

# Tu-I-SSP-04

## Ferromagnetic Josephson Junctions for Cryogenic Memory

Norman Birge, Michigan State University

in collaboration with Northrop Grumman Corporation  
and Arizona State University

(with thanks to D. Scott Holmes, Booz Allen Hamilton & IARPA)

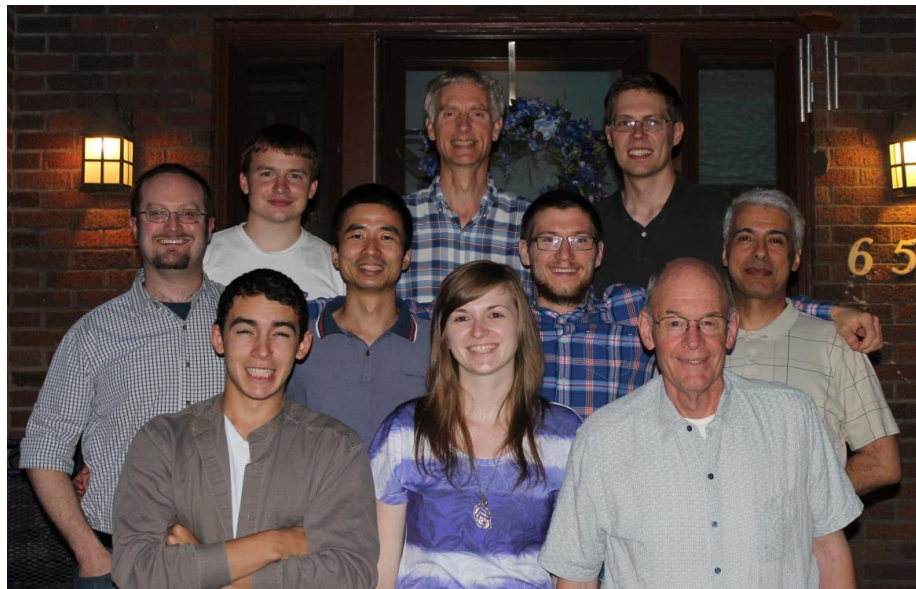
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# The group

**Eric Gingrich, Bethany Niedzielski, Joseph Glick, Yixing Wang, Bill Martinez, Josh Willard, Sam Edwards, Reza Loloee, William P. Pratt, Jr., plus many earlier students!**



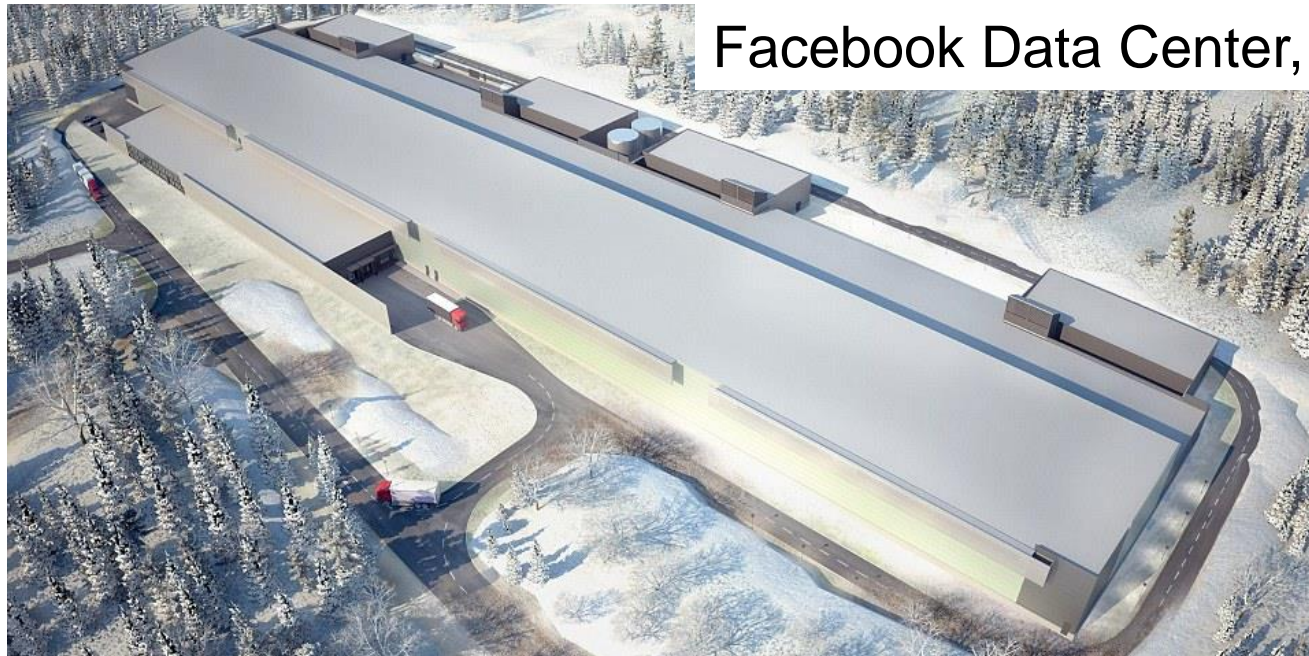
An old picture...

# Outline

- The need for energy-efficient computing
- Superconducting memory: JMRAM
- Superconducting/ferromagnetic hybrid systems
- Demonstration of phase control of an S/F/S Josephson junction – the basic memory device
- Future prospects

