

**3rd International Workshop on Superconducting Sensor and Detectors
IWSSD 2016 14-17 Nov. 2016 TUKUBA**

Phase slips and Superconductor- Insulator transition in Nitride nanowire

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Scientific Reports **6**, (2016) 27001
Scientific Reports **4**, (2014) 5740



National Institute of Information and Communications Technology

Outline



Background and Motivation

Phase slips (Quantum and Thermal Activation)

Superconductor-Insulator transition(SIT)

Duality analysis

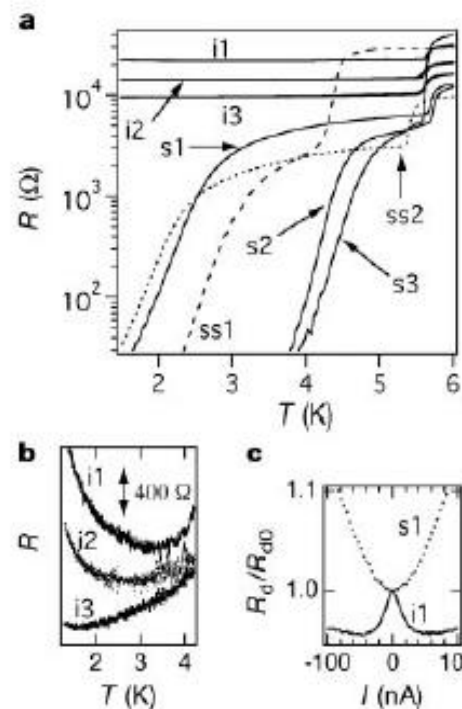
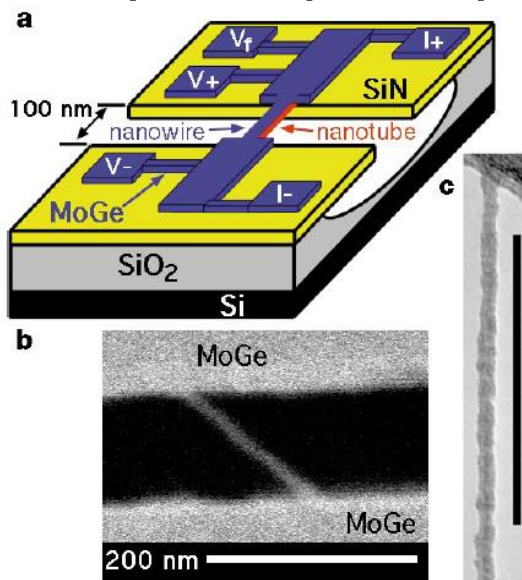
Future work and summary



Background



Quantum suppression of superconductivity (quantum phase slips)



A. Bezryadin, Lau, Tinkham
Nature 404 971 (2000)

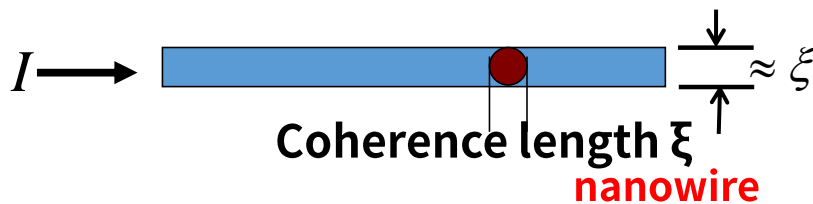


Phase slips



$$T < T_c$$

Normal conducting core due to fluctuations

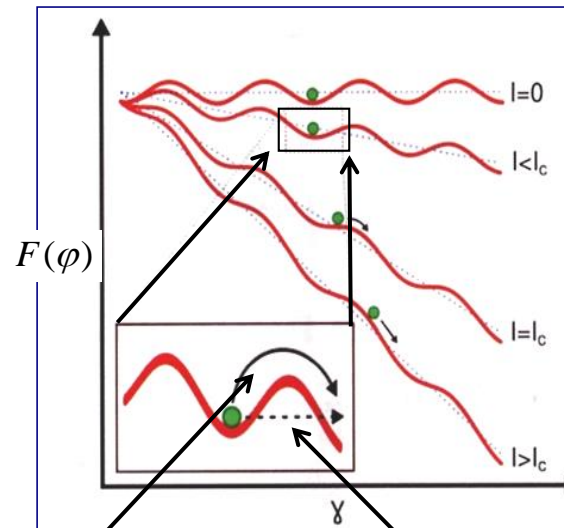


Potential barrier ΔF for Cooper pairs

Phase ϕ shifts occurs passing over ΔF

$$V = \left(\frac{\hbar}{2e}\right) \frac{d\phi}{dt}; \quad \text{Josephson equation}$$

$$R(T = 0) \neq 0$$



Thermal Activation
 →TA Phase Slip
 TAPS

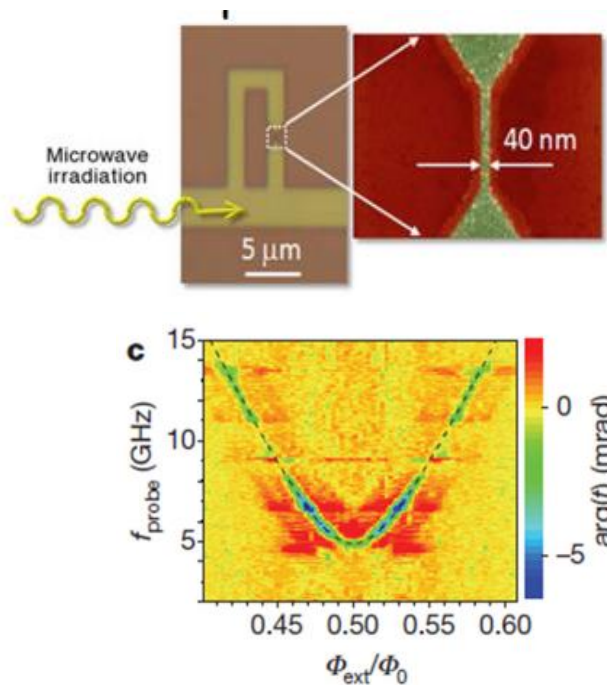
Quantum Tunneling
 →Q Phase Slip
 QPS



Coherent Quantum Phase slips

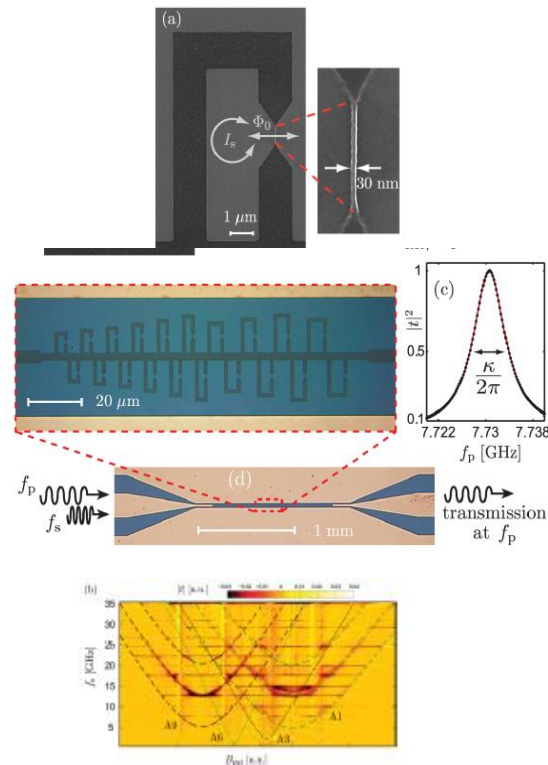


InOx nanowire



Astafiev et al.
 Nature 484 355-358 (2012)

NbN nanowire



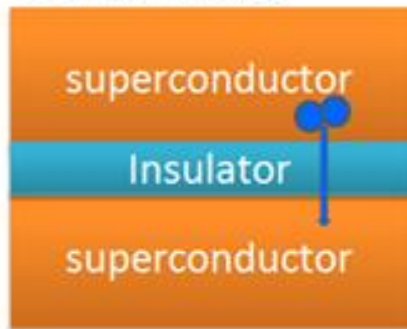
Peltonen et al.
 Phys. Rev. B (2013)



