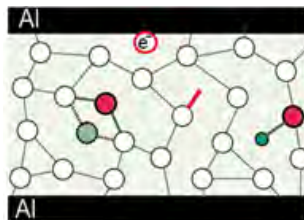
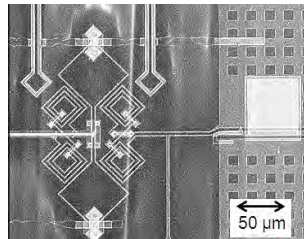
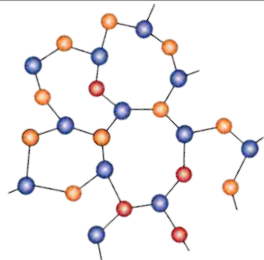


Spectroscopy and coherent Control of Defects in superconducting Films and Qubits

Jürgen Lisenfeld, Alexander Bilmes, Jan Brehm, Georg Weiss, and A.V. Ustinov

Physikalisches Institut, Karlsruhe Institute of Technology, Karlsruhe, Germany

ju.lis@kit.edu

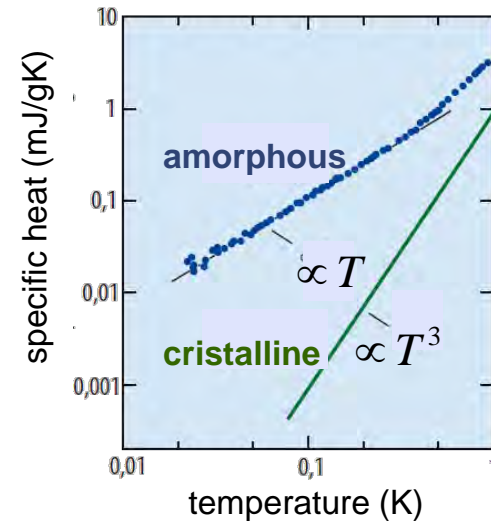


- **Two-Level-Systems (TLS):**
a major source of noise in quantum devices
- **Using superconducting Quantum Bits to study single TLS**
 - TLS spectroscopy and mechanical strain tuning
 - mutual TLS interactions, noise spectroscopy
 - TLS - electron interaction

Two-Level-Systems (TLS)

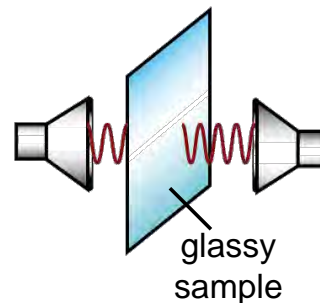
measurements on glasses revealed several peculiarities:

- specific heat disagrees with Debye model (Zeller & Pohl 1971)



■ common signatures in amorphous materials

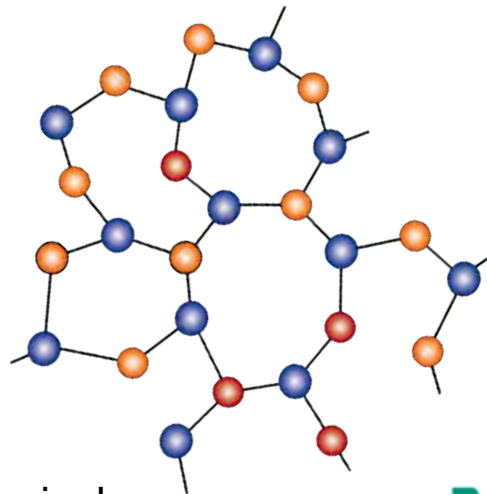
- ultrasound attenuation
- electric field response
- microwave response
- thermal conductivity
- heat capacity ...



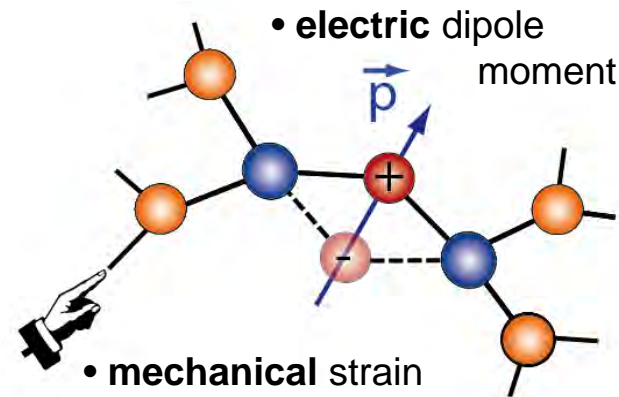
➔ conclusion: glasses must contain **intrinsic states** having excitation energies < 1 K and which **couple to both electric fields and phonons**.

Two-Level-Systems: Tunneling Atom Model^[1,2]

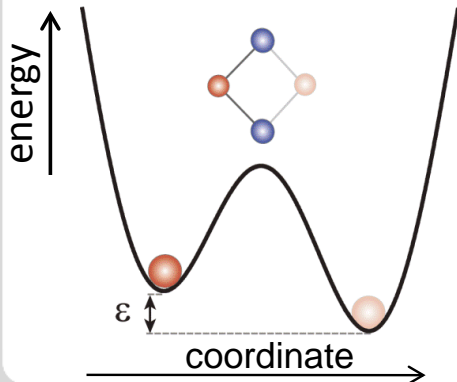
■ in amorphous materials, atoms may tunnel between two positions:



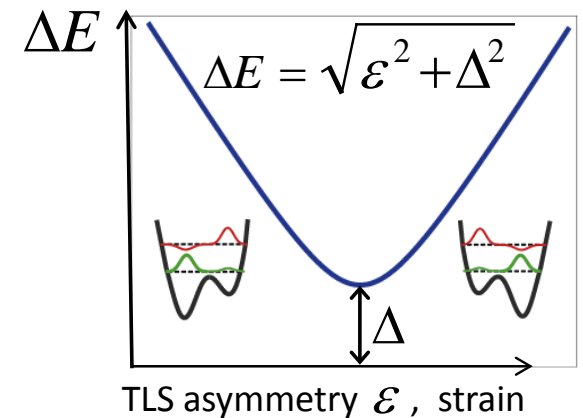
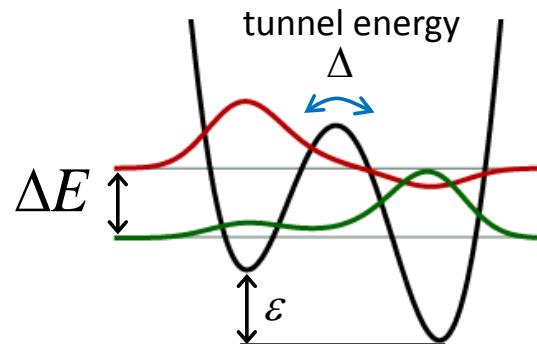
■ these "tunneling systems" couple via



■ classical



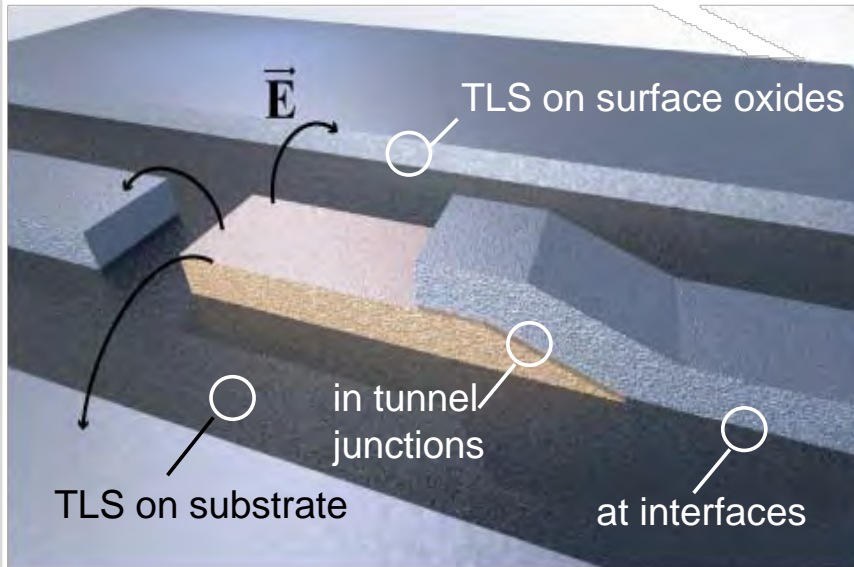
■ quantum



[1] W.A. Phillips, *J. Low Temp. Phys.* **7** 351 (1972)

[2] Anderson, Halperin, and Varma, *Philosophical Mag.* **25**, 1 (1972)