# ☐ The need for energy-efficient computing

## Facebook Data Center, Luleå, Sweden



Copyright  
Facebook

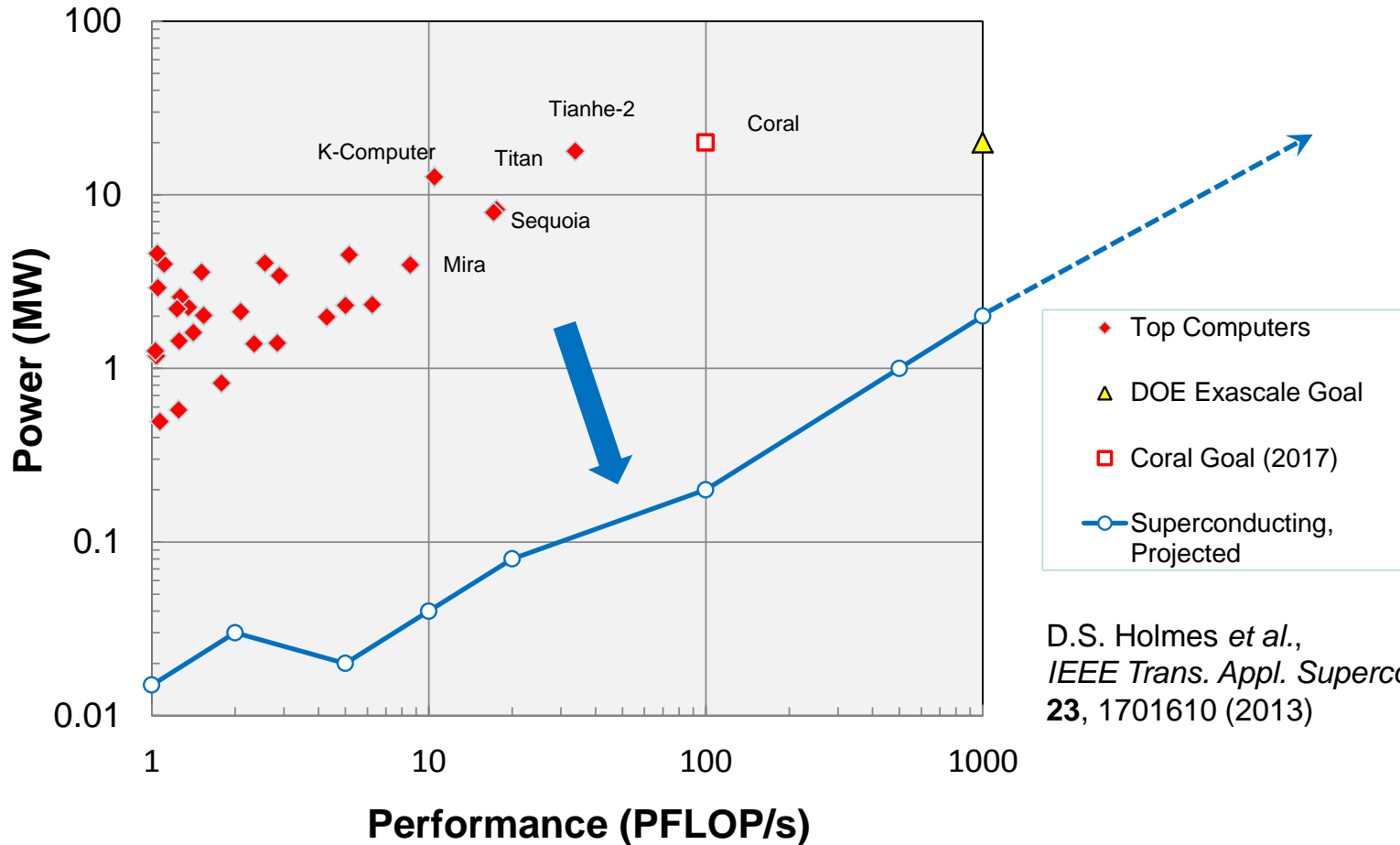
- Opened in 2013
- Cost: ~760 M\$
- Nearby Lule River generates 9% of Sweden's electricity (~4.23 GW)
- Average annual temperature: 1.3 °C

	Specifications
Performance*	27-51 PFLOP/s
Memory*	21-27 PB RAM 1900-6800 PB disk
Power	84 MW avg* (120 MW max)
Space	290,000 ft <sup>2</sup> (27,000 m <sup>2</sup> )
Cooling*	~1.07 PUE

\* estimated

Slide courtesy of Scott Holmes

# Superconducting computing looks promising



D.S. Holmes *et al.*,  
*IEEE Trans. Appl. Supercond.*  
**23**, 1701610 (2013)

# Our approach to superconducting memory:

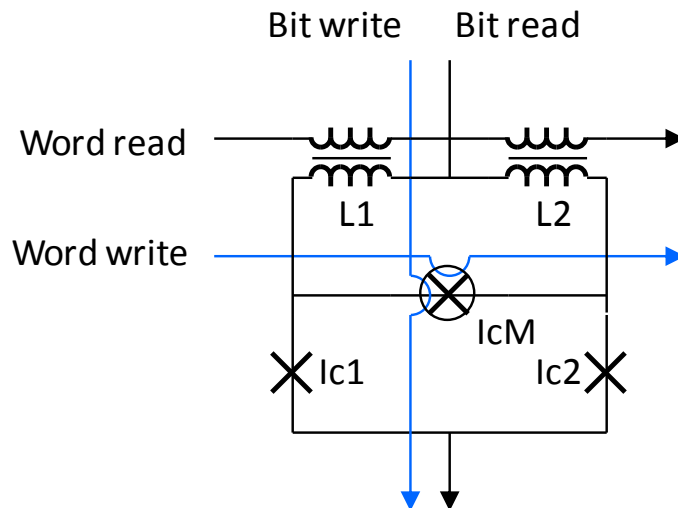


## Josephson Magnetic Random Access Memory (JM RAM)

Anna Y. Herr & Quentin P. Herr, US Patent 8,270,209 (2012)

A.Y. Herr, Q.P. Herr & Ofer Naaman, US Patent 9,208,861 (2015)

Northrop Grumman Corporation



Memory cell is a SQUID loop

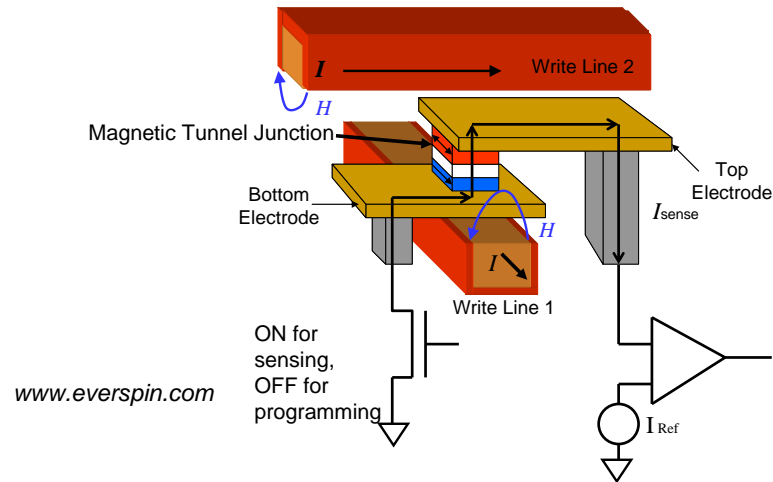
One junction has two stable phase states for "0" and "1"