Duality



Josephson Junction
 Cooper pair tunneling

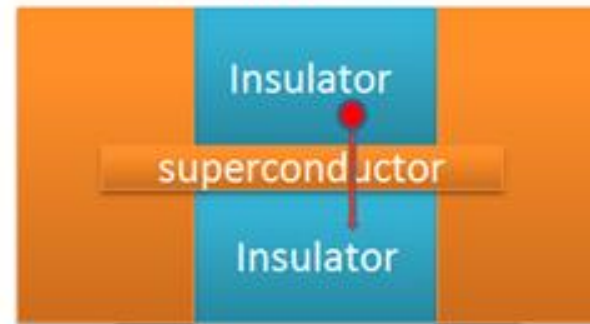


$$H_{JJ} = E_C(n - n_c)^2 - \left(\frac{E_J}{2} \sum_n |n+1\rangle\langle n| + h.c. \right)$$

Josephson junction

Charging energy E_C
 Josephson energy: E_J
 Phase : Φ
 Voltage : V

QPS Junction
 Phase tunneling



$$H_{QPS} = E_L(n - f)^2 - \left(\frac{E_S}{2} \sum_n |n+1\rangle\langle n| + h.c. \right)$$

QPS nanowire

Inductance energy E_L
 QPS energy: E_S
 Charge : $2e$
 Current : I



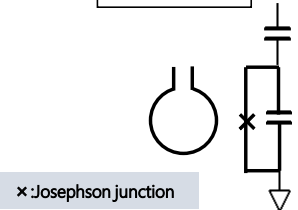
Some devices(ex. Qubits) without Josephson junction



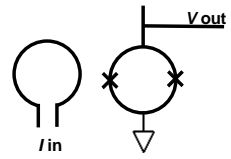
Duality devices

Josephson junction devices

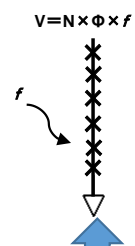
Charge Qubit



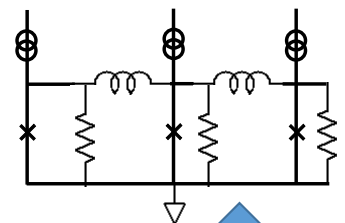
DC SQUID amp



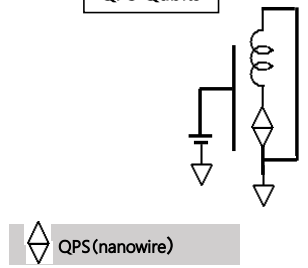
Josephson Voltage standard



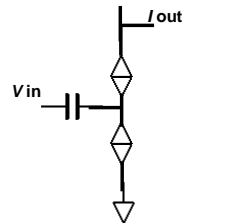
SFQ digital Circuits



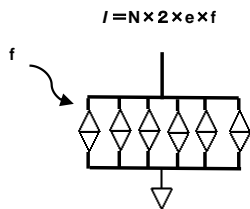
QPS Qubits



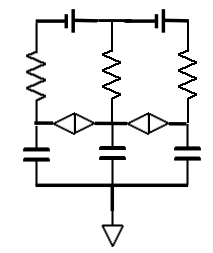
Hi reso. electrometer



QPS Current Standard



Cooperpair logic/memory

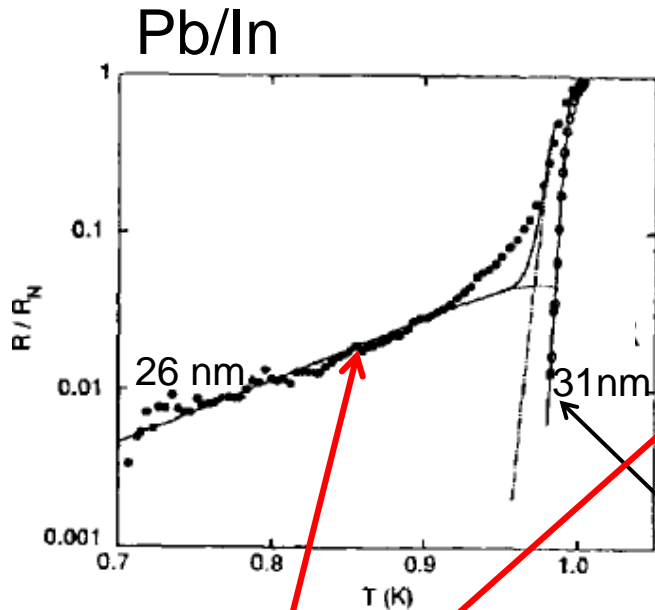


Duality

QPS devices



Phase Slip (R-T)



Physica B 203 (1994) 460-466

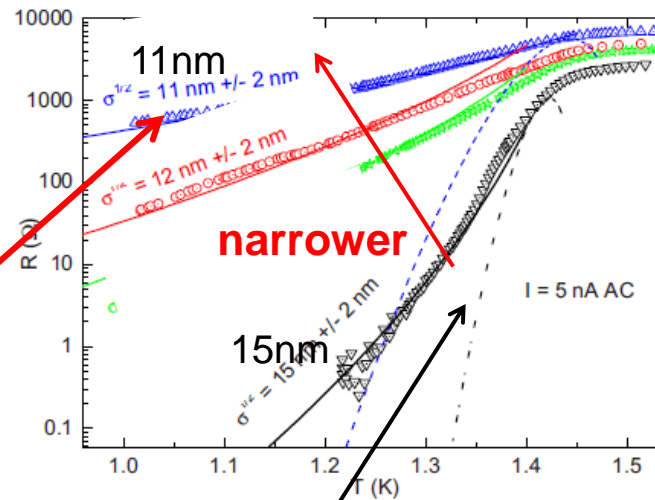
N. Giordano

Quantum phase slip (QPS)
• Schon and Zaikin (1990)



Al

PHYSICAL REVIEW B 77, 054508 (2008)
ZGIRSKI *et al.*



Thermal Activated phase slip (TAPS)

- Langer, Ambegaokar (1967)
- McCumber, Halperin (1970)

TA model & MQT model



J. S. Langer and V. Ambegaokar, Phys. Rev. **164**, 498 (1967); D. E. McCumber and B. I. Halperin, Phys. Rev. B **1**, 1054 (1970). **LAMH theory**

• thermal activation ;

$$R_{TA} = \frac{\phi_0 \Omega}{k_B T / \phi_0} \exp[-\Delta F / k_B T]$$

magnetic barrier

$$\Delta F = \sqrt{2} \sigma H_c^2 \xi / 3\pi$$

associated attempt frequency

$$\Omega = \frac{L}{2\pi^2 \xi \tau_{GL}} \left[\frac{3\pi \Delta F}{k_B T} \right]^{1/2}$$

GL relaxation time

$$\tau_{GL} = \frac{\pi \hbar}{8k_B (T_c - T)}$$

N. Giordano, Physica (Amsterdam) **203B**, 460 (1994)

change of energy scale $k_B T$ to \hbar / τ_{GL}

$$R_{MQT} = \left(\frac{\hbar}{4e^2} \right) \left(\frac{\hbar \Omega}{\hbar / \alpha \tau_{GL}} \right) \exp \left(- \frac{\Delta F}{\hbar / \alpha \tau_{GL}} \right)$$

D. S. Golubev and A. D. Zaikin, Phys. Rev. B **64**, 014504 (2001).

$$= \alpha R_Q \left(\frac{8\sqrt{3}L}{2\pi^{5/2}\xi_0} \right) 0.83 \left(\frac{LR_Q}{\xi_0 R_N} \right)^{1/2} \left[\left(1 - \frac{T}{T_c} \right)^{7/4} \frac{T_c}{T} \right] \exp \left[-0.83 \left(\frac{LR_Q}{\xi_0 R_N} \right) \left[\left(1 - \frac{T}{T_c} \right)^{1/2} \right] \right]$$

