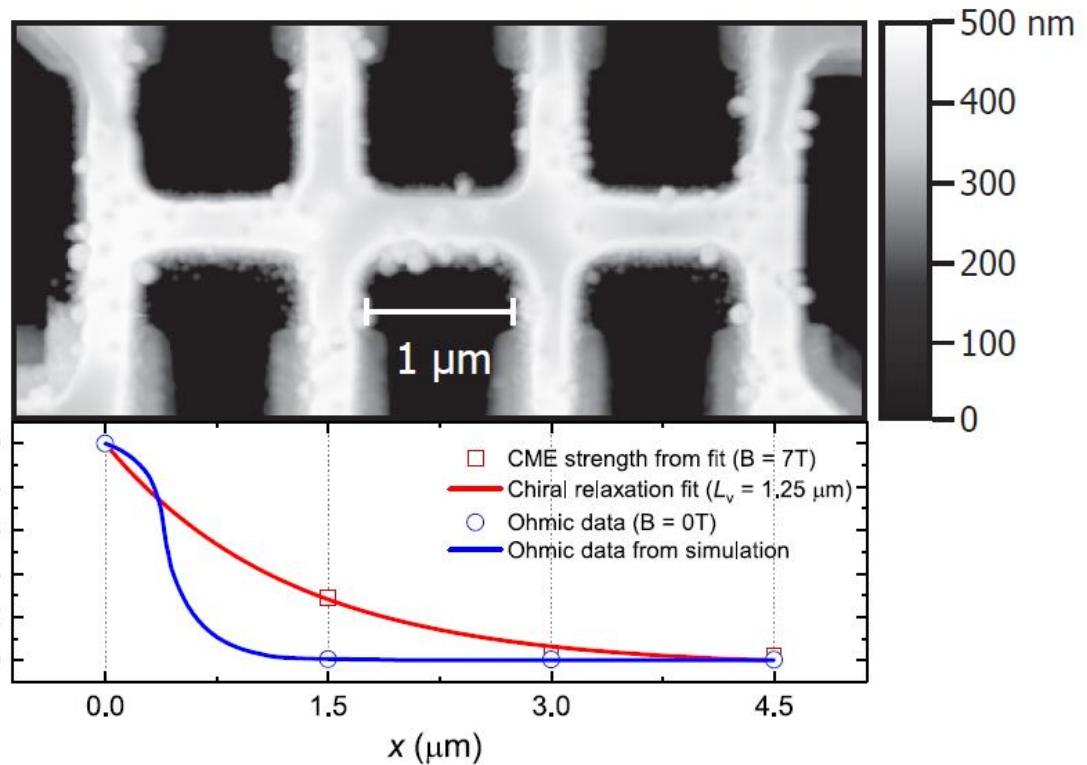
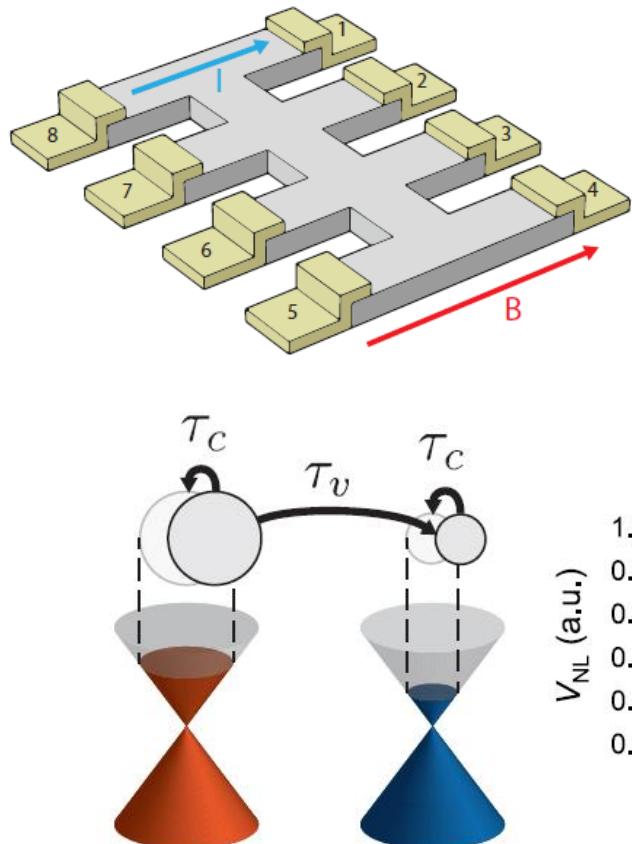
Analogy pion decay:



Negative longitudinal magnetoresistance (LMR)