Magnetic states are written using standard MRAM techniques

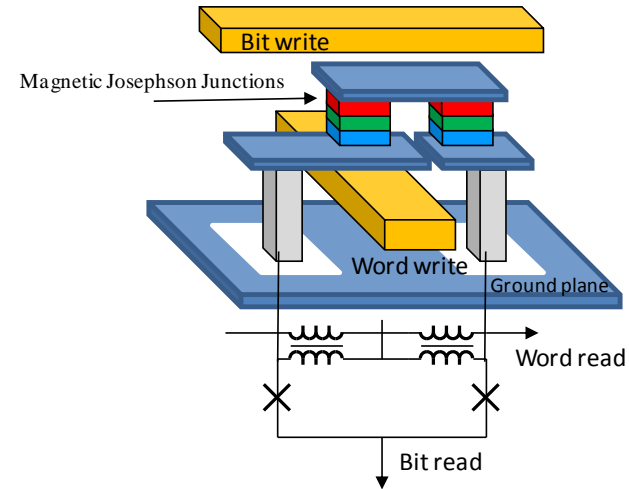
# JMRAM is a Superconducting MRAM



## MRAM (Everspin)



## JMRAM (Northrop Grumman)



**Memory cell is a Josephson junction containing a magnetic spin valve**

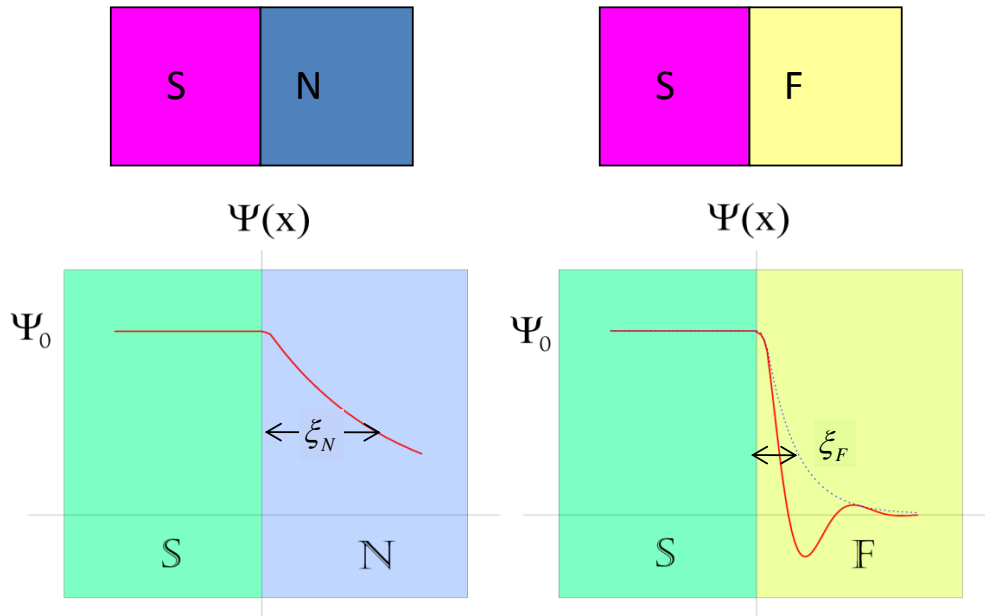
memory state – spin-valve state sets junction phase

write – magnetization reversal

read – Josephson effect

No idle/static power dissipation, read energy is dissipated only for logical “1”

# Superconductor/Ferromagnet proximity effect

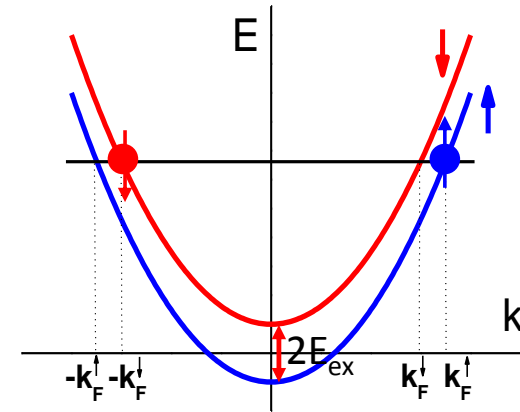


$$\xi_N = \sqrt{\frac{\hbar D_N}{2\pi k_B T}} \approx 1 \mu m$$

$$\xi_F \sim \sqrt{\frac{\hbar D_F}{E_{ex}}} \approx 1 nm$$

$D$  = diffusion constant

“FFLO-type” physics



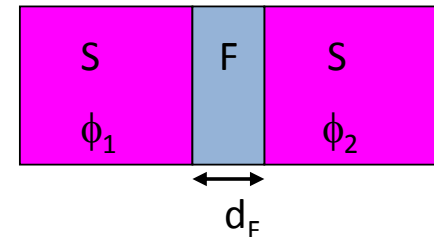
$$k_F^{\uparrow} - k_F^{\downarrow} \equiv Q \approx 2E_{ex} / \hbar v_F$$

$$\xi_F = Q^{-1} = \frac{\hbar v_F}{2E_{ex}} \quad \text{ballistic}$$

$$\xi_F = \sqrt{\frac{\hbar D_F}{E_{ex}}} \quad \text{diffusive}$$



# Consequence: S/F/S Josephson junctions oscillate between 0 and $\pi$ junctions as $d_F$ increases:



Buzdin, Bulaevskii, & Panyukov (1982).