TLS in microfabricated circuits and Josephson junctions



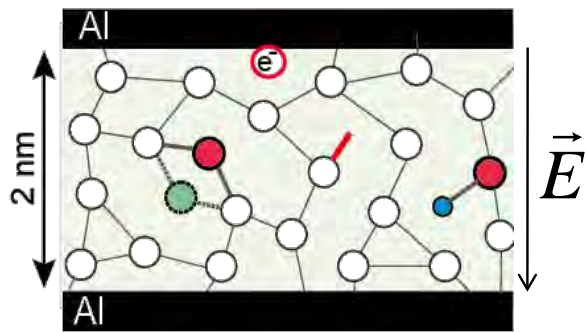
TLS are found

- in surface oxides
- in / on the substrate
- at interfaces
- in tunnel junctions

TLS generate noise & dissipation in

- MOSFETs & single-electron transistors
- micro-mechanical resonators
- single-photon detectors, nanowires
- superconducting resonators and qubits
- ...

in Josephson junctions:

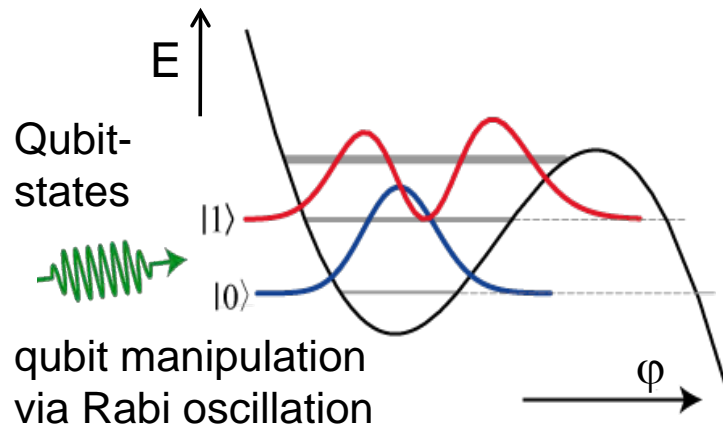
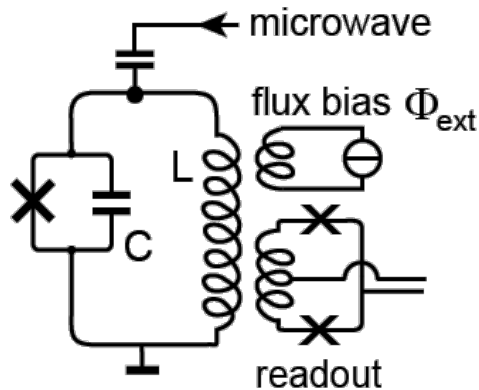


- **hydroxide defects** Appl. Phys. Lett. **97**, 252501 (2010)
- **dangling bonds**
- **electrons trapped at interfaces:** PRL **95**, 046805 (2005)
 Kondo- / Andreev Fluctuators PRB **84**, 235102 (2011)
- **phononically dressed electrons** PRB **87**, 144201 (2013)
- **tunneling atoms** Phys. Rev. Lett. **95**, 210503 (2005)

The Phase Qubit

J.M. Martinis et al., PRL 93, 077003 (2004)

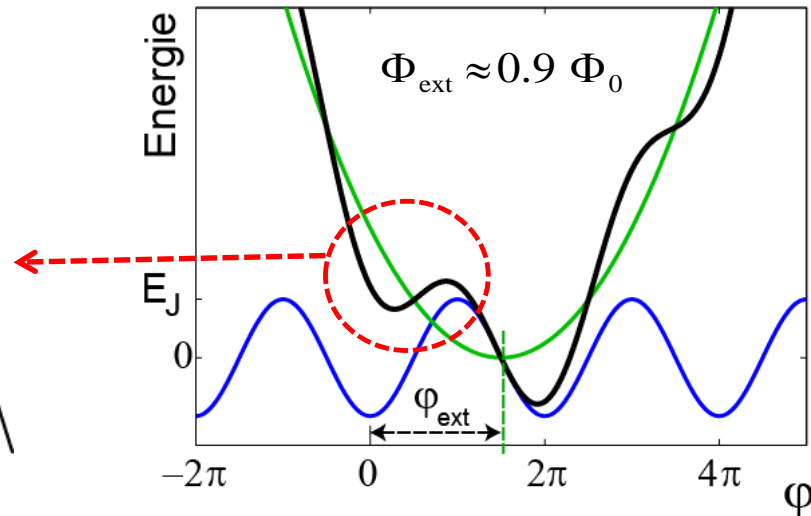
Complete circuit



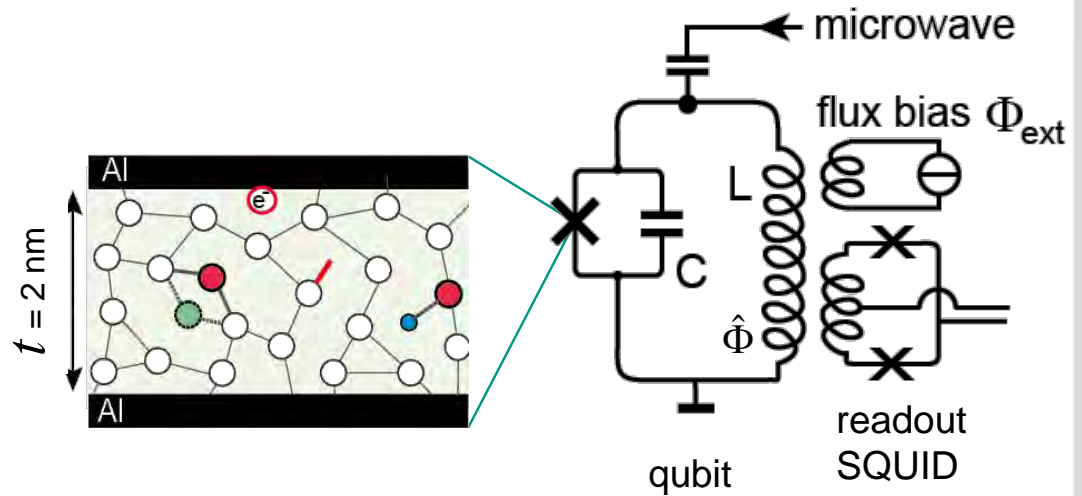
Hamilton-Operator

$$\hat{H} = \frac{\hat{Q}^2}{2C} + \frac{\hat{\Phi}^2}{2L} - E_J \cos \hat{\phi}$$

Potential for $E_J \gg Q^2/2C$:



Defect-Qubit - interaction



■ **Qubit-TLS interaction:**
 via TLS electrical dipole moment \vec{p}

Qubit:

$$\hat{H} = \frac{\hat{\Phi}^2}{2L} + \frac{\hat{Q}^2}{2C} + E_J$$

el. Field:

$$\vec{E} = \frac{\hat{Q}}{t C} \approx 1000 \text{ V/m}$$

Dipole interaction:

for $\bar{p} = 2 D = 2 \cdot 0.2 \text{ eÅ}$

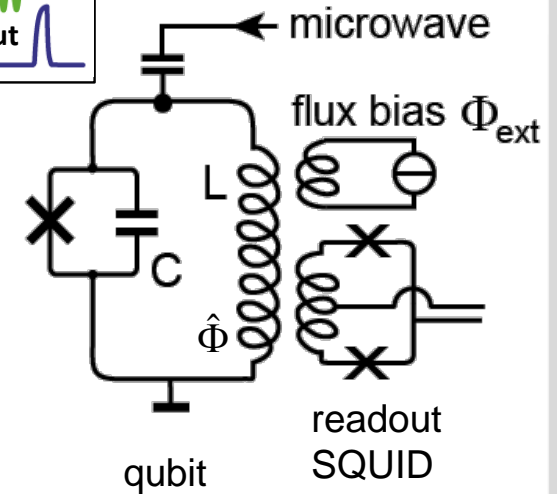
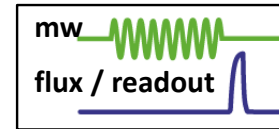
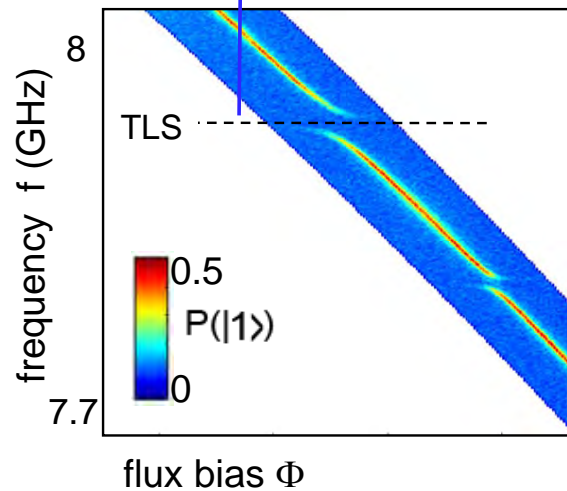
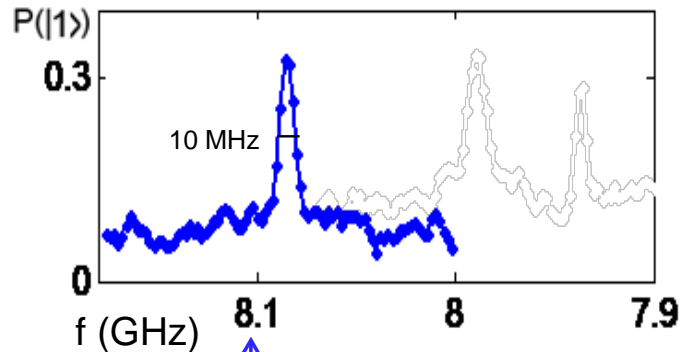
$$\Rightarrow g = \bar{p} |\vec{E}| \approx h \cdot 10 \text{ MHz}$$