3% ✓    4% ✗

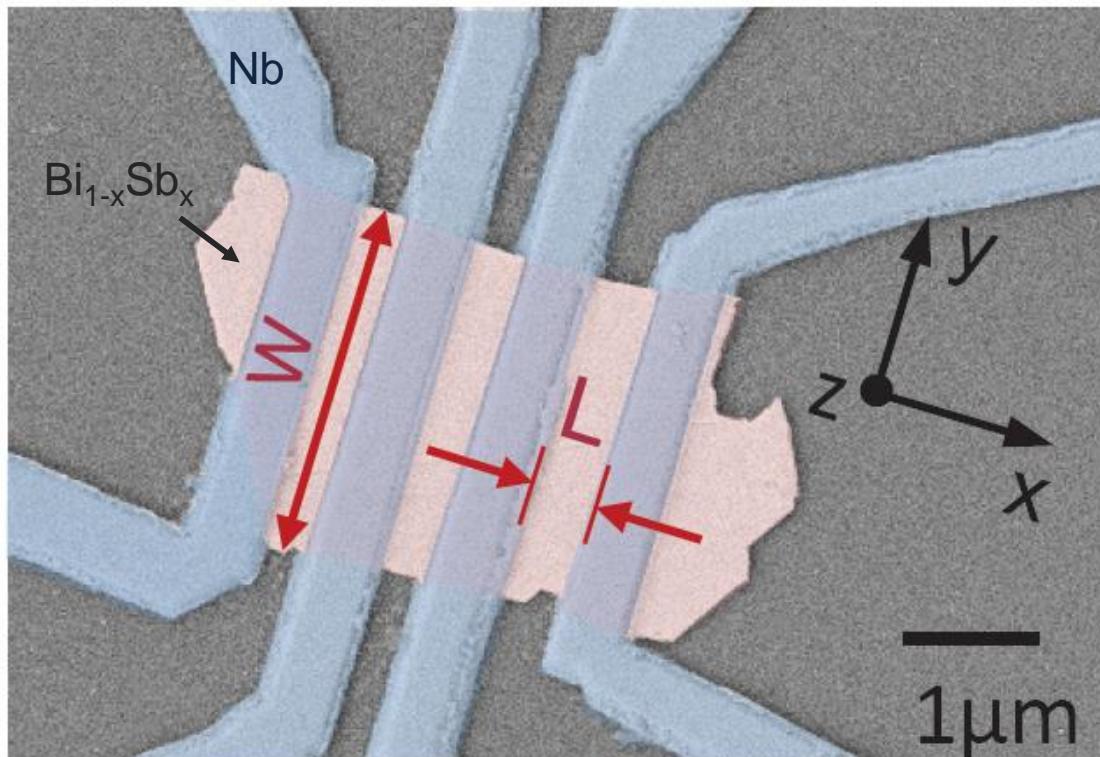
# Monopole diffusion



$$\frac{\partial \rho_{\pm}}{\partial t} = f(\Omega_{\mathbf{k}}) \frac{e^2}{4\pi^2\hbar^2} \mathbf{E} \cdot \mathbf{B}$$

Phys. Rev. B 99, 085124 (2019)

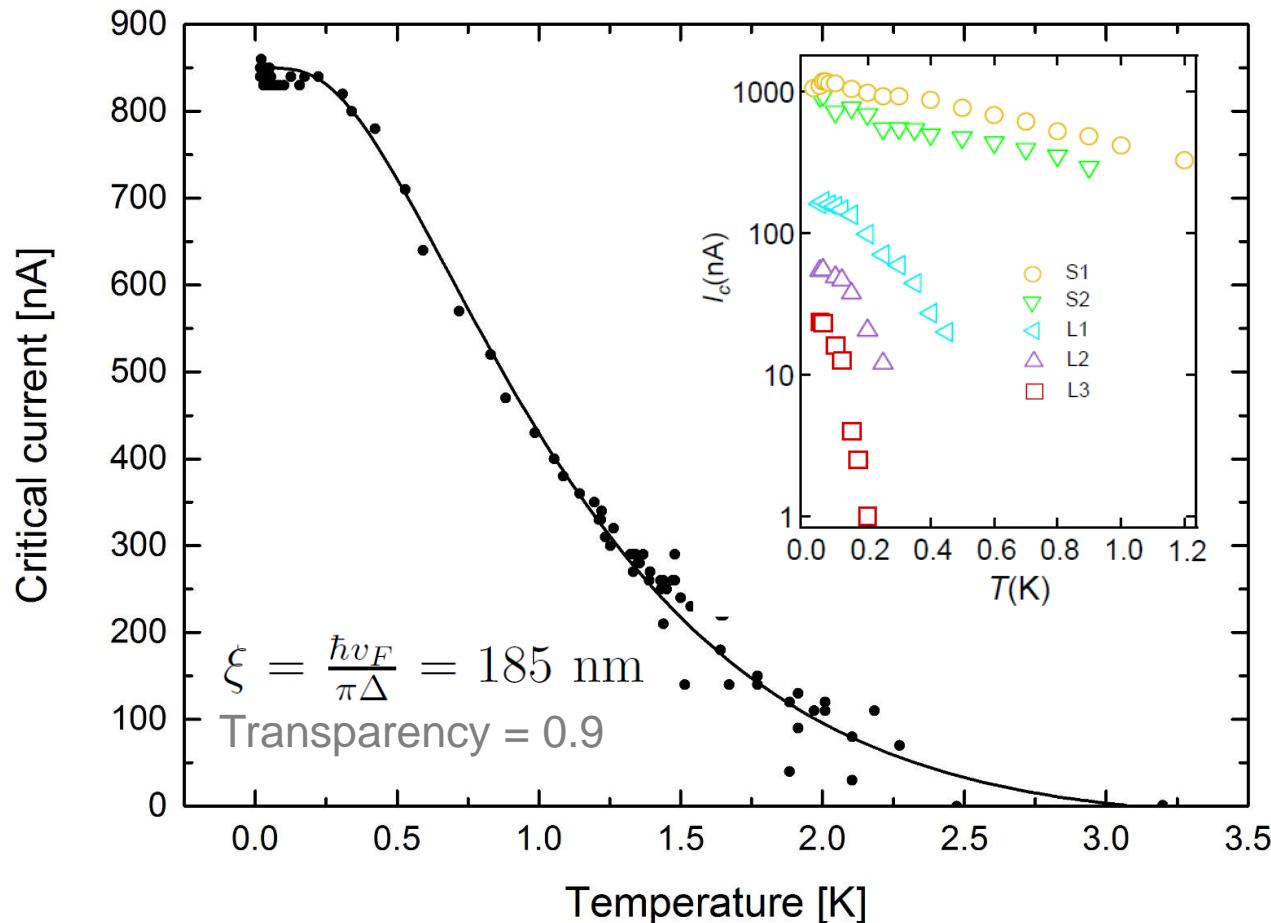
# $\text{Bi}_{1-x}\text{Sb}_x$ Josephson junctions