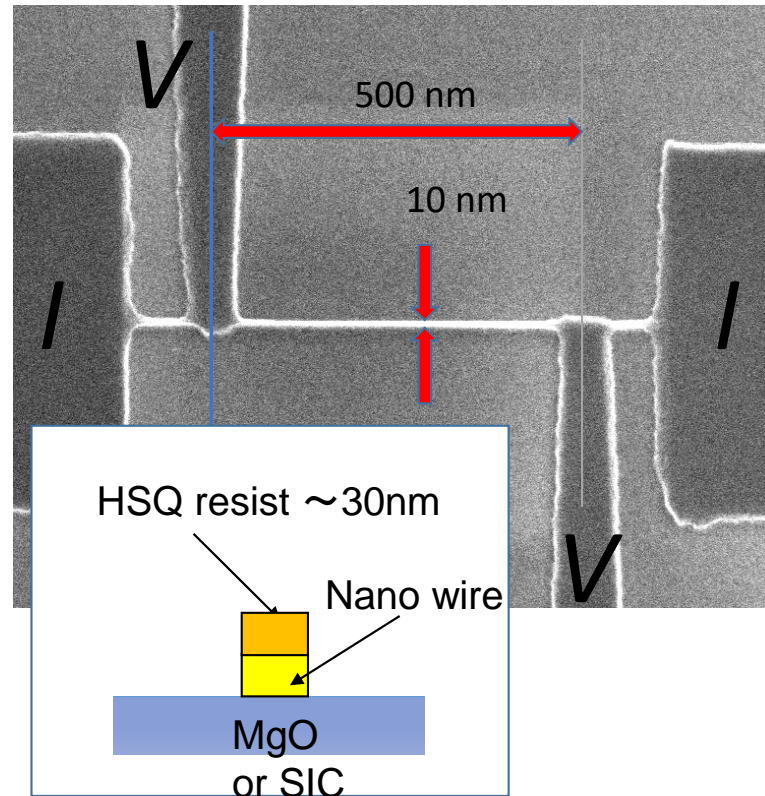


, where $R_Q = \hbar / 4e^2$

Experimental details



- Niobium Nitride (NbN), Niobium Titanium Nitride (NbTiN), Titanium Nitride (TiN)
Molybdenum Nitride (MoN)
MgO (100) substrate
3C-SiC/Si (100) substrate
- DC reactive magnetron sputtering method
- Reactive Ion Etching
- $d=2-10$ nm, $L = 250-1000$ nm
- $w = 10-300$ nm
- $R(T)$, $R(H)$, $I-V$ measurement



Nitride nanowire (NbN)



Epitaxial NbN nanowire

| | Lattice Type | Lattice Constants (nm) |
|--------|--------------|------------------------|
| MgO | NaCl-cubic | 0.439 |
| 3C-SiC | NaCl-cubic | 0.421 |
| NbN | NaCl-cubic | 0.436 |
| NbTiN | NaCl-cubic | 0.440 |

Lattice mismatch

$$\epsilon_{NbN-MgO} = \frac{a_{MgO} - a_{NbN}}{a_{NbN}}$$

1 %

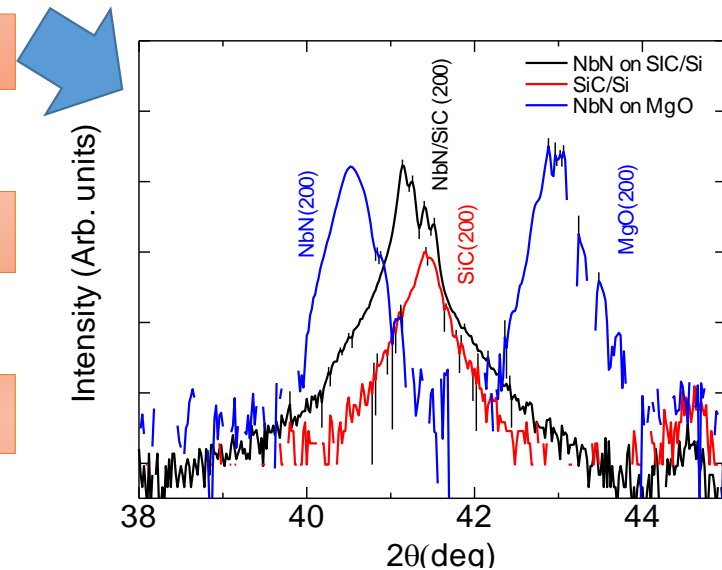
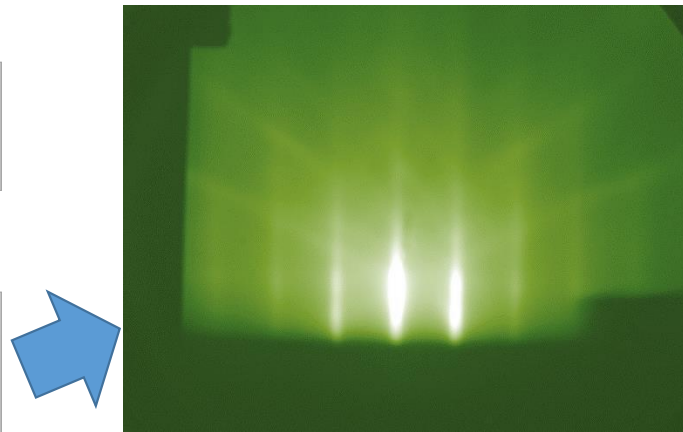
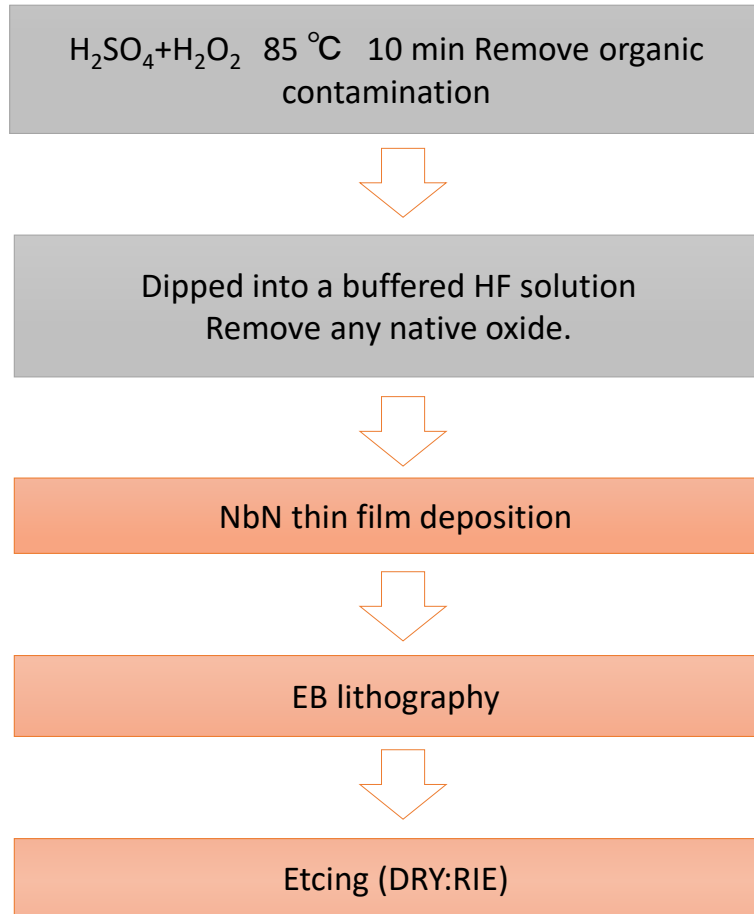
$$\epsilon_{NbN-SiC} = \frac{a_{SiC} - a_{NbN}}{a_{NbN}}$$

-0.7 %

Lattice mismatch between NbN and 3C-SiC/Si is smaller.