qubit-TLS coupling strength

Defect – Qubit - Interaction

Frequency Domain:

defects cause avoided level crossings



■ Qubit-TLS interaction: via TLS electrical dipole moment \vec{p}

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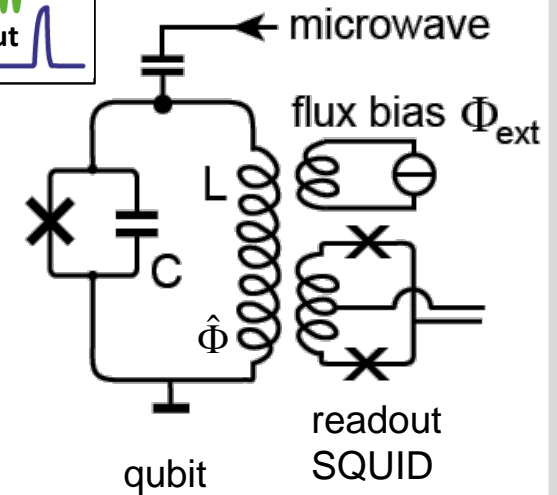
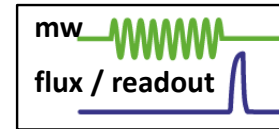
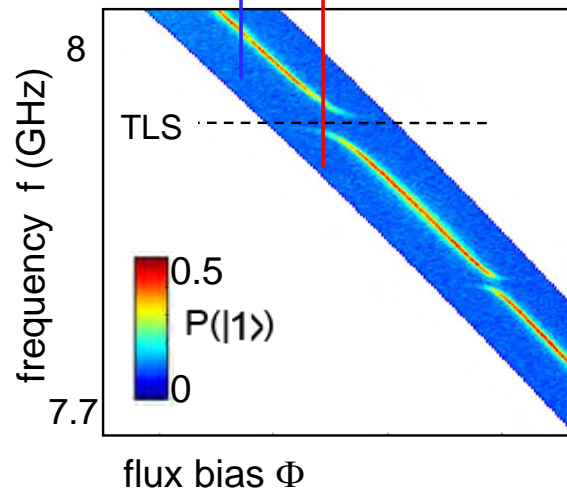
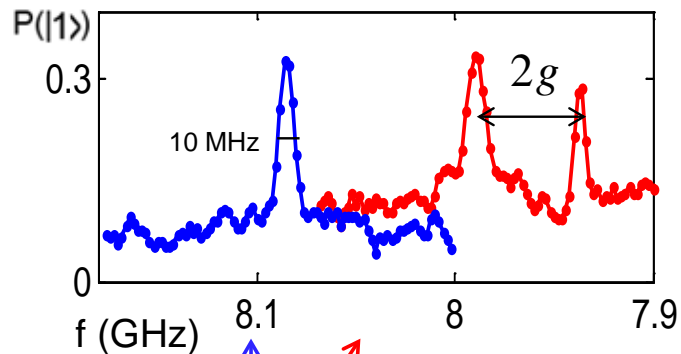
qubit-TLS coupling strength

cf.: R.W. Simmonds, K.M. Lang, D.A. Hite, D.P. Pappas, and J.M. Martinis, PRL **93**, 077003 (2004)
 Yoni Shalibo, Matthew Neeley, John M. Martinis, Nadav Katz et al., PRL **105**, 177001 (2010).

Defect – Qubit - Interaction

Frequency Domain:

defects cause avoided level crossings



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Dipole interaction:

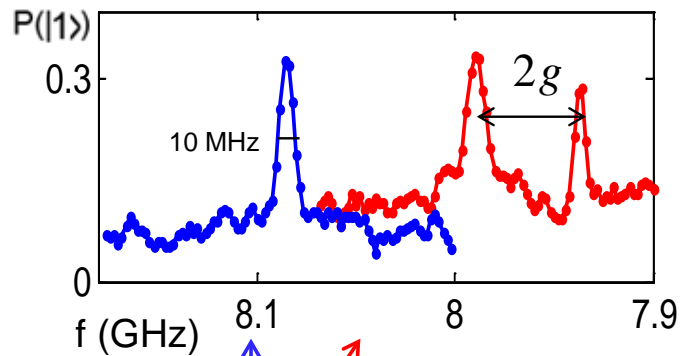
for $\vec{p} = 2 D = 2 \cdot 0.2 \text{ eÅ}$

$\Rightarrow g = \vec{p} \cdot \vec{E} \approx h \cdot 10 \text{ MHz}$
qubit-TLS coupling strength

Defect – Qubit - Interaction

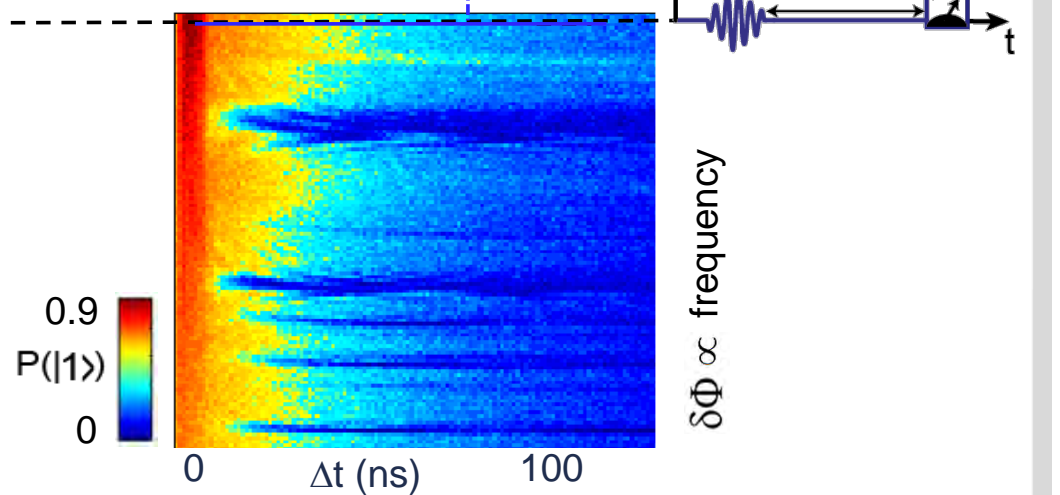
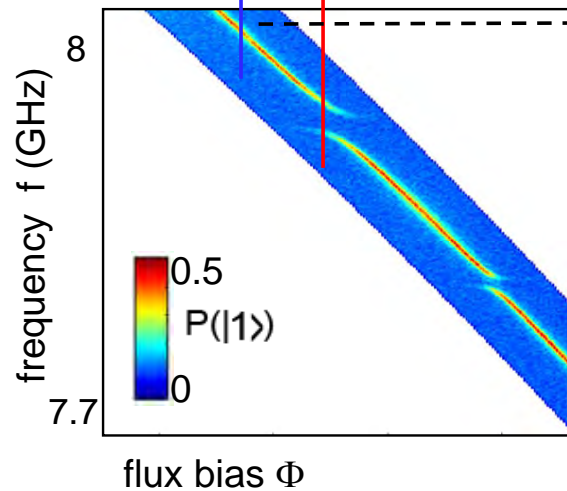
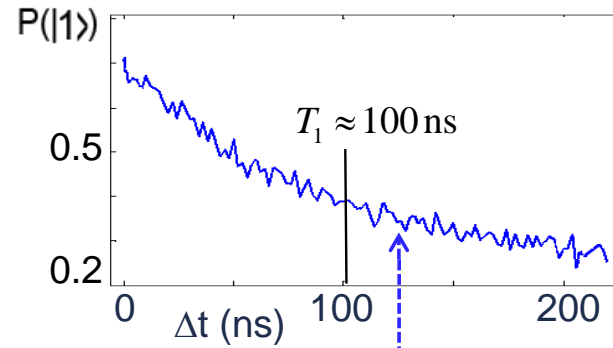
Frequency Domain:

defects cause avoided level crossings



Time Domain:

qubit decays due to energy relaxation

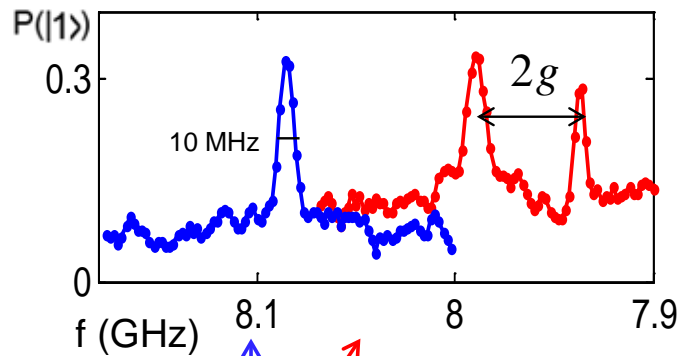


cf.: R.W. Simmonds, K.M. Lang, D.A. Hite, D.P. Pappas, and J.M. Martinis, PRL **93**, 077003 (2004)
 Yoni Shalibo, Matthew Neeley, John M. Martinis, Nadav Katz et al., PRL **105**, 177001 (2010).

Defect – Qubit - Interaction

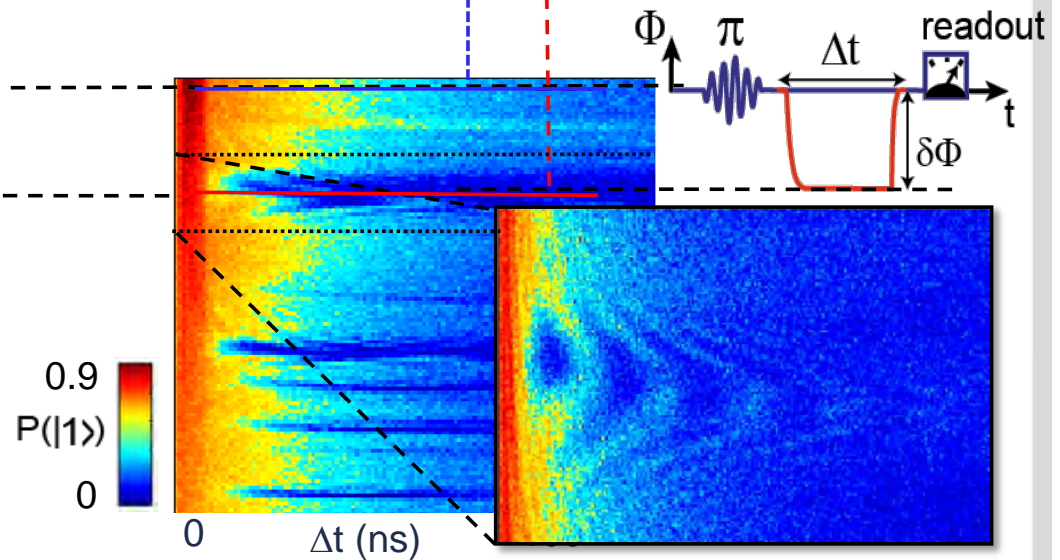
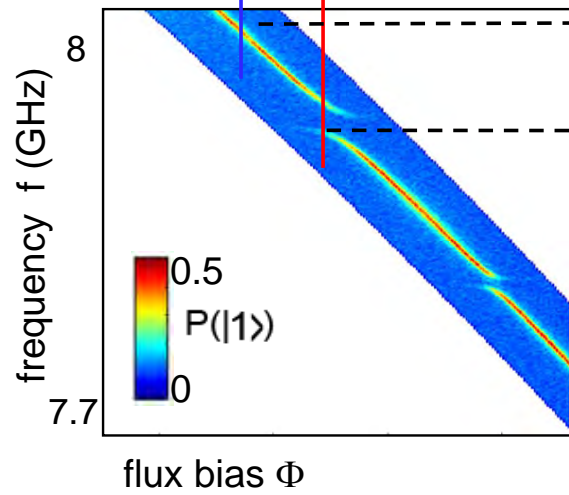
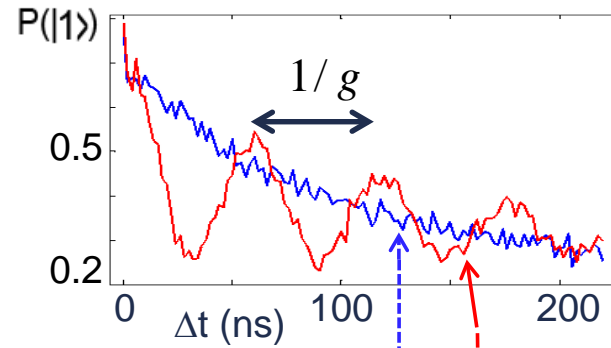
Frequency Domain:

defects cause avoided level crossings



Time Domain:

energy oscillates between qubit and defects

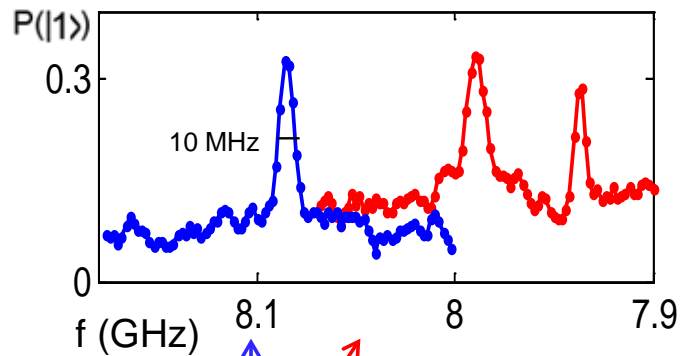


cf.: R.W. Simmonds, K.M. Lang, D.A. Hite, D.P. Pappas, and J.M. Martinis, PRL **93**, 077003 (2004)
 Yoni Shalibo, Matthew Neeley, John M. Martinis, Nadav Katz et al., PRL **105**, 177001 (2010).

Defect – Qubit - Interaction

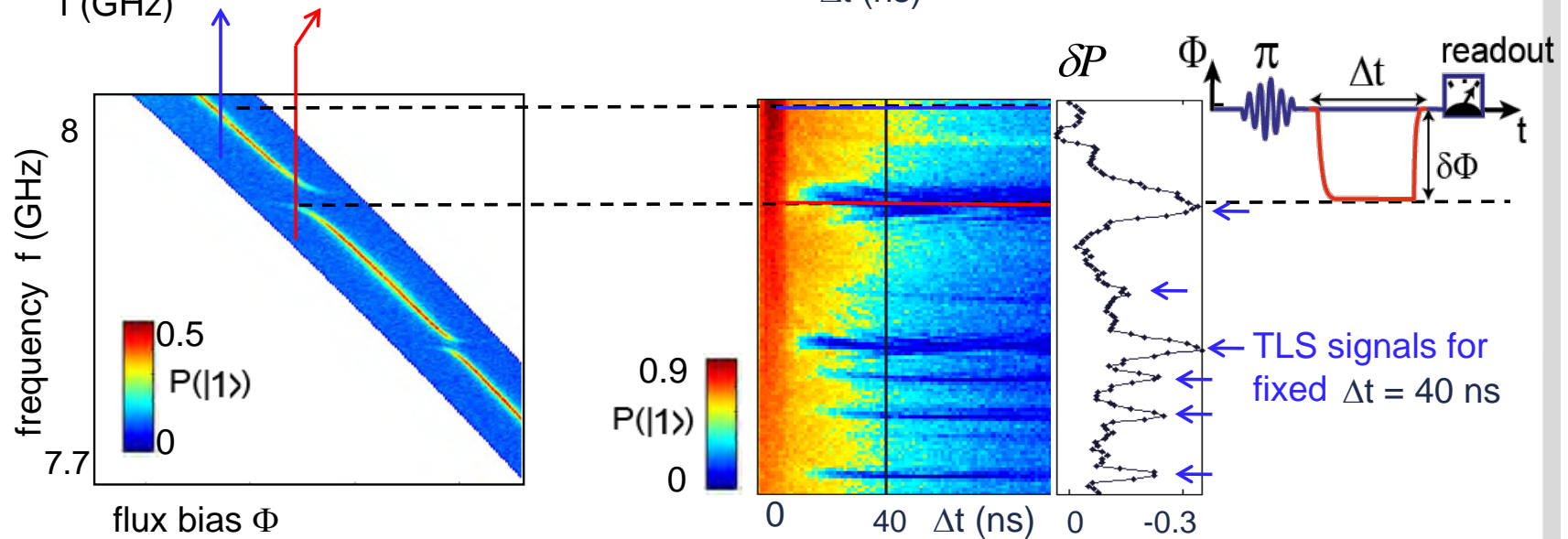
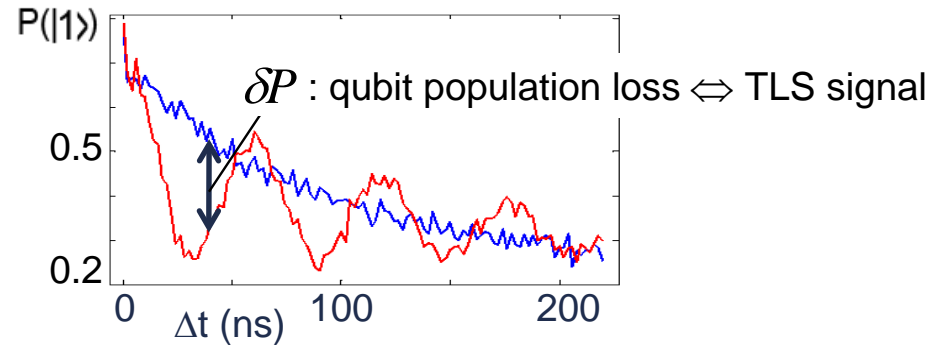
Frequency Domain:

defects cause avoided level crossings



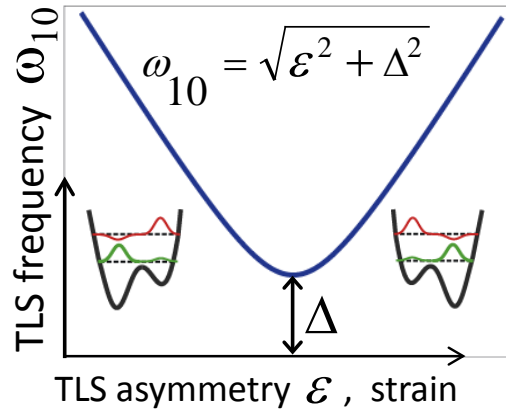
Time Domain:

energy oscillates between qubit and defects

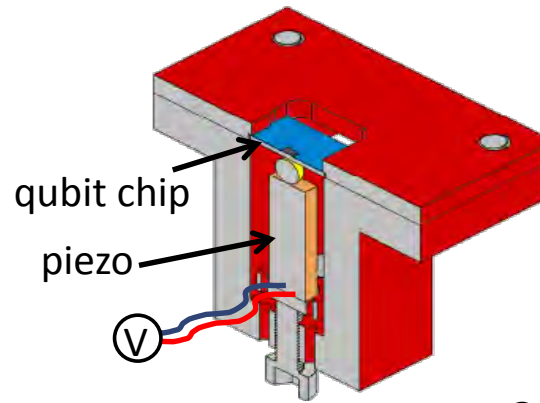


TLS Strain Spectroscopy

G. Grabovskij, J. Lisenfeld, G. Weiss, A.V. Ustinov et al., *Science* **338**, 232 (2012)



TLS strain tuning by deforming the sample using a piezo



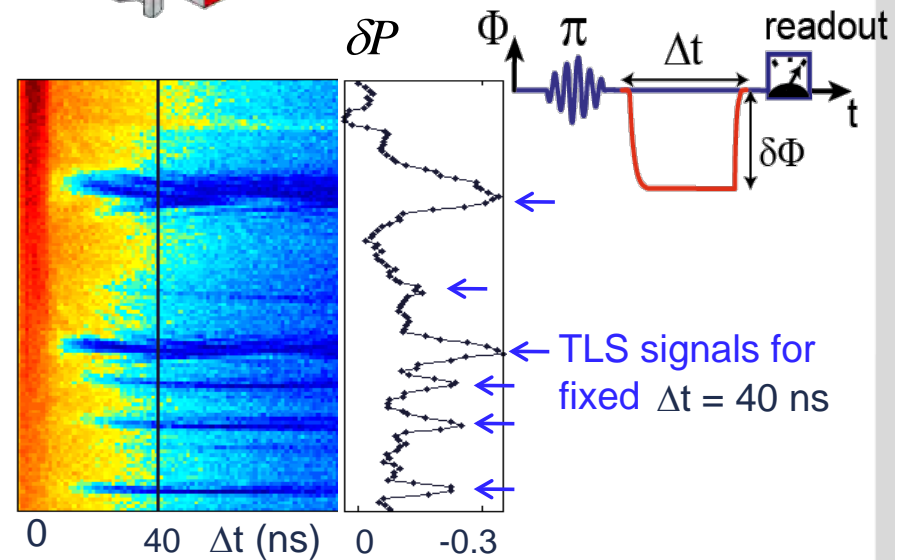
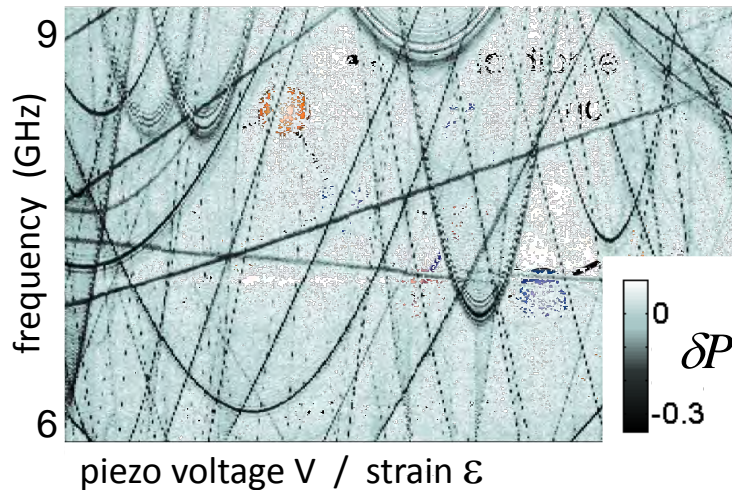
- tiny deformations

$$\frac{\Delta L}{L} \approx 10^{-7} / V$$

(compress 1 nm by 10^{-16} m)

change TLS asymmetry:

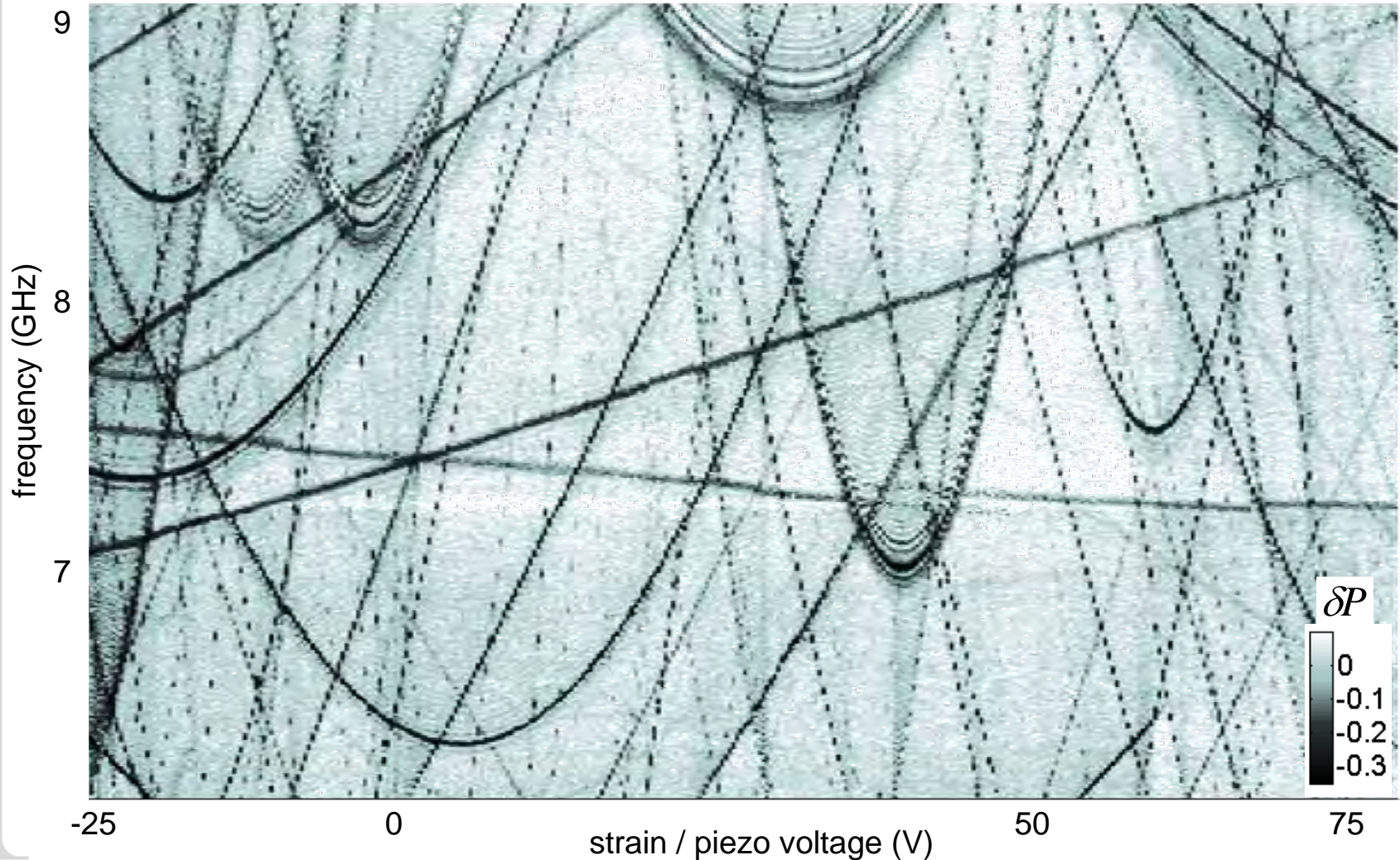
$$\varepsilon \approx 200 \text{ MHz/V}$$