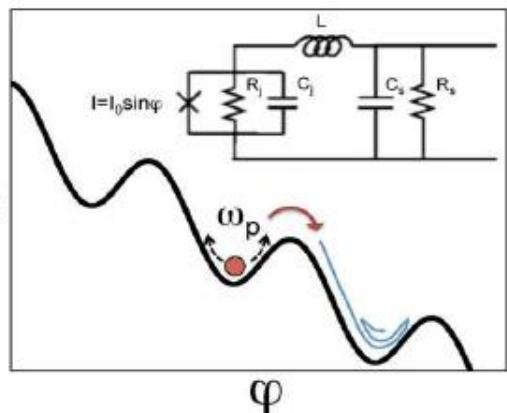
Nature Mater. 17, 875 (2018)

# Bi<sub>1-x</sub>Sb<sub>x</sub> Josephson junctions, x=0.03

Ballistic 500 nm Josephson junctions:

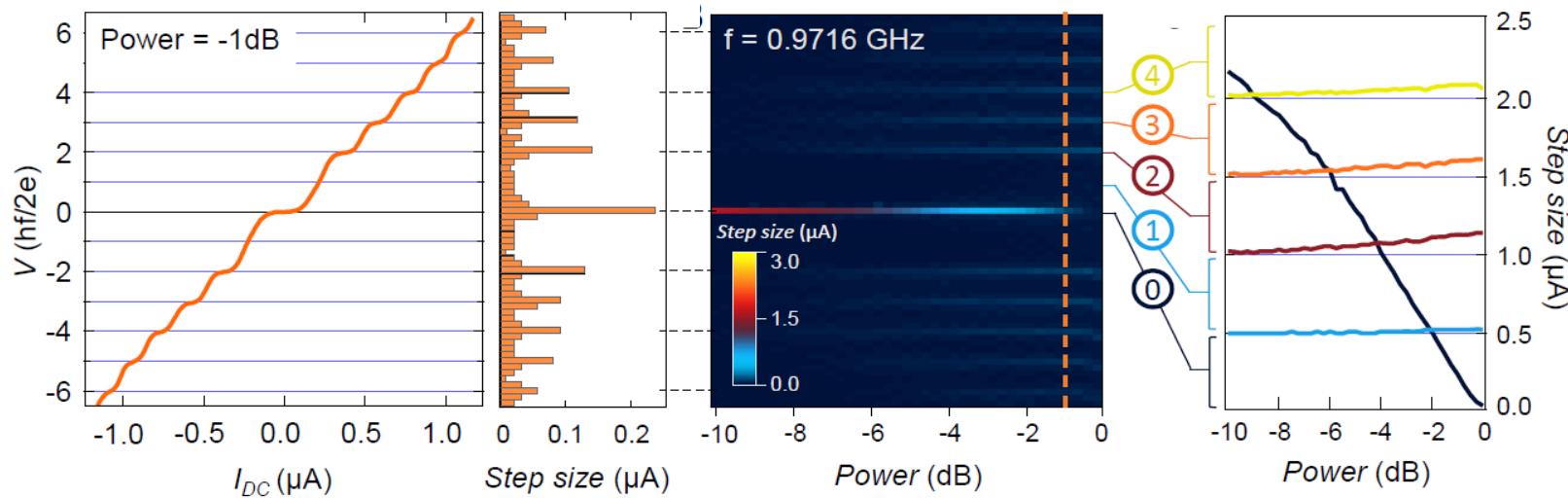
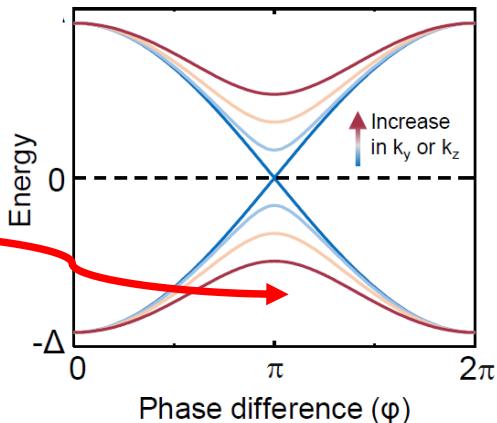


# Observation of $4\pi$ -periodic supercurrent

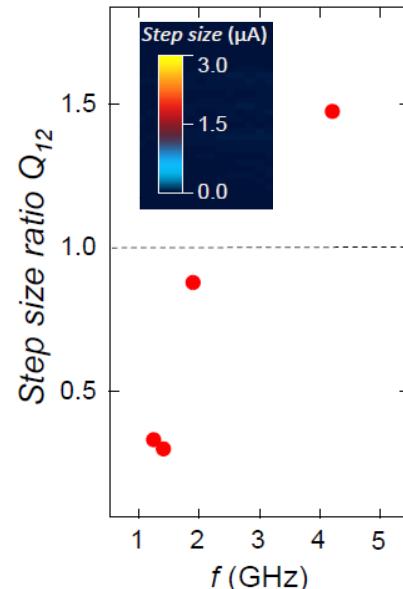
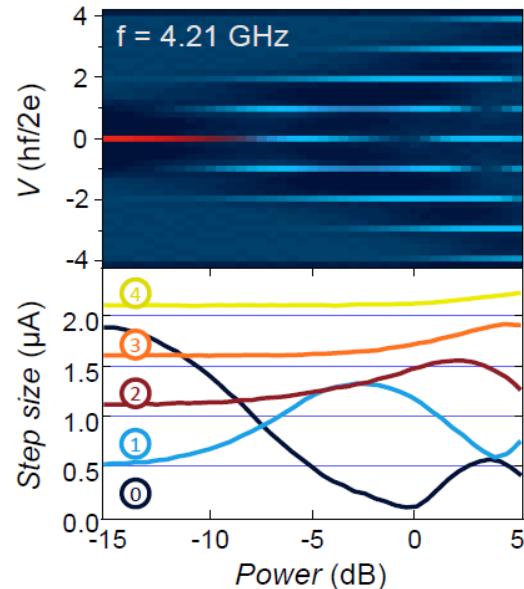
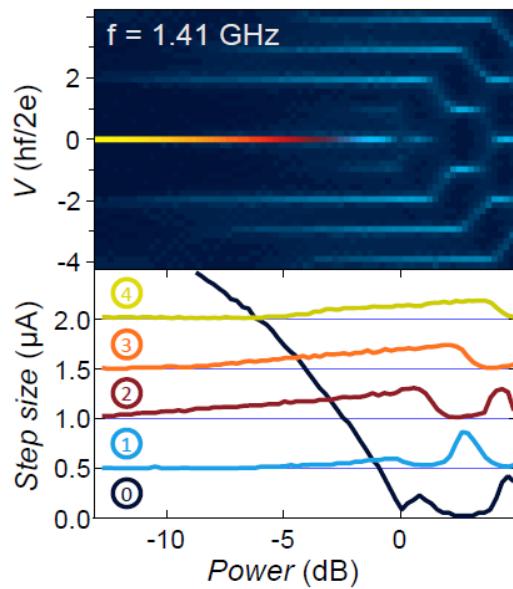