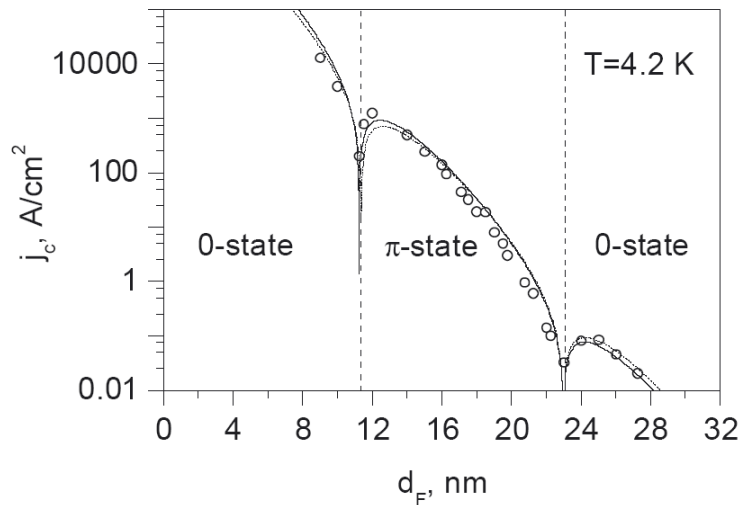


$$\xi_F = \sqrt{\frac{\hbar D_F}{E_{ex}}}$$

0-state:  $I_s = I_c \sin(\phi_2 - \phi_1)$

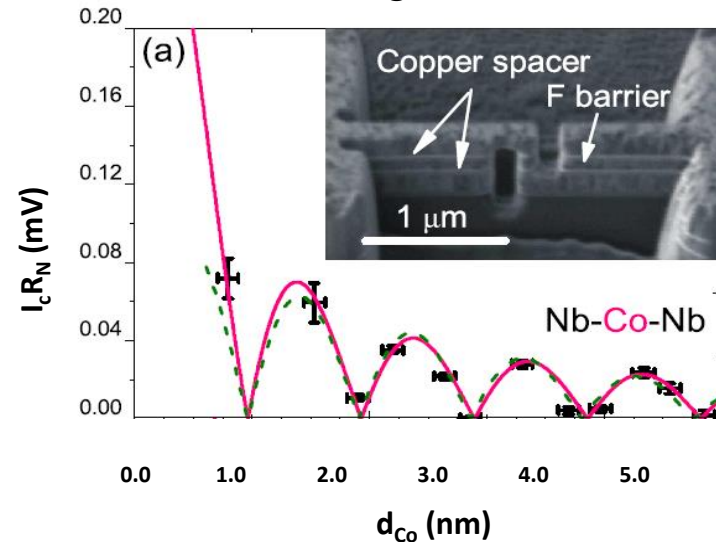
$\pi$ -state:  $I_s = I_c \sin(\phi_2 - \phi_1 + \pi)$

Weak F:  $\text{Cu}_{48}\text{Ni}_{52}$  alloy



Ryazanov *et al.*, PRL **86**, 2427 (2001);  
 PRL **96**, 197003 (2006).

Strong F: Co



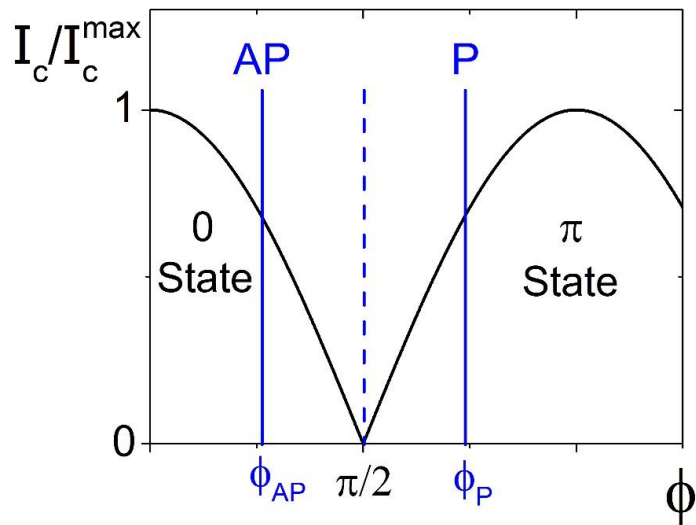
Robinson, Piano, Burnell, Bell, Blamire,  
 PRL **97**, 177003 (2005)

Can we control  $I_c$  or phase state of a single Josephson junction?

## ☞ Add a second ferromagnetic layer: S/F<sub>1</sub>/F<sub>2</sub>/S

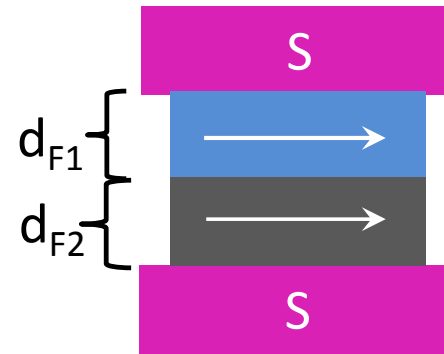
I<sub>c</sub> and phase depend on relative magnetization direction

Electron pair accumulates phase  $\phi$  while traversing junction

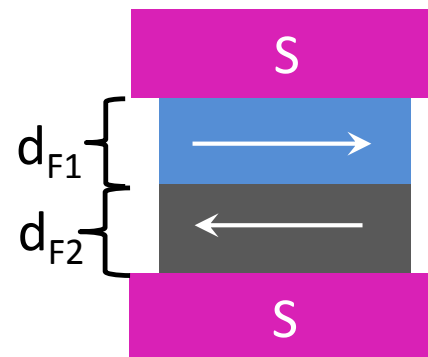


$$I_c \propto \exp[-(d_{F1}/\xi_{F1} + d_{F2}/\xi_{F2})] \left| \cos(d_{F1}/\xi_{F1} \pm d_{F2}/\xi_{F2}) \right|$$

**Parallel state:**  $\phi = \frac{d_{F1}}{\xi_{F1}} + \frac{d_{F2}}{\xi_{F2}}$



**Antiparallel state:**  $\phi = \frac{d_{F1}}{\xi_{F1}} - \frac{d_{F2}}{\xi_{F2}}$



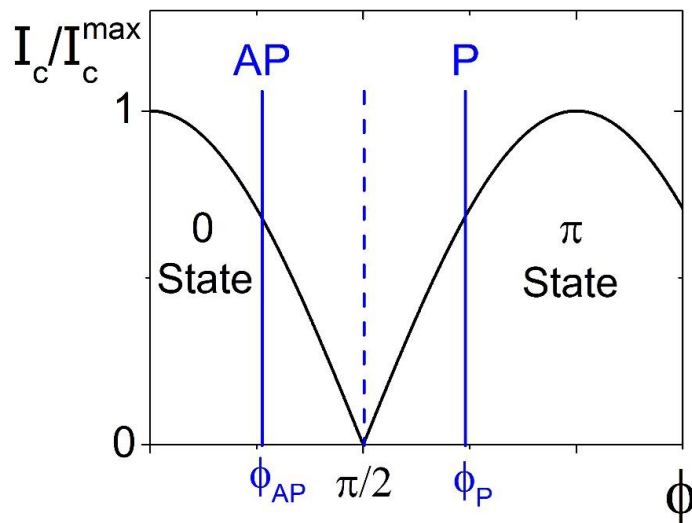
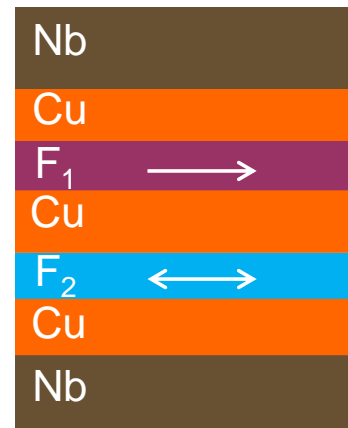
# S/F<sub>1</sub>/N/F<sub>2</sub>/S Josephson Junction Composition

Nb(100)/Cu(5)/NiFe(1.5)/Cu(10)/Ni(1.2)/Cu(5)/Nb(150)