TLS Strain Spectroscopy

G. Grabovskij, J. Lisenfeld, G. Weiss, A.V. Ustinov et al., *Science* **338**, 232 (2012)

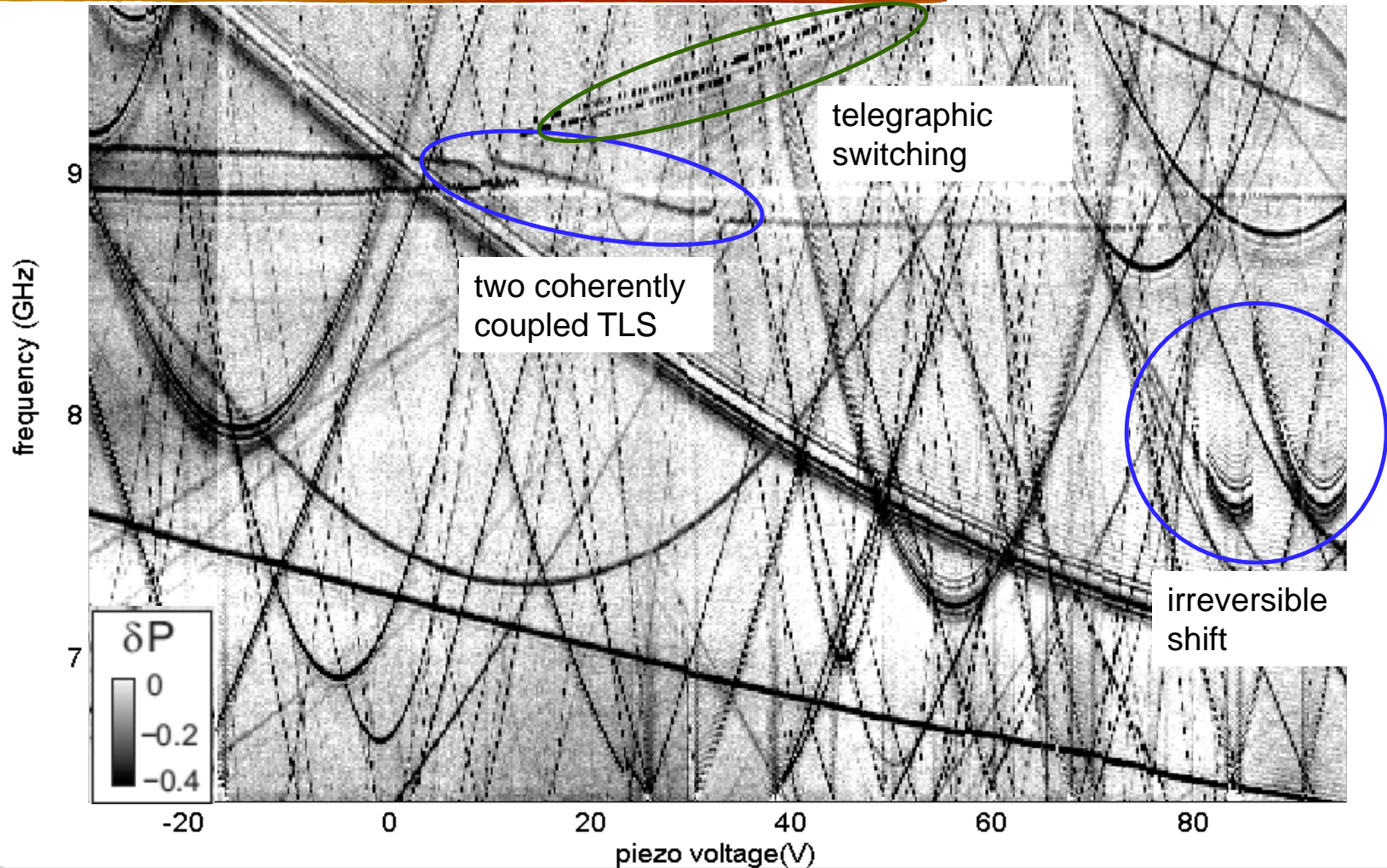
J. Lisenfeld, G. Grabovskij, J. Cole, C. Müller, G. Weiss, and A.V. Ustinov, *Nature Commun.* **6**, 6182 (2015)



TLS Strain Spectroscopy

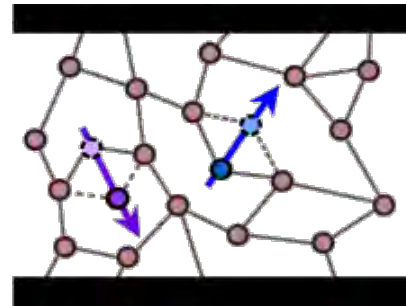
G. Grabovskij, J. Lisenfeld, G. Weiss, A.V. Ustinov et al., *Science* **338**, 232 (2012)