Microwave irradiation:

$$n \frac{h}{2e} \nu$$



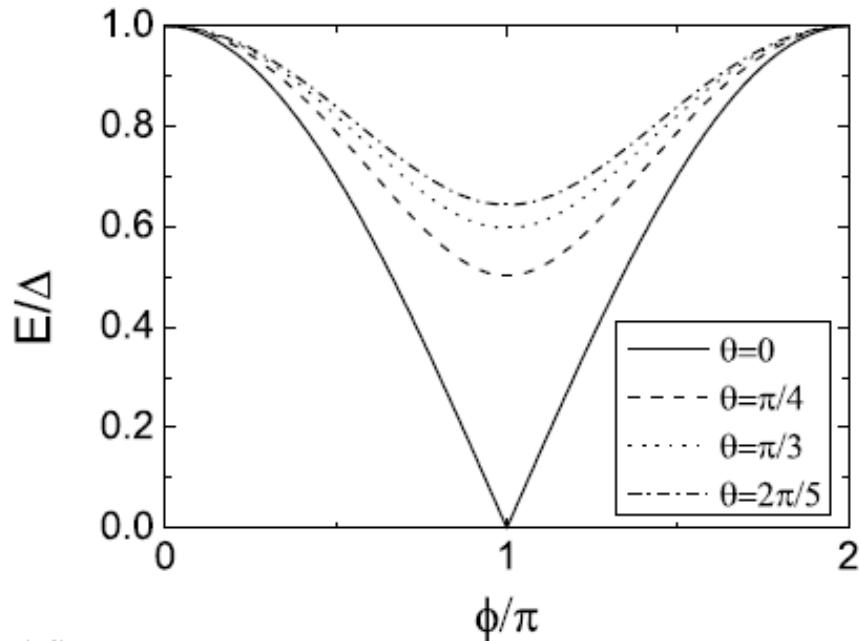
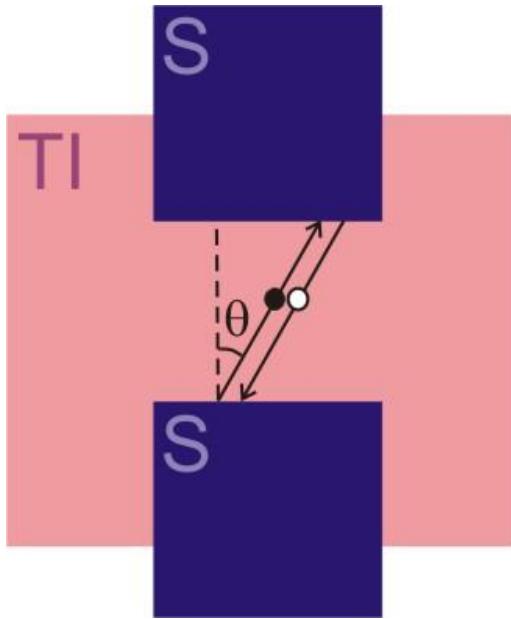
# Recovery of $2\pi$ -periodicity



$$f_c = \frac{2e}{h} R_N I_c^{4\pi} \quad I_c^{4\pi} = 0.2 \text{ } \mu\text{A}$$