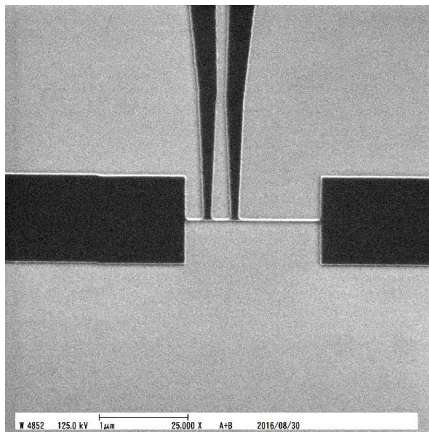
Sample preparation



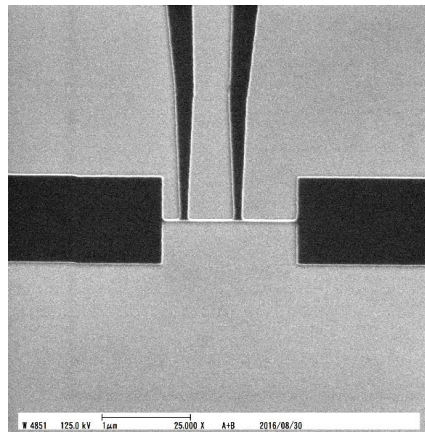
Nitride nanowire (NbN)



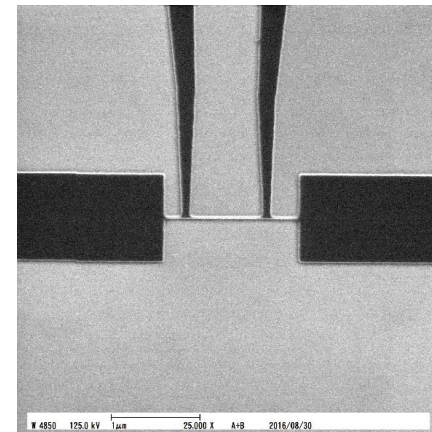
Using four-probe method to eliminate the contact resistance.



L = 300 nm



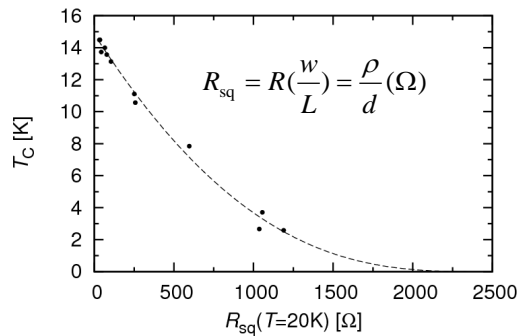
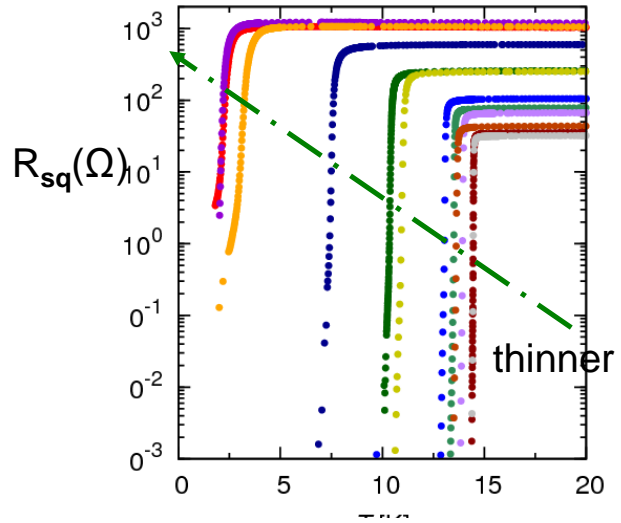
L = 600 nm



L = 900 nm

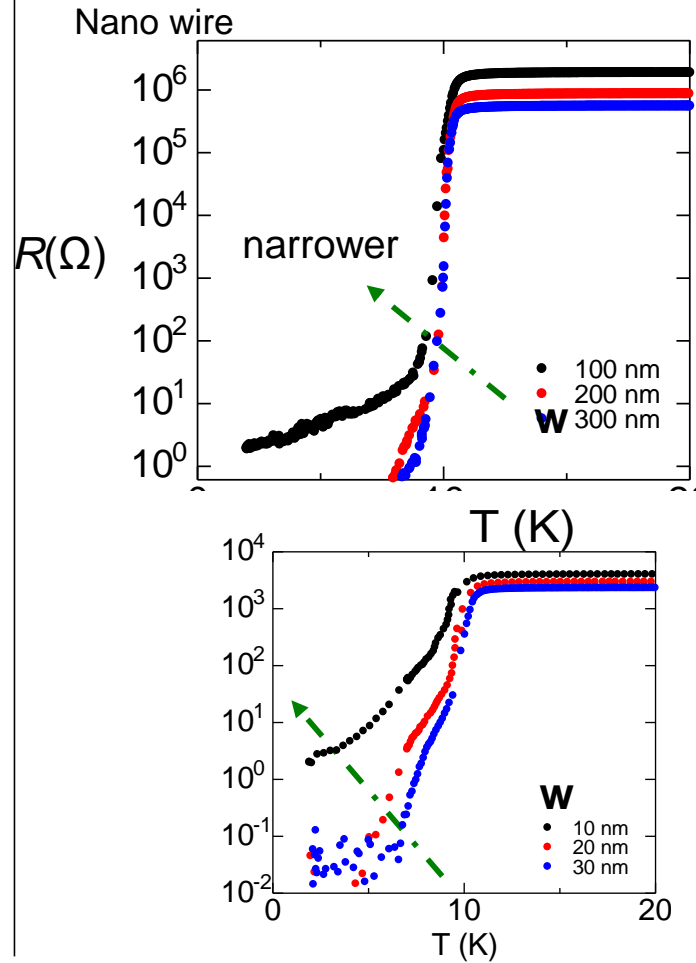


Size dependence

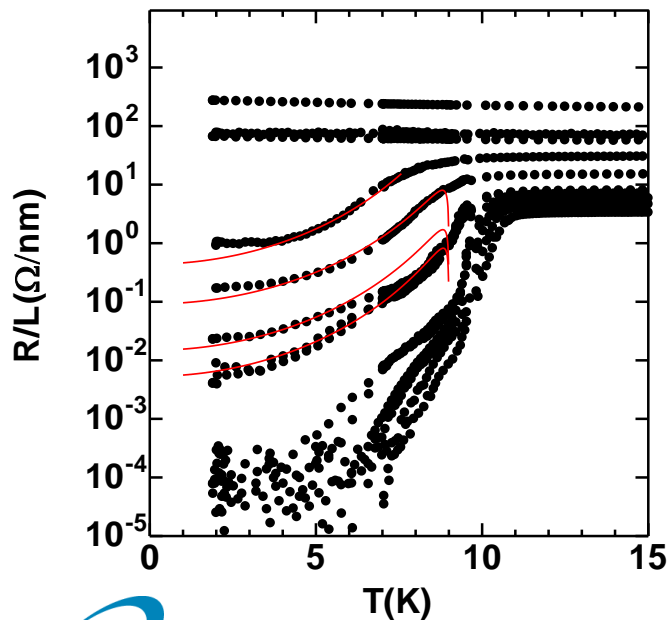


NbTiN 2D film $w=100\mu\text{m}$

$$R_{sq} = R\left(\frac{w}{L}\right) = \frac{\rho}{d}(\Omega)$$



QPS fitting



$$R_{QPS}(T) = \left(\frac{h}{4e^2} \right) \left(\frac{\hbar\Omega}{\hbar/\alpha\tau_{GL}} \right) \exp \left(- \frac{\Delta F}{\hbar/\alpha\tau_{GL}} \right)$$

$$R_{QPS}(T) = \alpha R_Q \left(\frac{8\sqrt{3}L}{2\pi^{5/2}\xi_0} \right) 0.83 \left(\frac{LR_Q}{\xi_0 R_N} \right)^{1/2} \left[\left(1 - \frac{T}{T_C} \right)^{7/4} \frac{T_C}{T} \right]$$

$$\exp \left[-0.83 \left(\frac{LR_Q}{\xi_0 R_N} \right) \left[\left(1 - \frac{T}{T_C} \right)^{1/2} \right] \right]$$