J. Lisenfeld, G. Grabovskij, J. Cole, C. Müller, G. Weiss, and A.V. Ustinov, *Nature Commun.* **6**, 6182 (2015)

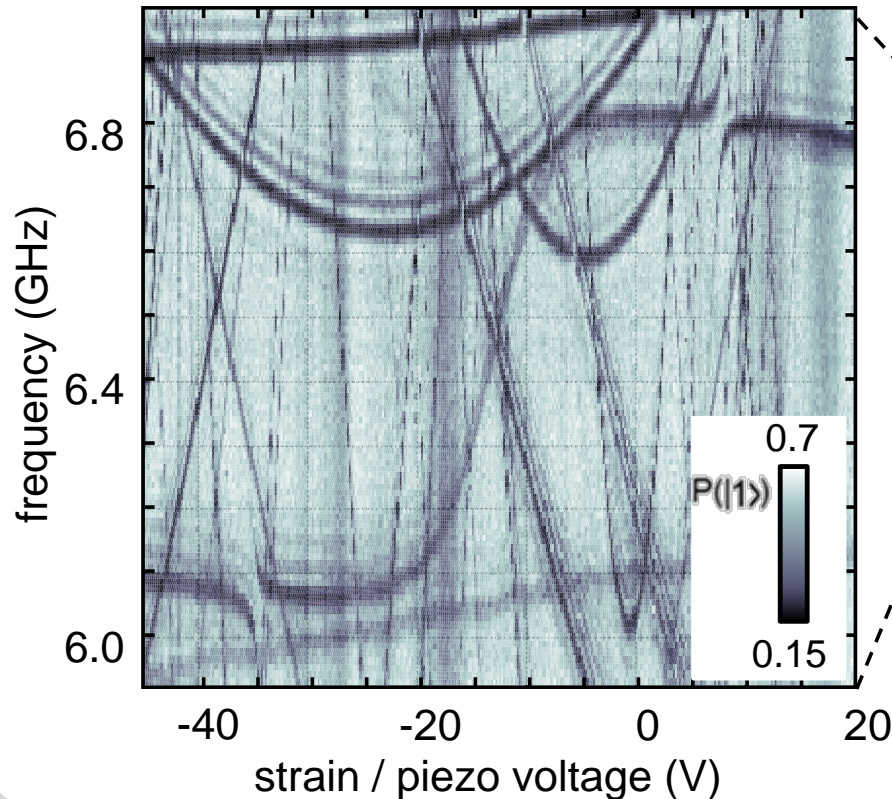


coherently interacting defects

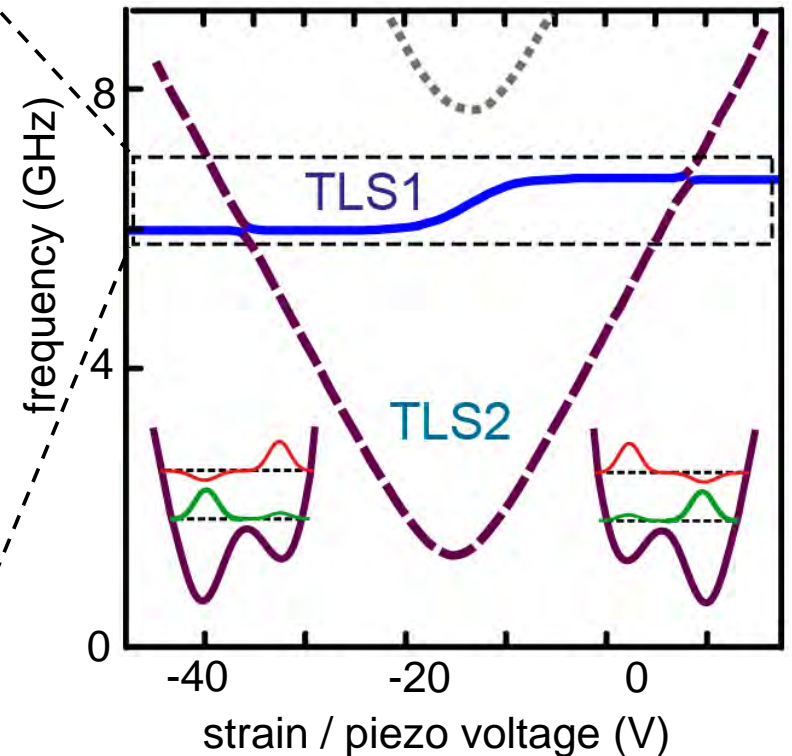
J. Lisenfeld, G. Grabovskij, C. Müller, J.H. Cole, G. Weiss,
and A.V. Ustinov, Nature Communications **6**, 6182 (2015)



signature in defect spectroscopy



simulated spectrum of 2 coupled TLSs



coherently interacting defects

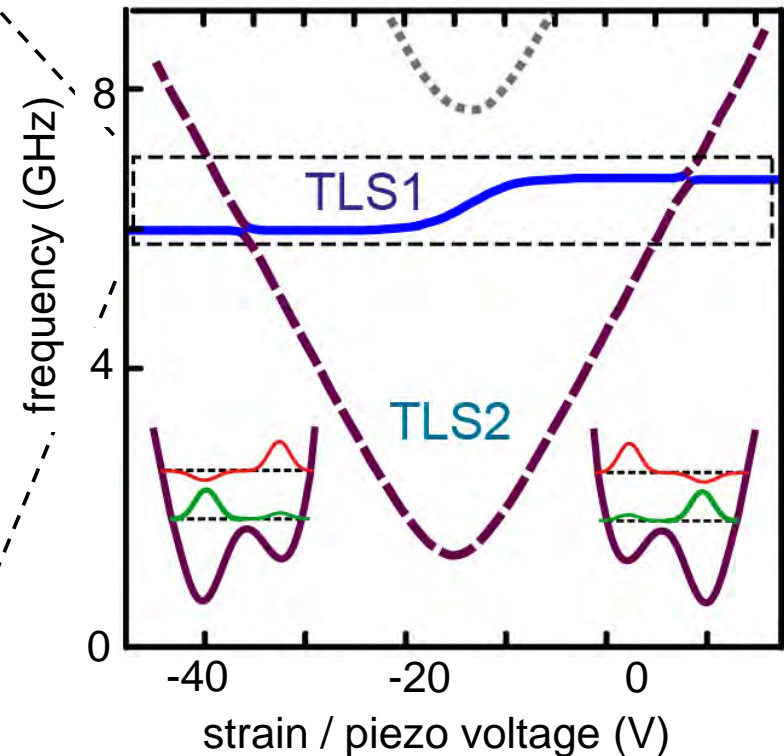
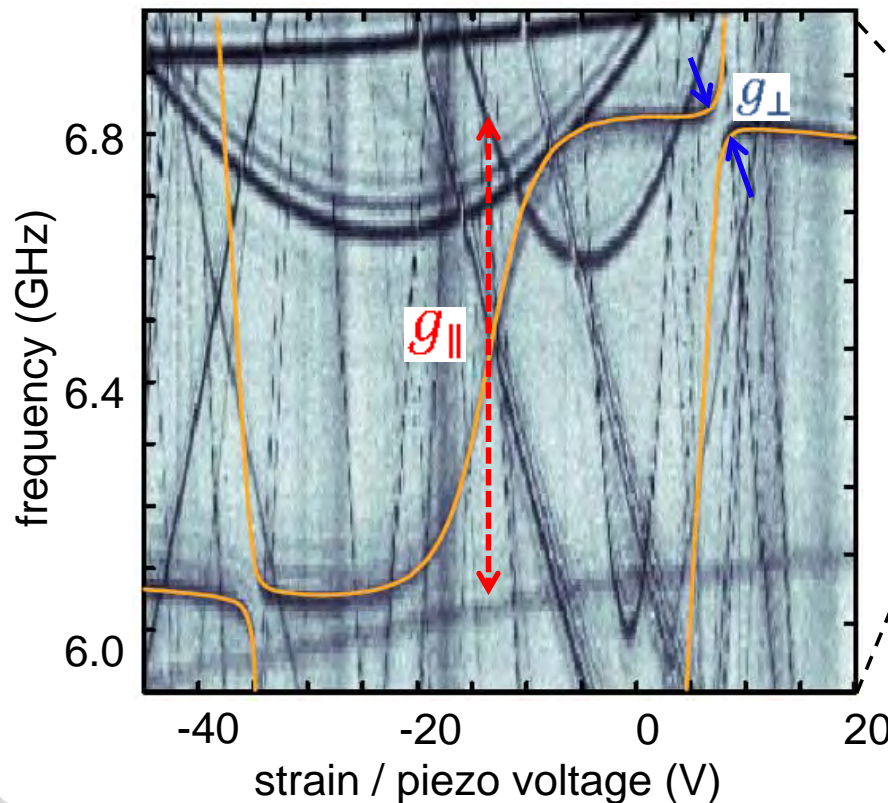
J. Lisenfeld, G. Weiss, A.V. Ustinov et al.,
 Nature Communications **6**, 6182 (2015)



interaction Hamiltonian:

$$H_{\text{int}} = g \sigma_{z_1} \sigma_{z_2} \quad \text{rotate to eigenbasis by angle } \theta, \text{ where } \cos \theta = \frac{\varepsilon}{E}, \quad \sin \theta = \frac{\Delta}{E}$$