From the critical frequency we obtain:  $I_{4\pi} > 20\% \text{ of } I_c$

# Non-perpendicular trajectories

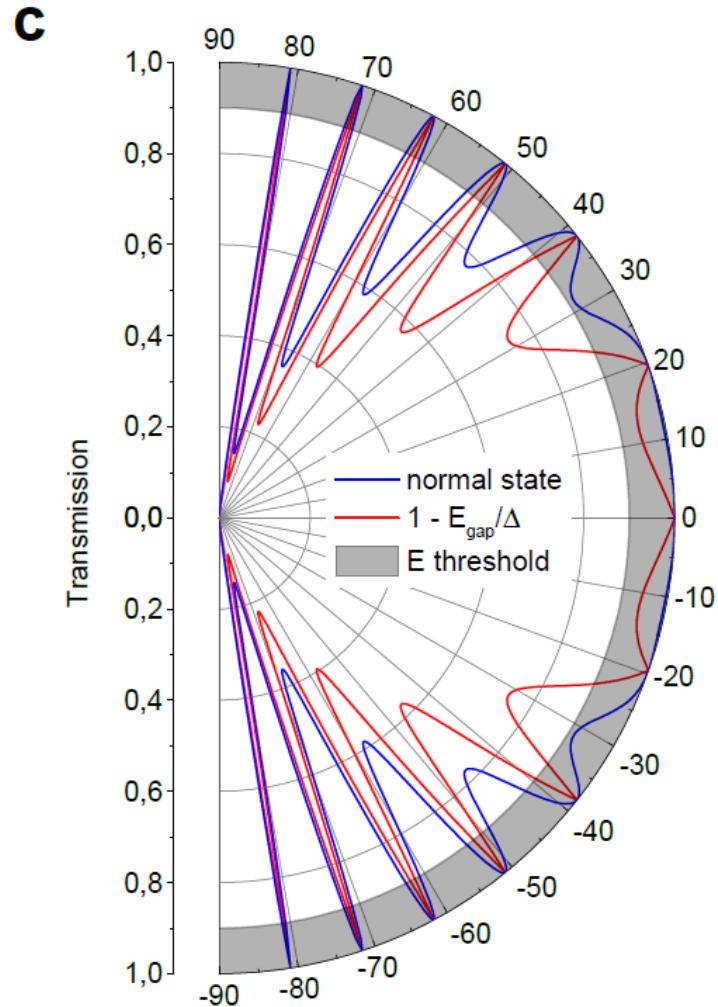
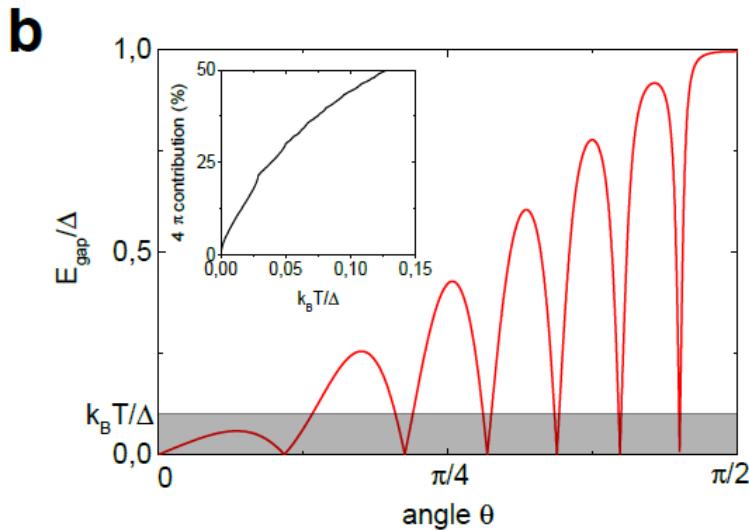
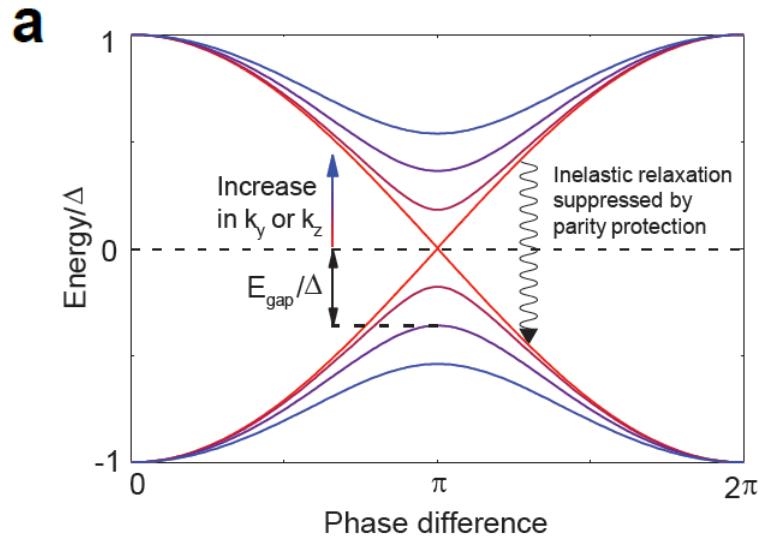


Towards single channel:

1. Quantization of transverse momentum
2. Including a magnetic layer

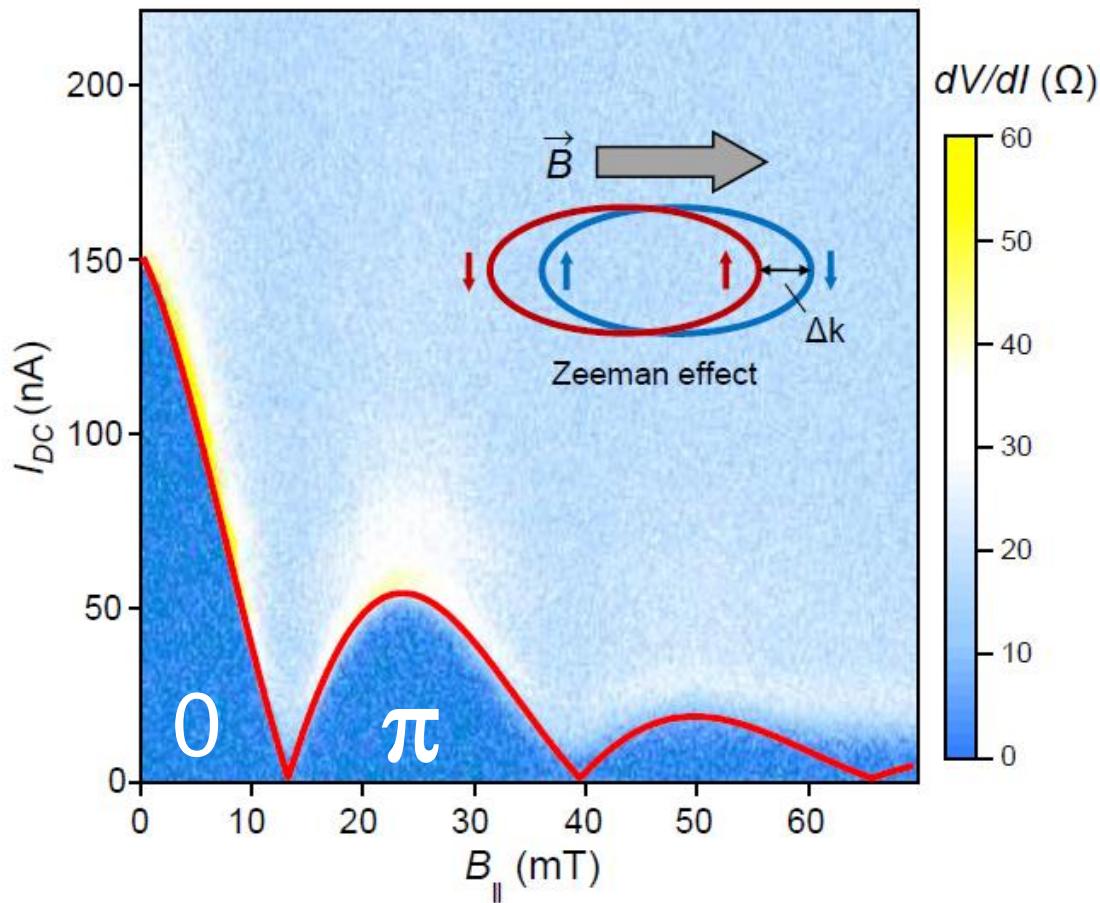
*M. Snelder, M. Veldhorst, A.A. Golubov, A. Brinkman, Phys. Rev. B 87, 104507 (2013)*