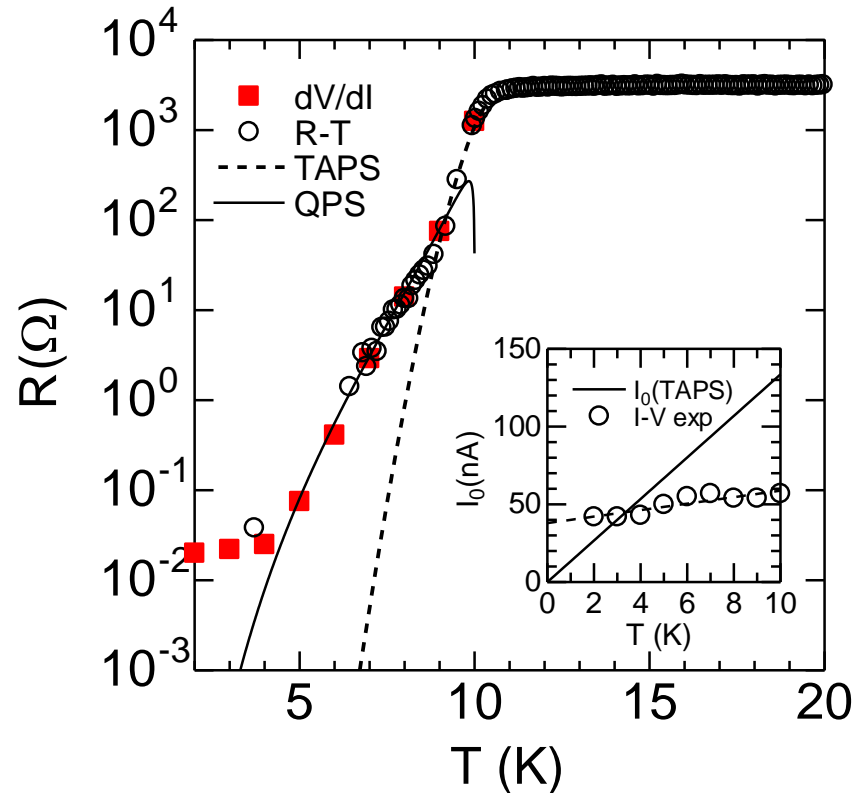
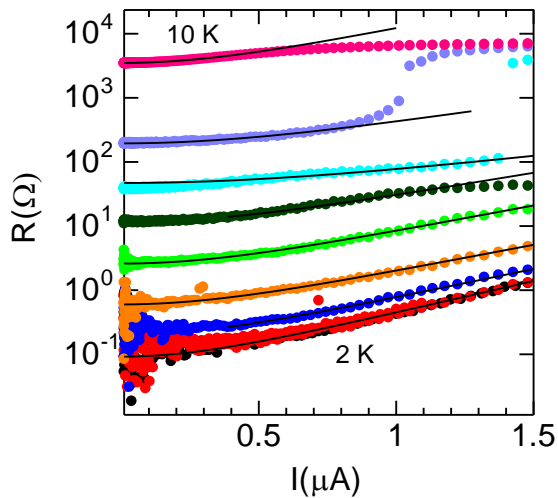
$$I_{QPS} = \left(\frac{2e}{\beta\pi\tau_{GL}} \right)$$

$$V_{QPS} = I_{QPS} R_{QPS} \sinh \left(\frac{I}{I_{QPS}} \right)$$

Zaikin et al., Phys. Rev. Lett. 78, 1552 (1997); D. S. Golubev and A. D. Zaikin, Phys. Rev. B 64, 014504 (2001).



TAPS fitting (I-V & R-T Curve)



$$dV/dI = R_{TAPS}(T) \cosh(I/I_0)$$

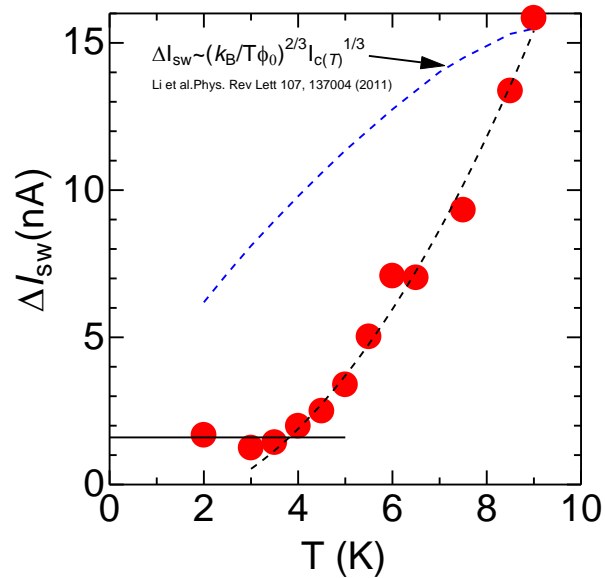
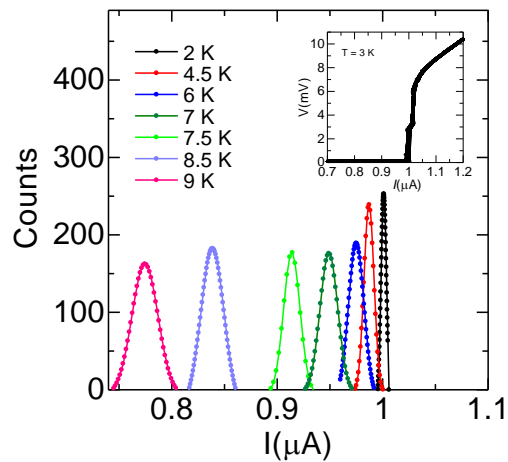
$$I_0(T) = 4ekT/h \quad (\text{By TAPS, } I_0 \text{ depend on only Temperature.})$$



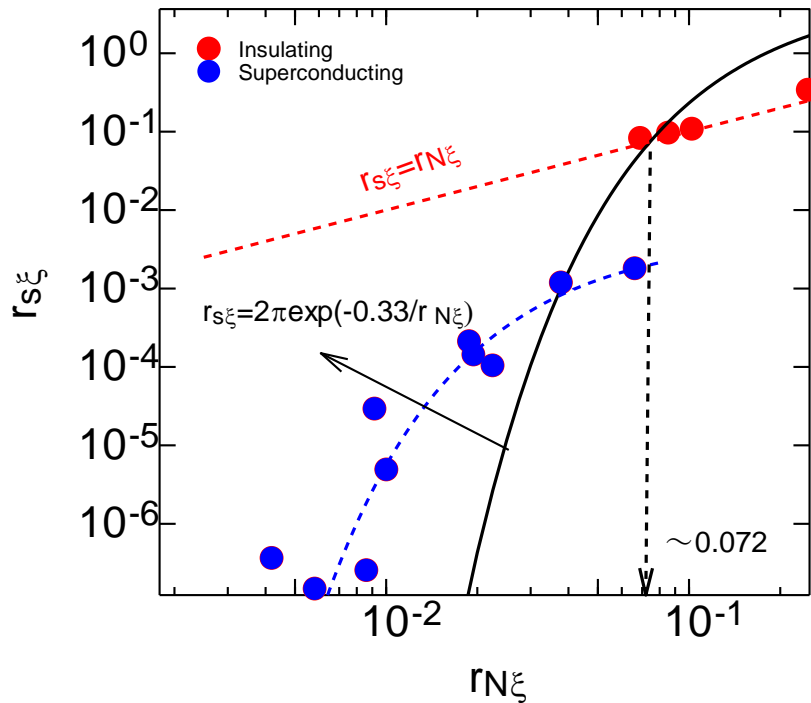
Switching Current



I_{sw} distribution at different temperature



Estimation of Cross over for SIT based on QPS model



superconducting state resistance R_s
 measured in units of R_Q for length in
 units of ξ