F<sub>1</sub> = Ni (1.2 nm)  
fixed layer

F<sub>2</sub> = NiFe (1.5 nm)  
free layer



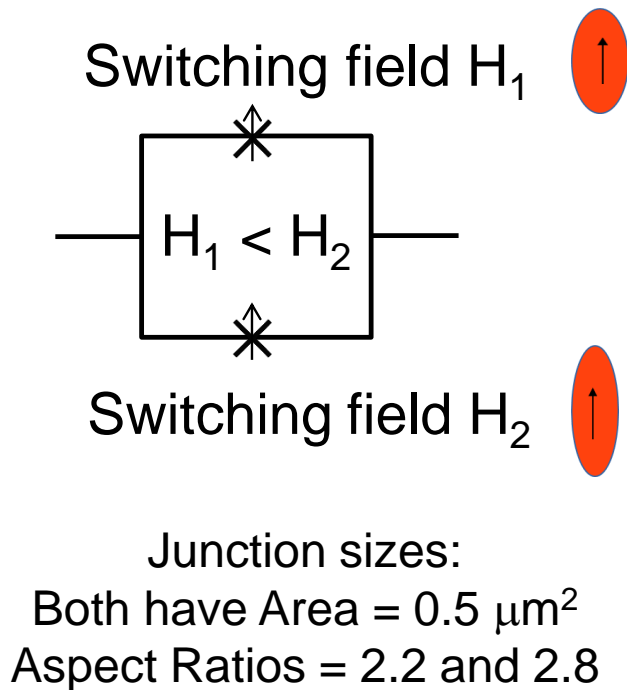
Choose NiFe thickness to put F<sub>2</sub> close to 0 -  $\pi$  transition.

Ni provides additional small push to right or left.



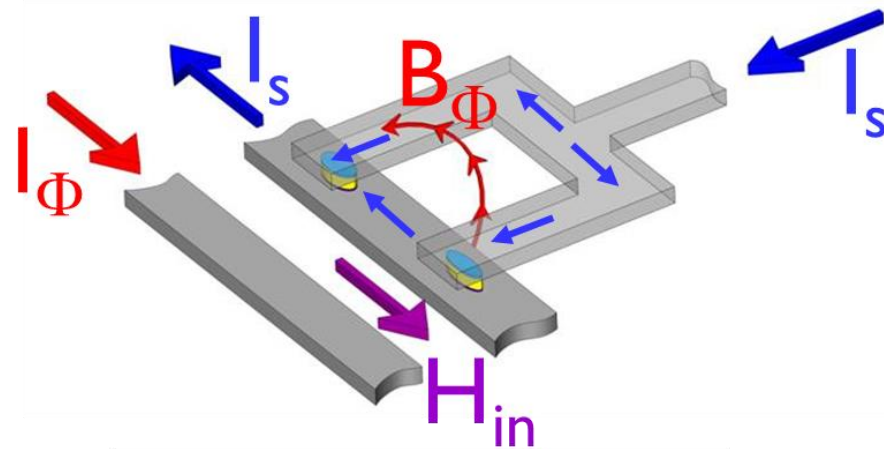
# Two junctions in a SQUID loop used to measure relative junction phase