# 'Majorana' percentage

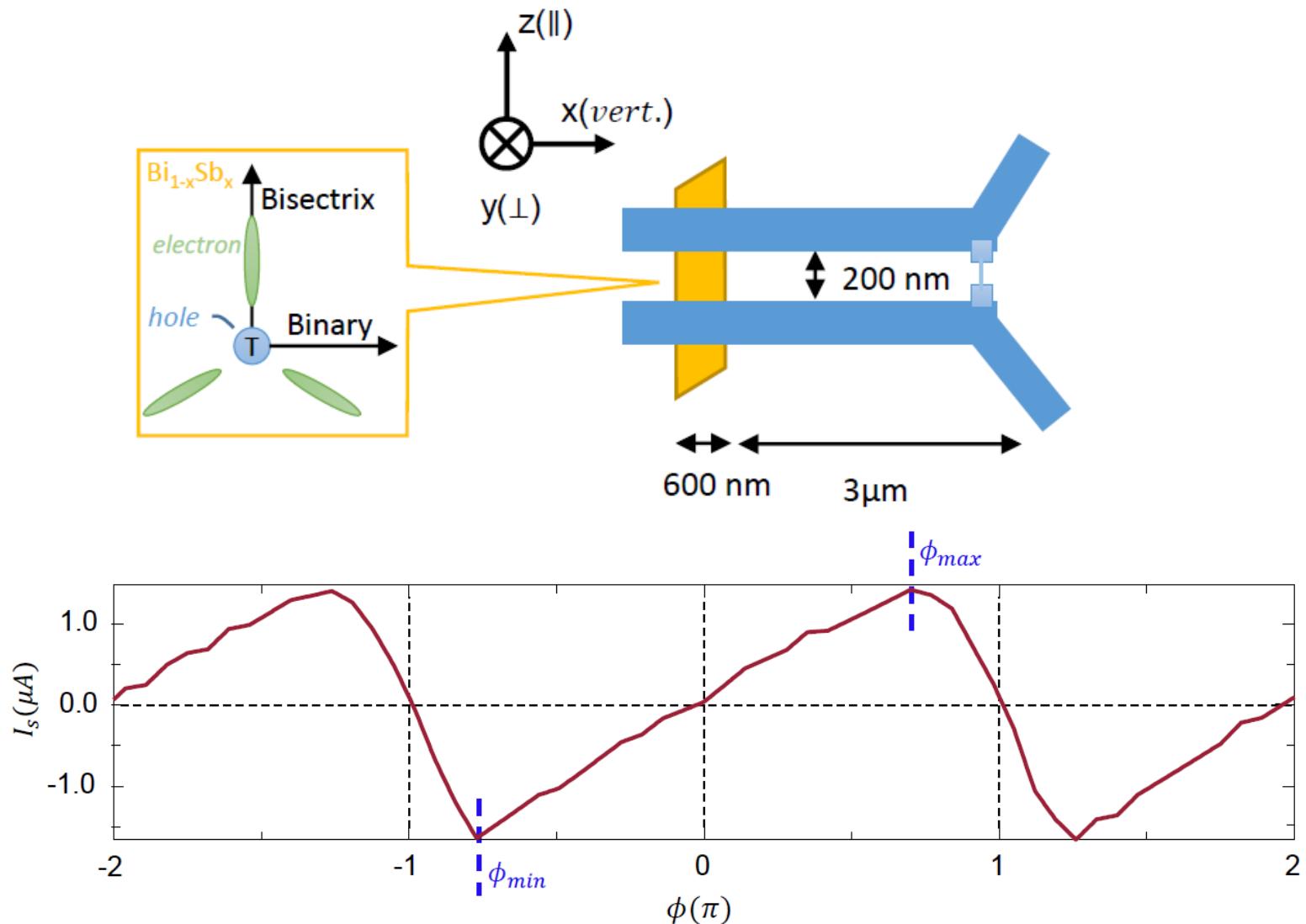


# Giant Zeeman effect in Dirac cone

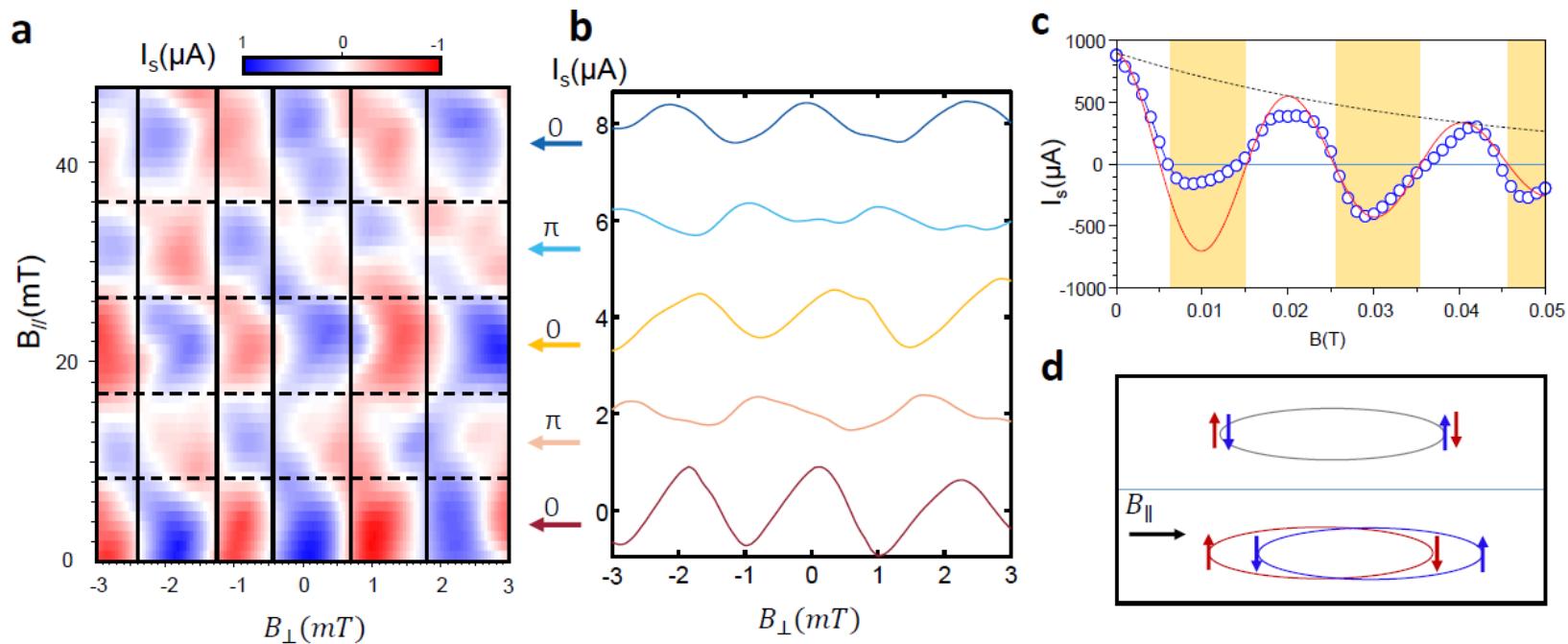
$$\Delta k_x = \frac{g\mu_B B_x}{\hbar v_F} \text{ with } g = 1000$$



# Asymmetric SQUID for current-phase relation



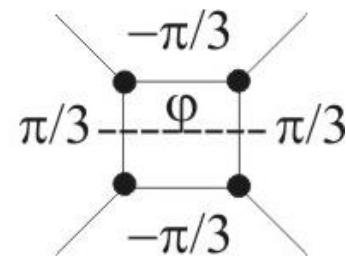