$$(R_{QPS} = R_s)$$

$$\begin{aligned} (R_s / R_Q)(\xi / L) &\equiv r_{S\xi} \\ &= 2\pi \exp[-0.33(R_Q / R^N)(L / \xi)] \\ &= 2\pi \exp[-0.33 / r_{N\xi}] \end{aligned} \quad (1)$$

$$\text{at cross over, } r_{S\xi} = r_{N\xi} \quad (2)$$

$$\Rightarrow r_{N\xi} = 1/13.3 \approx 0.072$$

$$(R^N / L)_{c.o.} = R_Q / 13.3\xi \approx 485\Omega / \xi(0)$$

$\xi(0) \approx 9nm$ for presents pecimens

$$(R^N / L)_{c.o.} \approx 54\Omega / nm$$



PS & Inductance energy : E_s, E_L

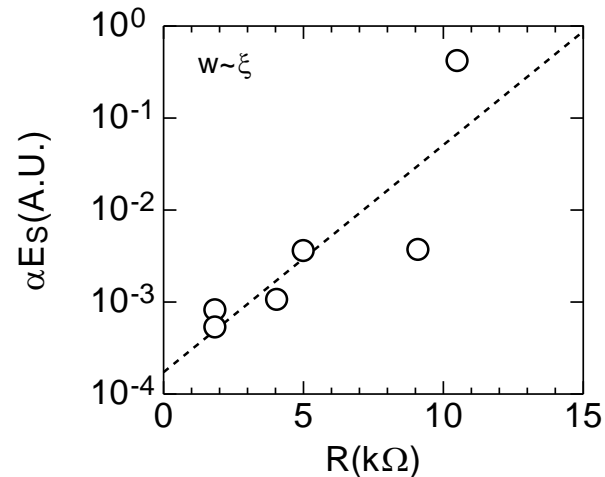
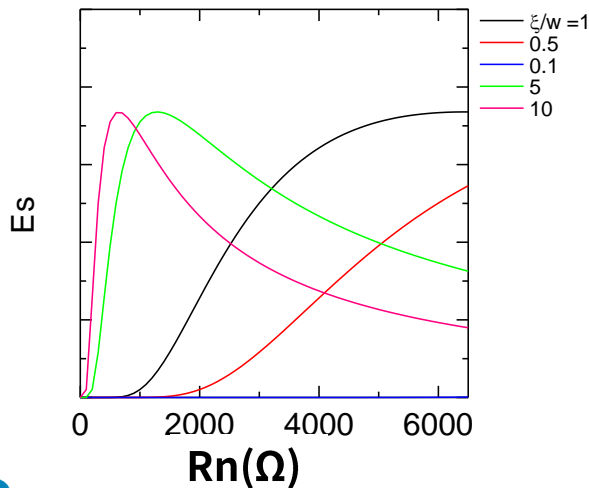


According the theory Quantum phase slip is dependent on sheet resistance.

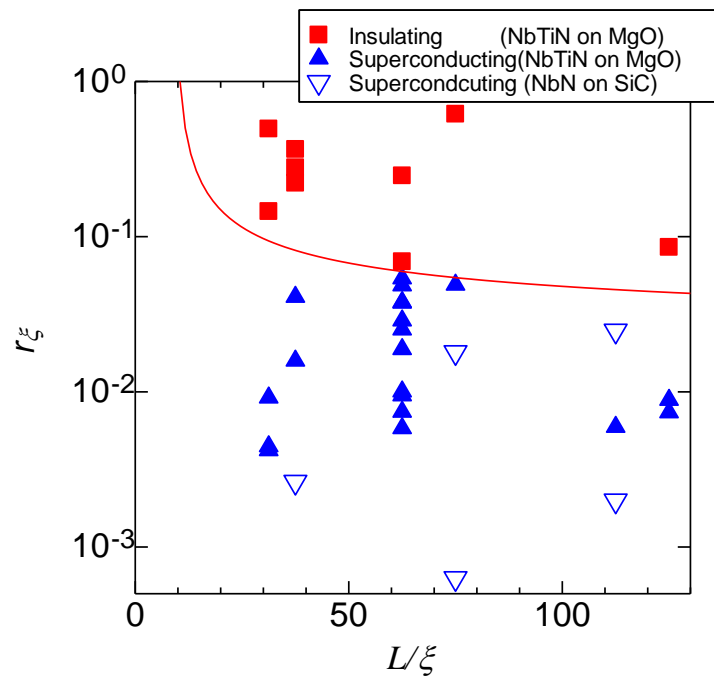
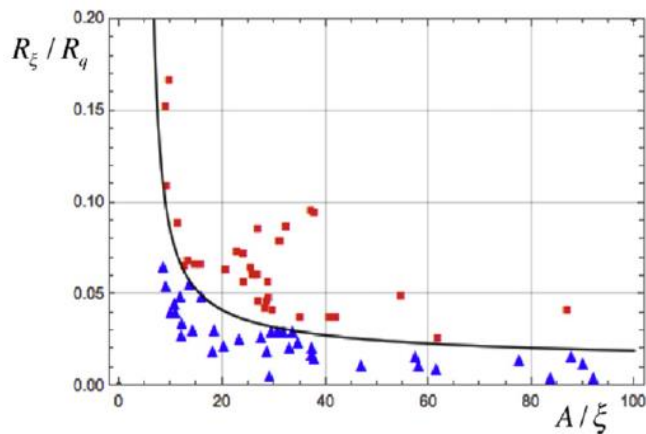
But it remains unclear whether the characteristics of the QPS depend on disorder

$$E_S = a \frac{A}{\xi} k_B T_C \frac{R_Q}{R_\xi} \exp\left(-b \frac{R_Q}{R_\xi}\right) \quad L = \frac{\hbar R_n}{\pi \Delta}$$

$$R_\xi = R_n \frac{\xi}{w} \quad E_L = (h/2e)^2 / 2L$$



Phase diagram of S-I transition



$$E_s / E_{Li} = (a\lambda^2 / 17.4) \exp(-b / r_\xi) = \alpha_c$$

$$r_\xi = (R^N / R_Q) / (L / \xi)$$

$$r_\xi(\lambda) = b / \ln(a\lambda^2 / 17.4\alpha_c)$$

$$\equiv b / \ln(c\lambda^2 / 17.4)$$

MoGe nanowire $E_s/E_L = 0.3$

NbSi nanowire $E_s/E_L = 0.6$

Mooij et al. New J.Phys 17 (2015) 033006.

**NbN and NbTiN nanowire
 $0.2 < E_s/E_L < 0.5$**



To realize CQPS



Required property of Material

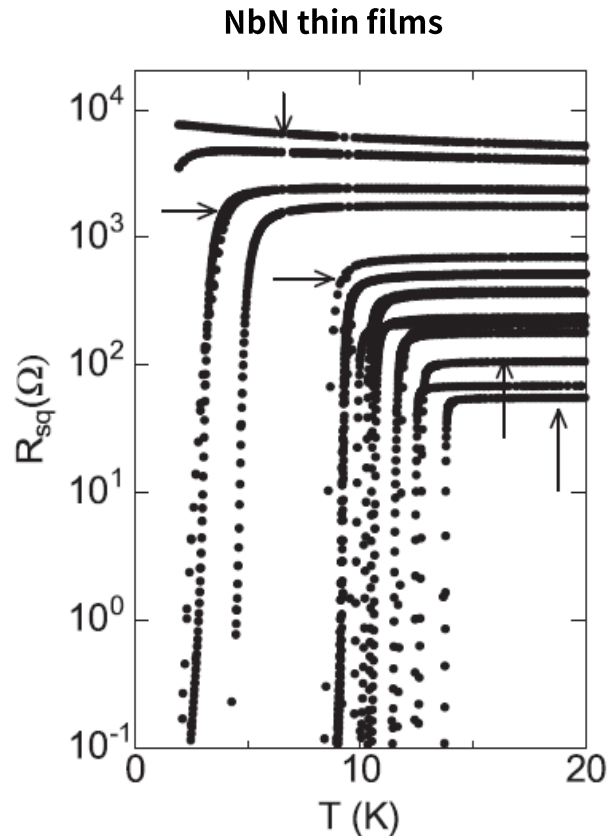
- High degree of disorder
- homogeneous system
- Oxide or **Nitride**