**Schematic:**



**Cartoon of Actual Device:**

On-chip current line couples magnetic flux into SQUIDs

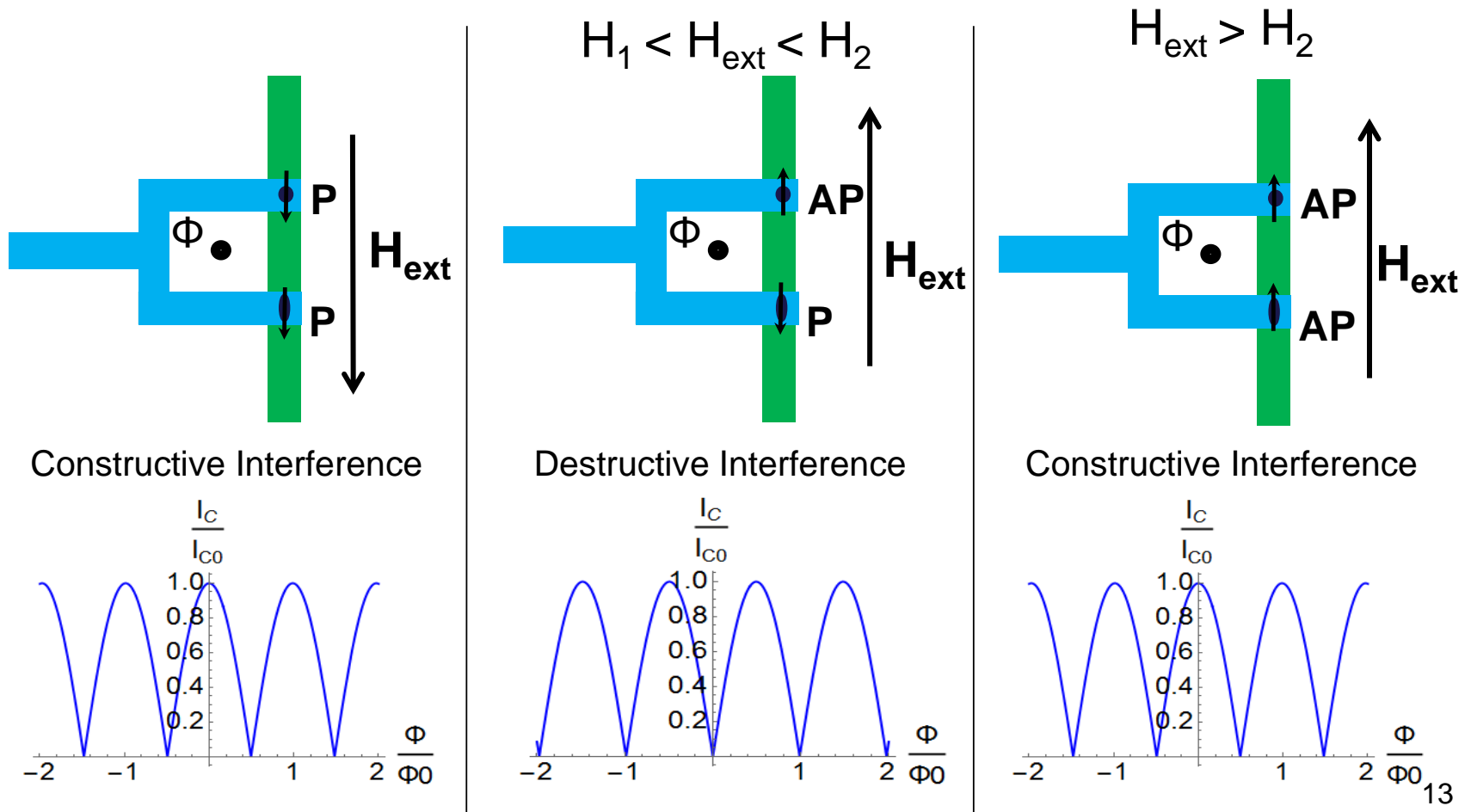


Switch magnetization with in-plane field

Different aspect ratio pillars have different switching fields

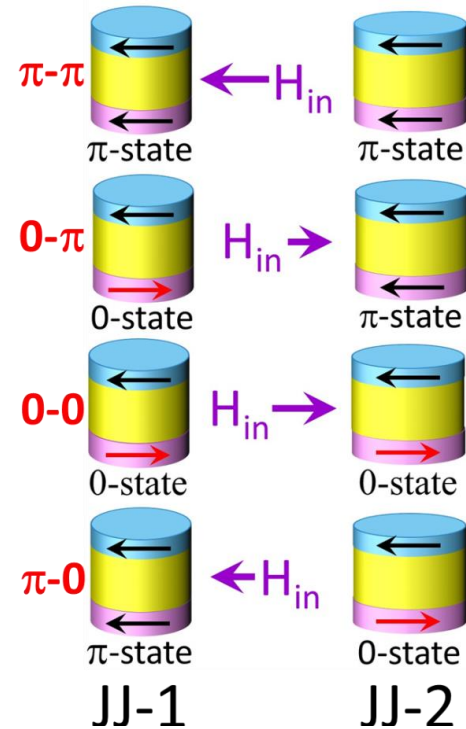
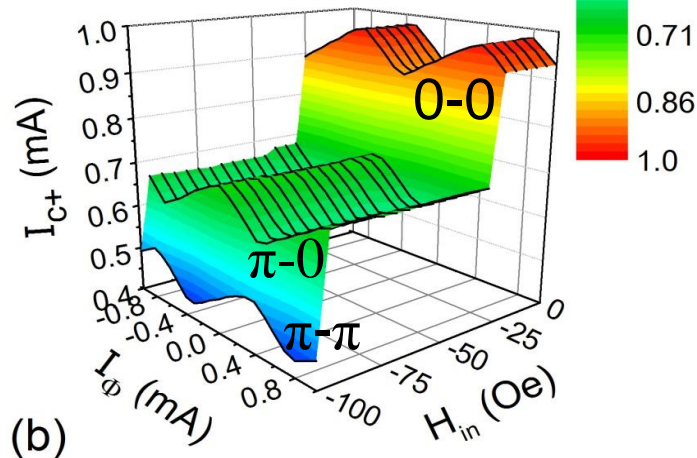
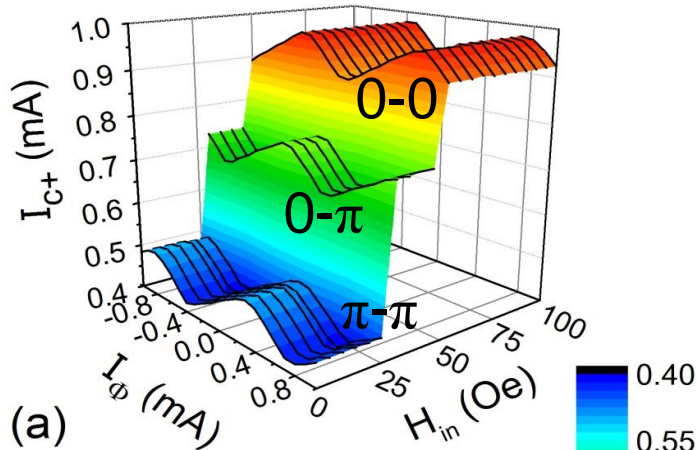
# At Relatively Low External Fields, Two Phase Changes Should Be Observable

Initialize with large field in  $-z$  direction, then slowly increase  $H_{\text{ext}}$  in  $+z$  direction