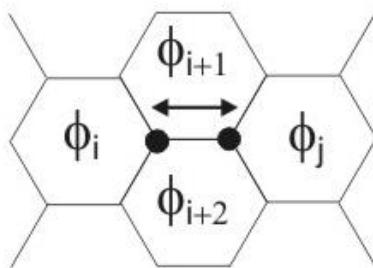
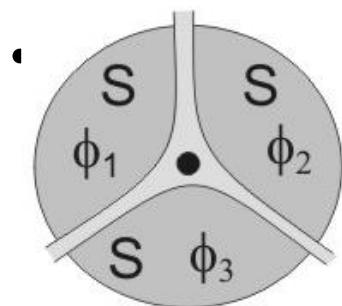
# 0 – pi transitions



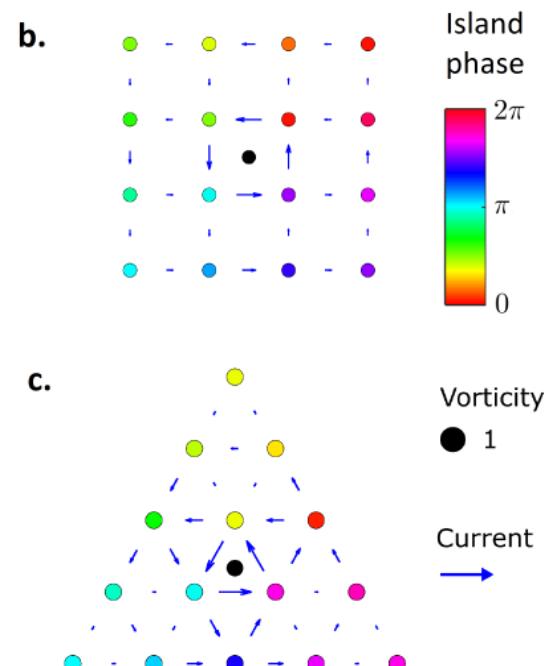
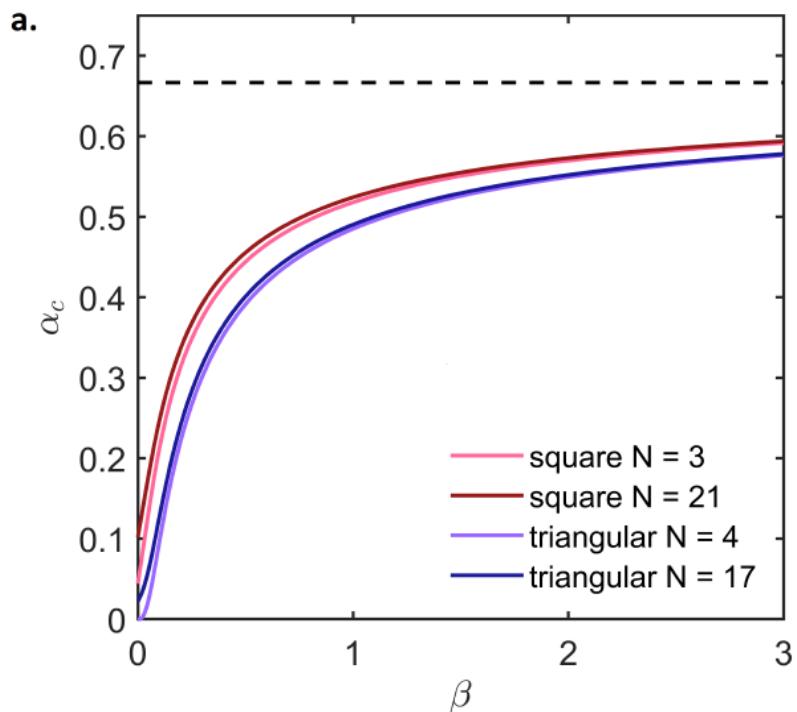
Phys. Rev. Lett. **123**, 026802 (2019)