Phase slips energy(E_s)

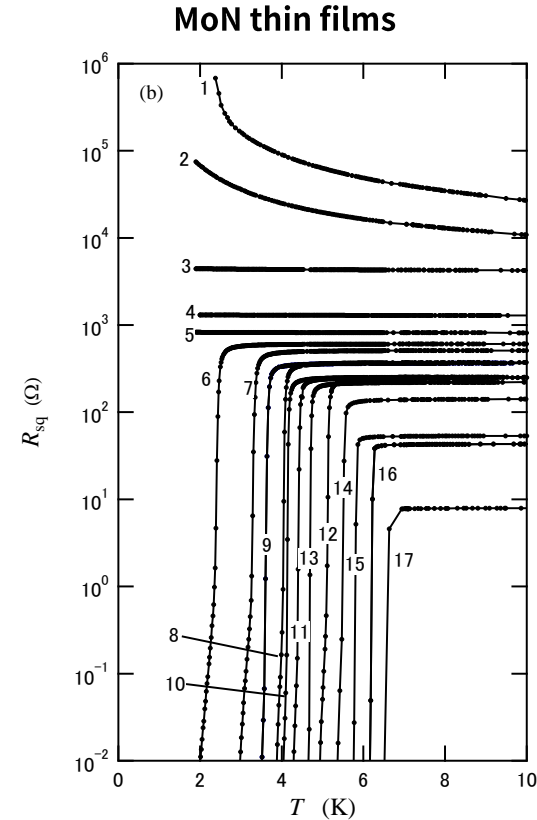
- **Close to SI transition**
- narrow wires



S-I transition for 2D



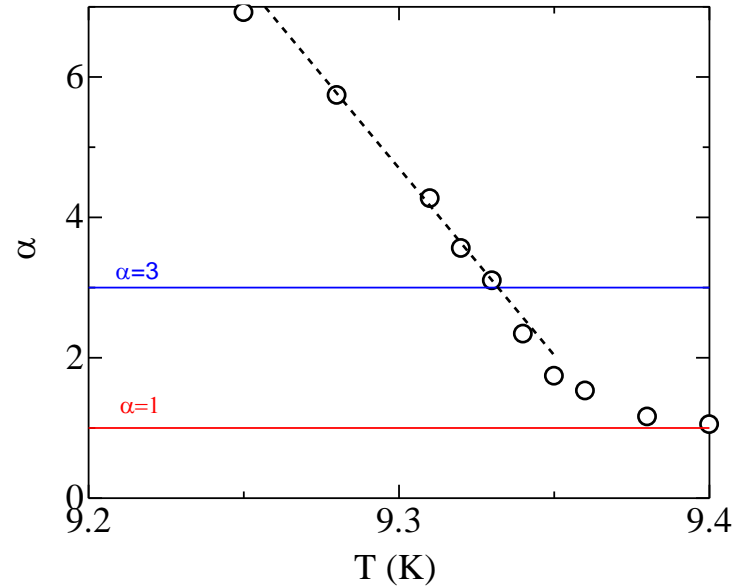
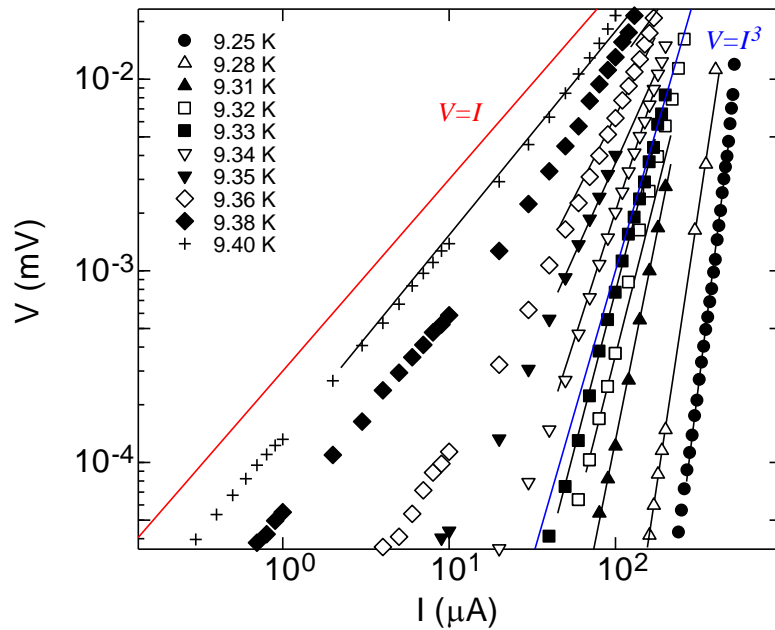
K. Makise et al. *Mater. Res. Express*
2 (2015) 106001.



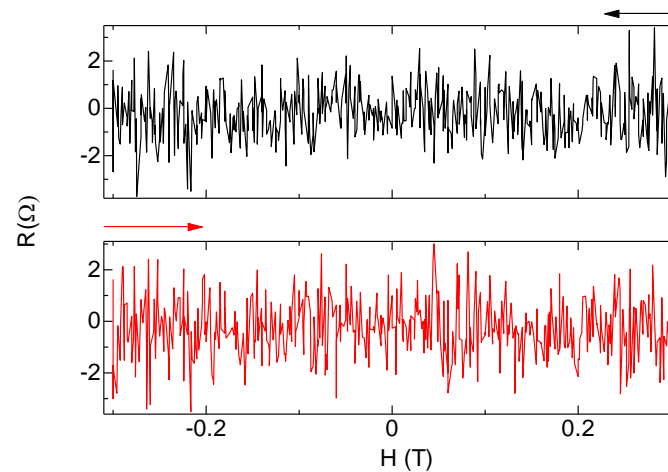
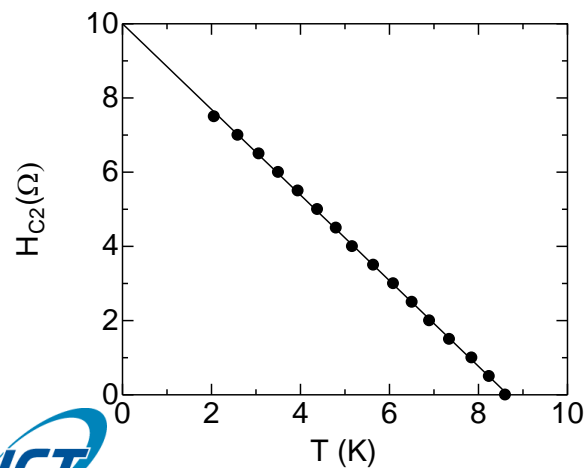
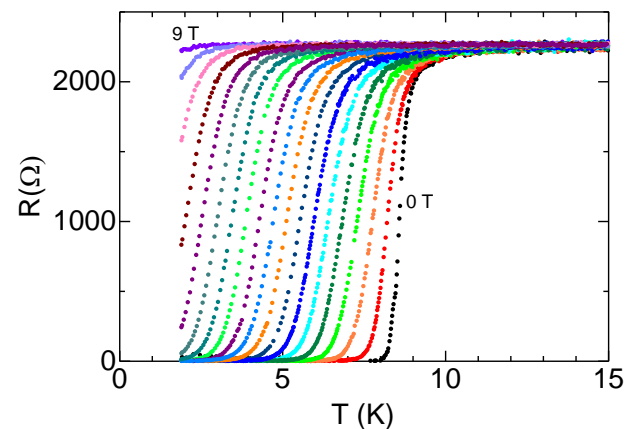
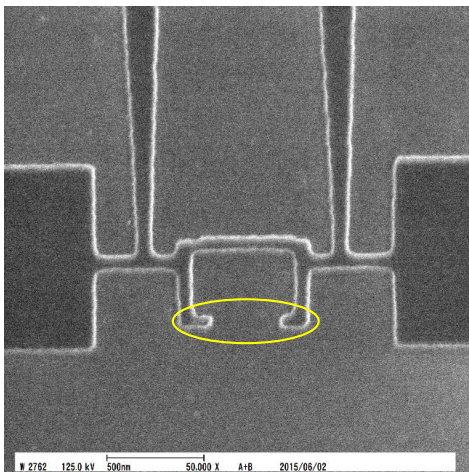
K. Makise et al. *J of phys. Cond. Matt.*
(2016) Accepted.