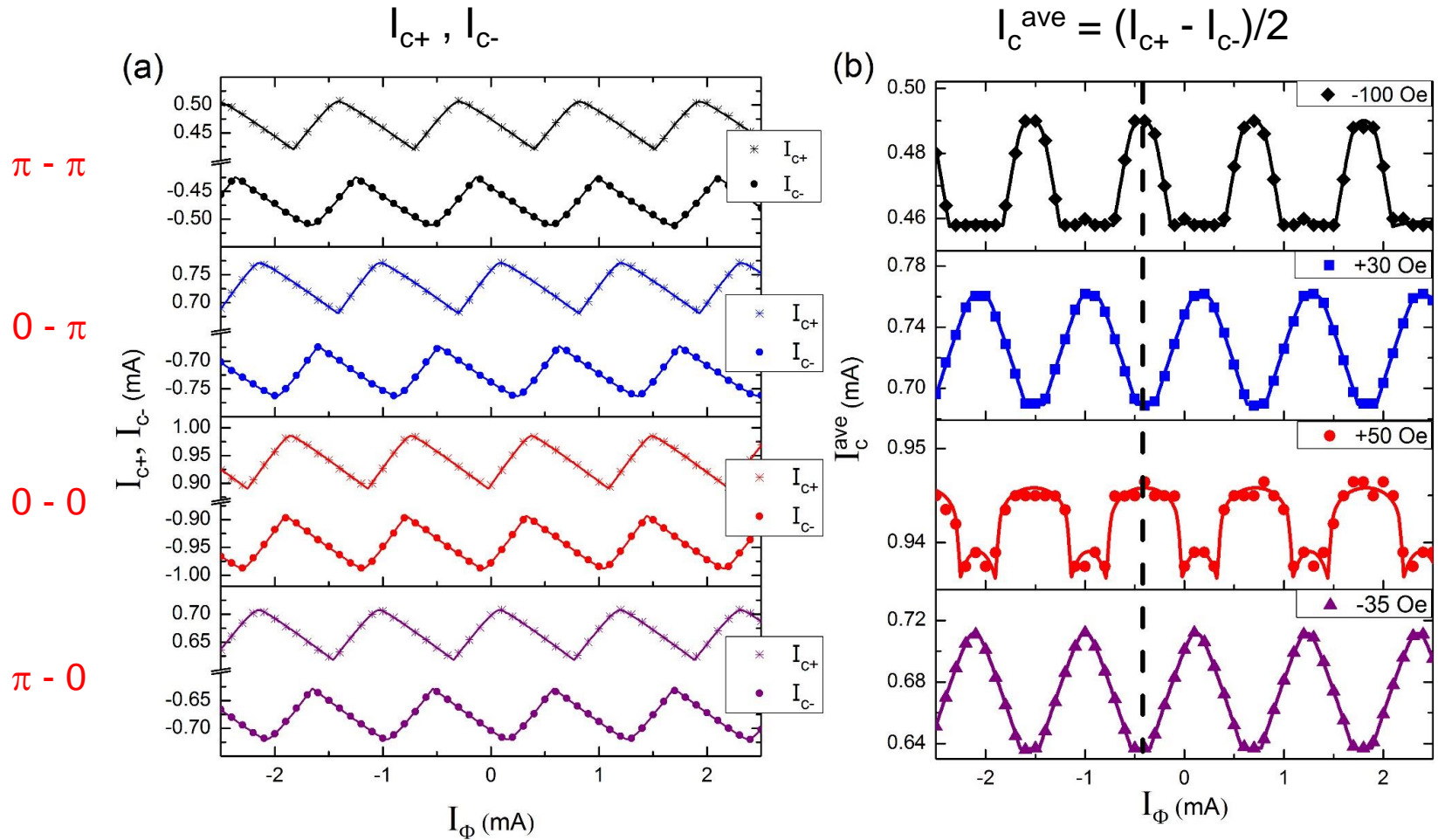


# Data show clean switching between the four expected states

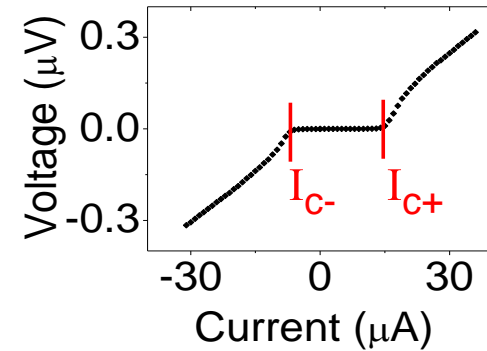
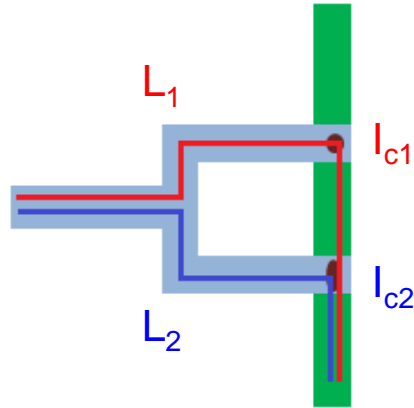
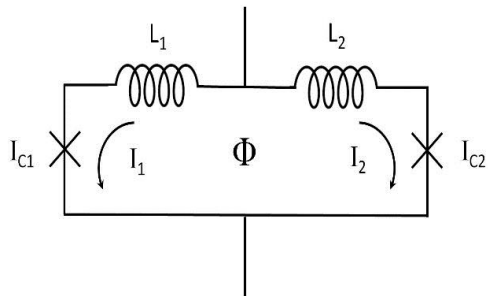


Switching Fields: +30 Oe, +50 Oe, -35 Oe, -100 Oe

# Data cuts for the four magnetic states



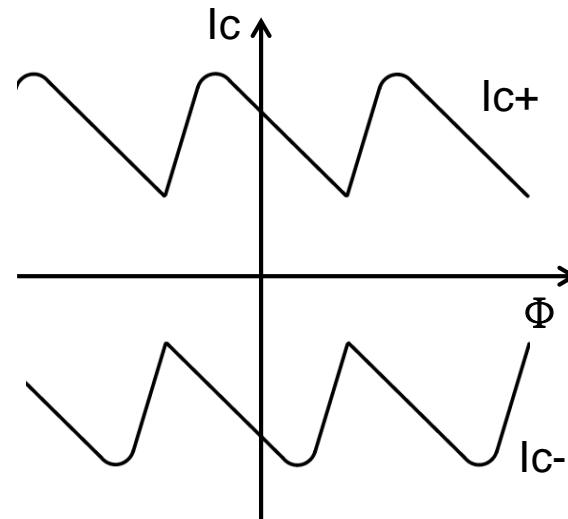
# $I_c(\Phi)$ curves have tilted ratchet shape when loop inductances and/or critical currents are asymmetric



$I_c^+(\Phi)$  and  $I_c^-(\Phi)$  oscillations are asymmetric when  $L_1 \neq L_2$  &  $I_{c1} \neq I_{c2}$

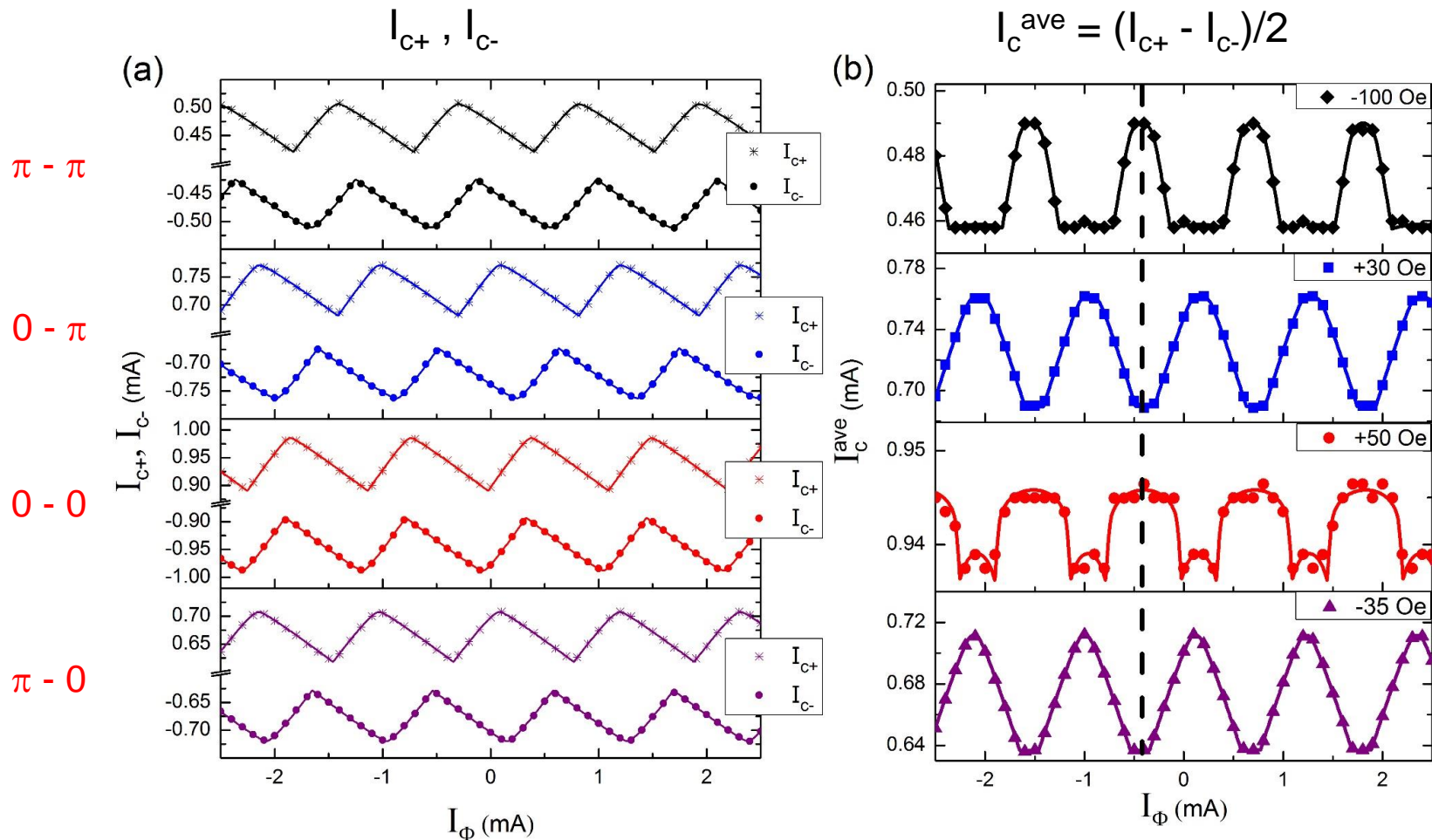
$I_c^+$  and  $I_c^-$  shift by equal amounts and in opposite directions along the  $\Phi$  axis