# Conclusion and outlook

- Josephson supercurrent observed in several topological Josephson junctions
- Largest Majorana signal in the 3D Dirac semimetal
- Thin films & gating needed



# Outlook: Majorana vortex lattice

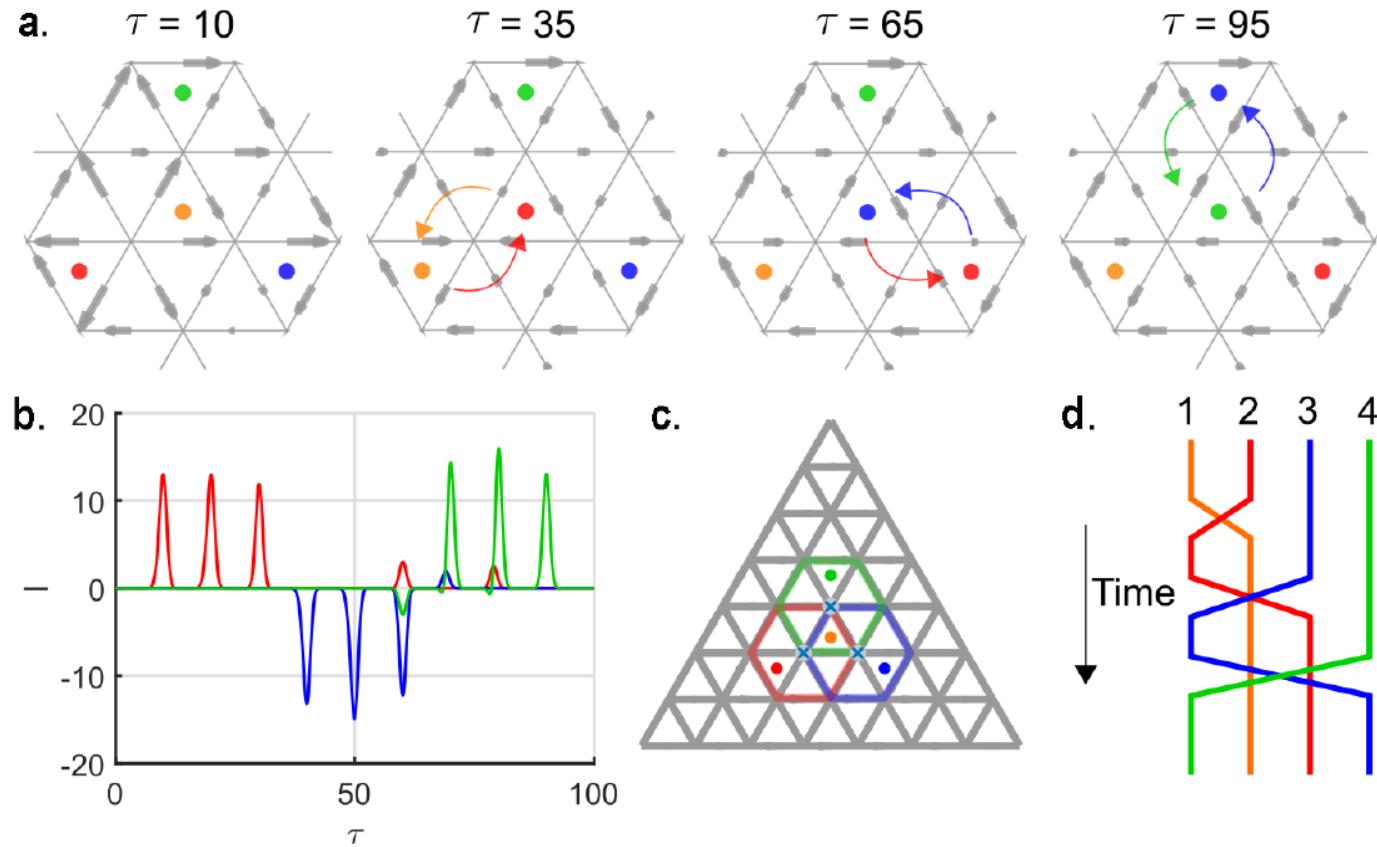


$$\alpha = I_c^{4\pi} / (I_c^{2\pi} + I_c^{4\pi})$$

$\beta$  = dimensionless self inductance of the loop

# Outlook: Majorana vortex lattice

Using *local current pulses* instead of flux biasing



Compatible with GHz rapid single flux quantum (RSFQ) technology?

# Outlook: Majorana vortex lattice

Using *global* current pulses and trap sites