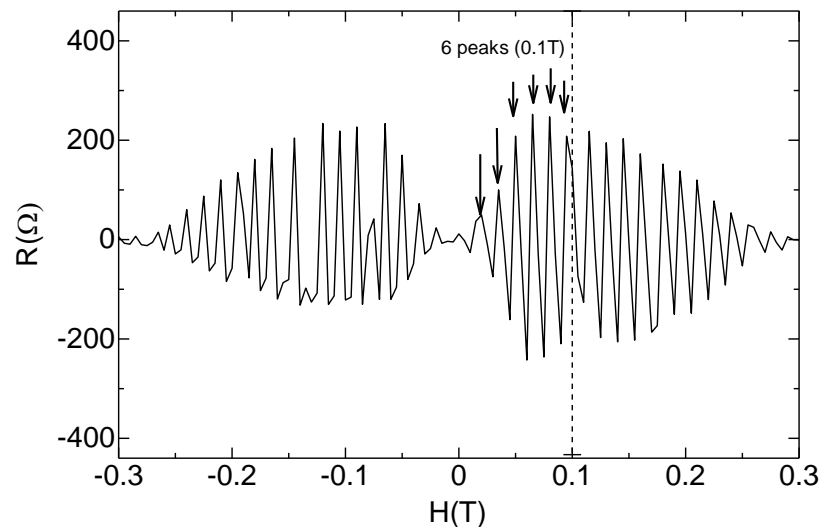
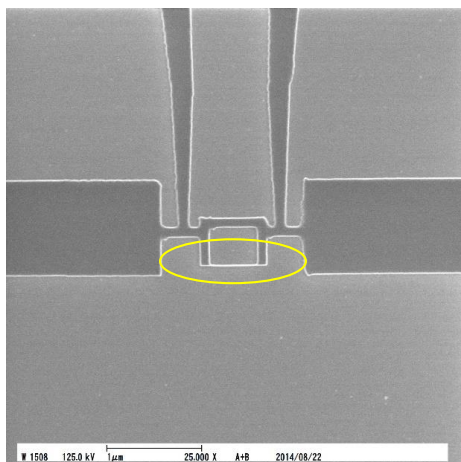
BKT transition (NbTiN films)



Phase slips ring in Mantic field



Phase slips ring in Mantic field



Area of Ring

$$S = 500 \times 250 \text{ nm}^2$$

$$a_{\square} = (\Phi_0/B)^{1/2} = (2.07 \times 10^{-15}/0.1)^{1/2} \\ = 143 \text{ nm}$$

$$Vn = S / (a_{\square} \times a_{\square}) = 6.11 \sim 6$$

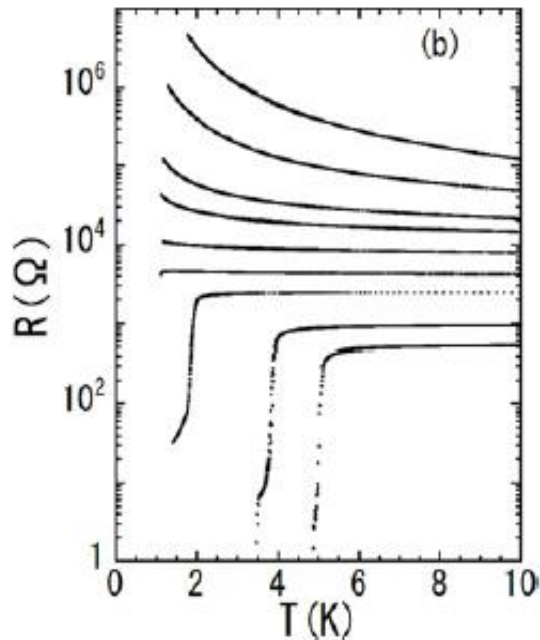
→QPS Qubits



MoRu and MoRuN films

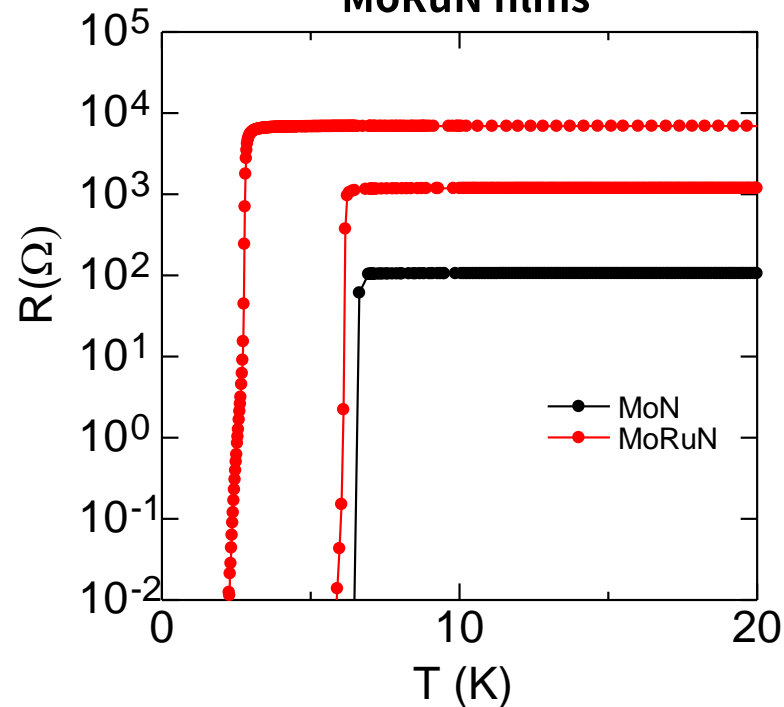


MoRu films



A. Hirakawa, and K. Makise
J. phys. Cond. Matt. 20
(2008)

MoRuN films

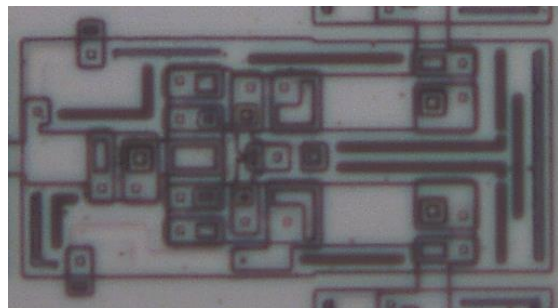


- High degree of disorder
- homogeneous system (amorphous)
- Nitride



Integrated monolithic QPS devices

(All Nitride SC)

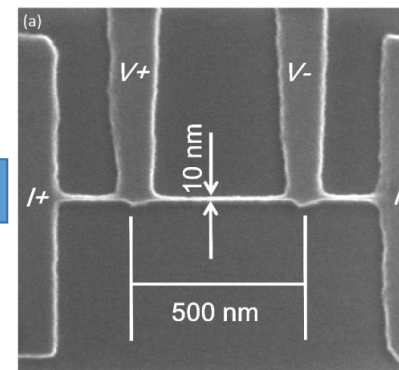


Read out
NbN-SFQ



CPW resonator **TiN**
MW of Single Photon level
Internal Q above 10^6

QPS **NbN, NbTiN, TiN, MoRuN**
→CQPS



Summary



- **We investigated the transport properties of superconducting NbTiN SNW's in a wide range of RN/L using four-probe method to eliminate the contact resistance.**
- **The characteristic with resistive tail below T_c for SNWs with high values of R_N can be well explained by the QPS theory.**
- **The analysis based on the model for the SNW which is being dual element upto Josephson junction, suggests that the separation of the superconducting and insulator phases may be controlled by the ratio of QPS amplitude energy E_s and inductive energy of SNW E_L .**
- **For the present NbTiN series, we observed that SIT may occur at $0.2 < E_s/E_L < 0.5$.**