Analyze  $I_c^+$  and  $I_c^-$  peak shifts to extract JJ phase shifts



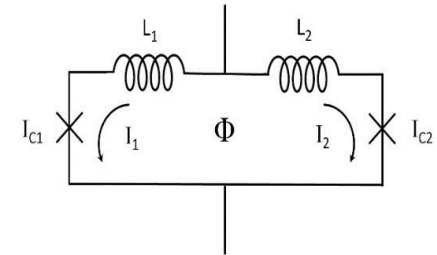


# Quantitative fits to SQUID modulation data for the four magnetic states



## Quantitative Analysis Consistently Assigns the Inductance and Critical Currents of Each State

state	$I_{c1}$ (mA)	$I_{c2}$ (mA)	$L_1$ (pH)	$L_2$ (pH)
$\pi - \pi$	0.292	0.217	5.73	11.38
$0 - \pi$	0.565	0.203	5.64	11.33
$0 - 0$	0.567	0.419	5.63	11.55
$\pi - 0$	0.294	0.420	5.71	11.56
		ave	5.68	11.46
		$\sigma$	0.05	0.12



FastHenry simulations:

$L_1 \approx 7$  pH,  $L_2 \approx 13$  pH

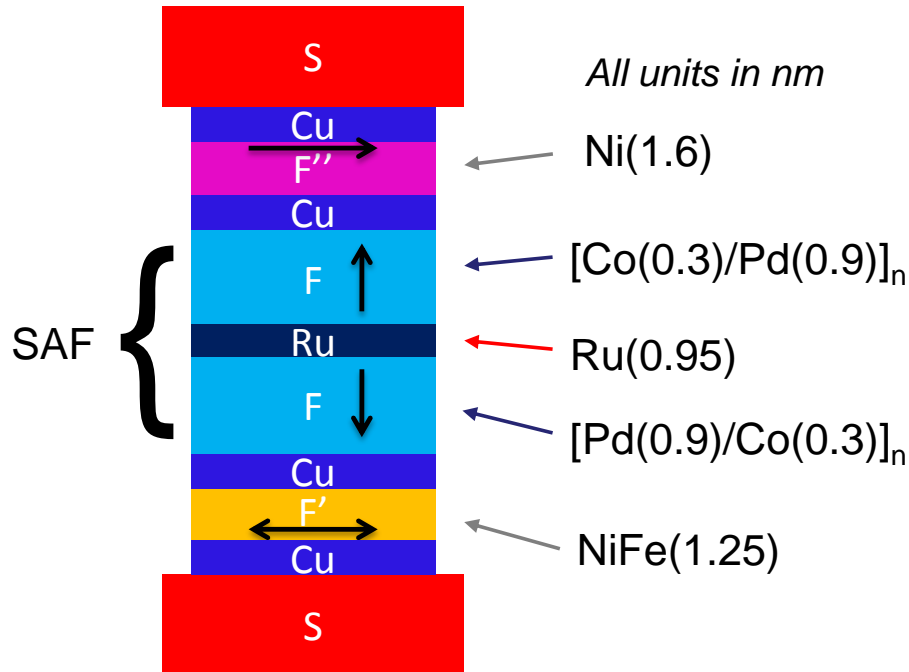
Fitting parameters from independent fits of 4 magnetic states are highly consistent

- Exception: critical current of JJ #2 changes slightly in  $\pi$  state when JJ #1 switches from  $\pi$  to 0 state



## New result!

# Controllable $0-\pi$ switching with spin-triplet supercurrent



Spin-triplet supercurrent decays very slowly in F

$0 - \pi$  switching occurs by spin rotation rather than accumulated pair phase

Spin-triplet JJ requires three F layers with non-collinear magnetizations between adjacent layers

Data are not yet available for public dissemination, but we plan to submit them for publication soon: J.A. Glick et al. (2017)



# What needs to be done

- Memory (see talk Fr-C-DIG-03 by Ofer Naaman)
  - Optimize performance of magnetic materials
    - Lower  $M_{\text{sat}} \Rightarrow$  lower  $E_{\text{switch}}$
    - Reduce extrinsic sources of anisotropy in thin films
    - Find better material for fixed layer (Ni has issues)
    - Minimize underlayer roughness
  - Develop read/write electronics & interface to SFQ logic
    - (see poster We-SDM-08 by Quentin Herr)
- Make the rest of the computer! (see talk Fr-I-DIG-02 by Anna Herr)

# Conclusions

- Magnetic Josephson junctions have demonstrated potential for ultra-low-power cryogenic memory
- Much more work needs to be done!

# Bibliography

- E.C. Gingrich, B.M. Niedzielski , J.A. Glick , Y Wang , D.L. Miller , R. Loloee R, W.P. Pratt Jr, N.O. Birge, “**Controllable  $0-\pi$  Josephson junctions containing a ferromagnetic spin valve**,” *Nature Phys.* **12**, 564 (2016). DOI: [10.1038/nphys3681](https://doi.org/10.1038/nphys3681)
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- D.S. Holmes, A.L. Ripple, and M.A. Manheimer, “**Energy-efficient superconducting computing – power budgets and requirements**”, *IEEE Trans. Appl. Supercond.*, **23**, 1701610 (2013). DOI: [10.1109/TASC.2013.2244634](https://doi.org/10.1109/TASC.2013.2244634)
- Q.P. Herr, A.Y. Herr, O.T. Oberg, and A.G. Ioannidis, “**Ultra-low-power superconductor logic**”, *J. Appl. Phys.* **109**, 103903 (2011). DOI: [10.1063/1.3585849](https://doi.org/10.1063/1.3585849)