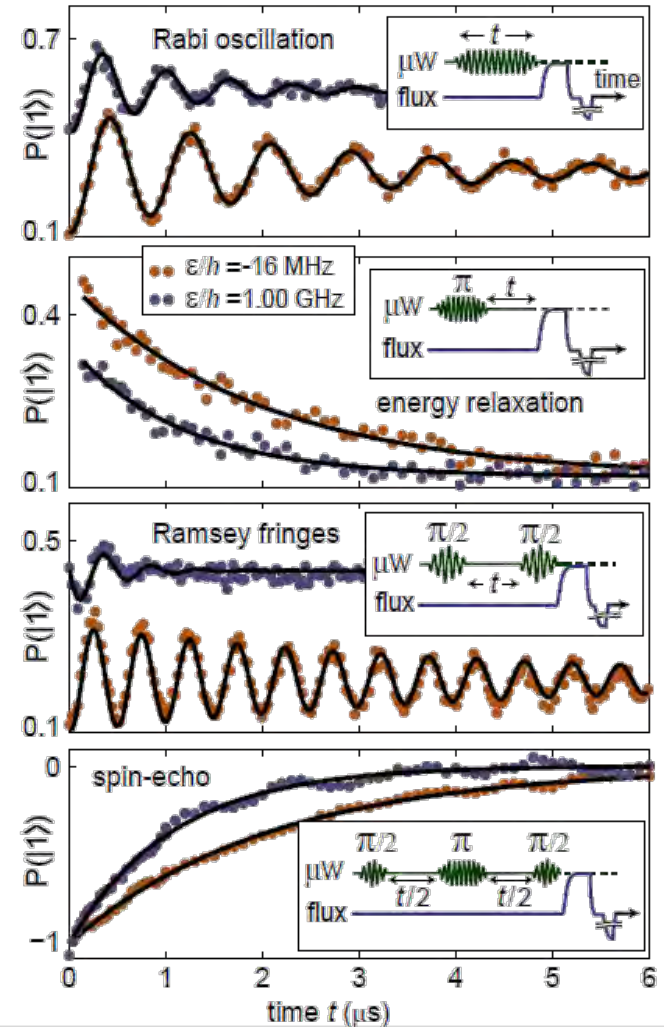
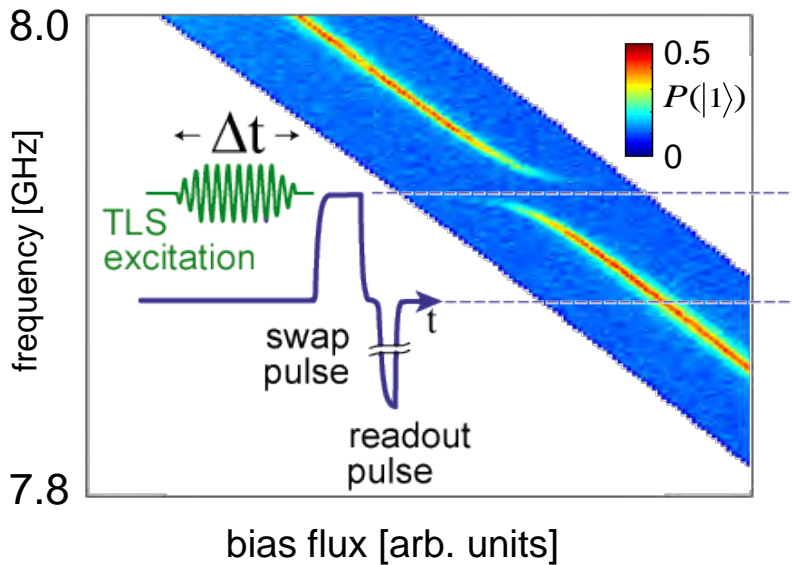
$$\rightarrow \hat{H}_{\text{int}} = \underbrace{g \cos \theta_1 \cos \theta_2}_{g_{\parallel}} \hat{\sigma}_{z_1} \hat{\sigma}_{z_2} + \underbrace{g \sin \theta_1 \sin \theta_2}_{g_{\perp}} \hat{\sigma}_{x_1} \hat{\sigma}_{x_2} + \underbrace{\propto (\hat{\sigma}_x \hat{\sigma}_z + \hat{\sigma}_z \hat{\sigma}_x)}_{\text{minor contributions}}$$



coherent control of individual TLS

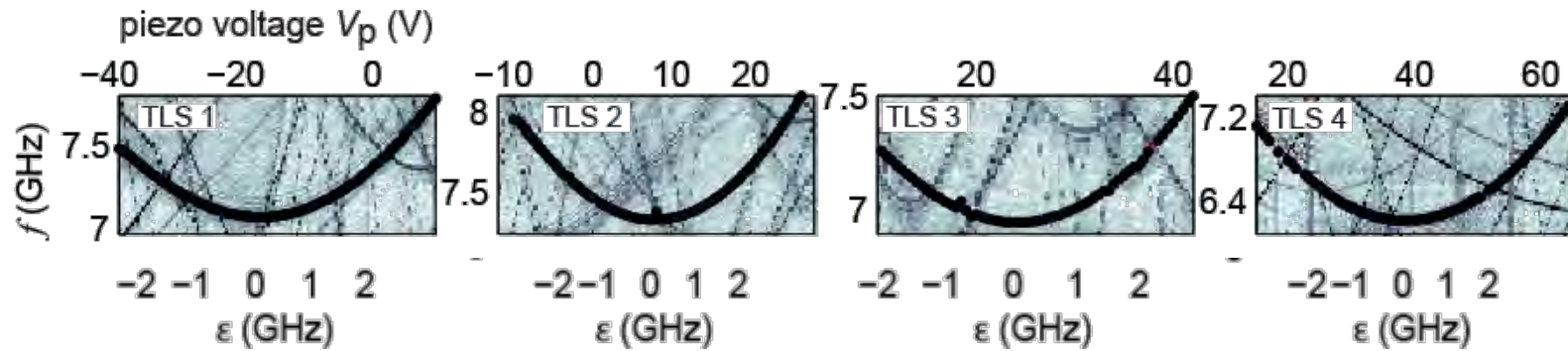
J. Lisenfeld et al., Phys. Rev. B **81**, 100511 (2010)

- TLS resonantly absorbs photons via a **Raman transition** involving a virtual qubit excitation
- ➔ full NMR-like TLS control via microwave pulse sequences
- ➔ TLS operation not degraded by qubit decoherence



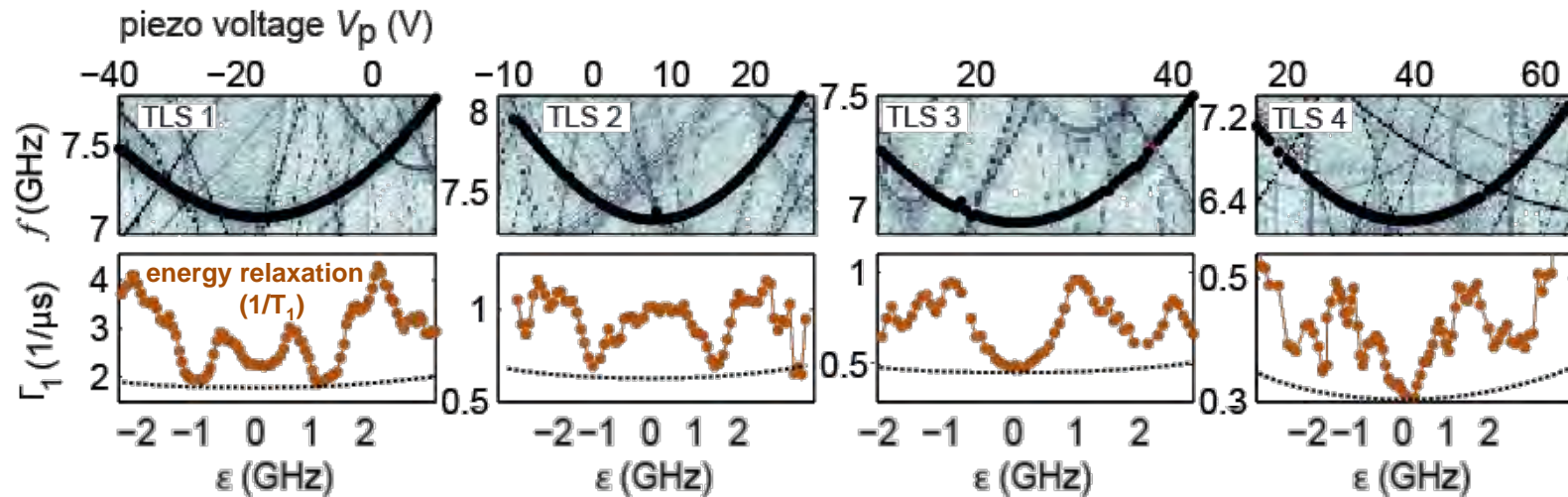
Strain-dependence of TLS coherence times

J. Lisenfeld, G. Schön, A. Shnirman, G. Weiss, A.V. Ustinov et al., Scientific Reports **6**, 23786 (2015)



Strain-dependence of TLS coherence times

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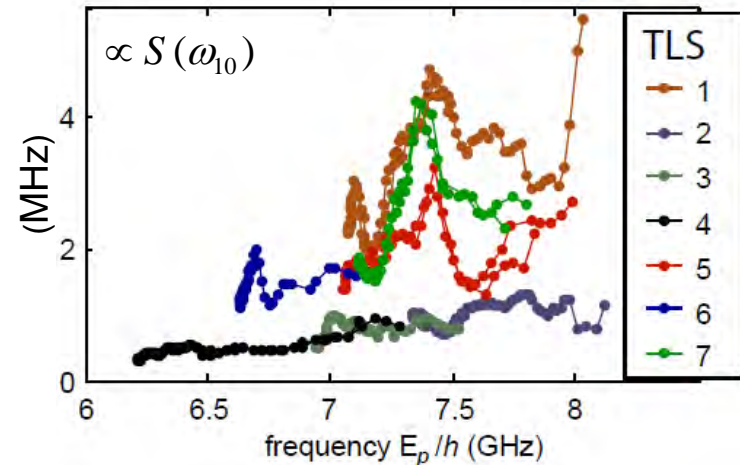


➔ symmetric pattern in Γ_1 can not originate in mutual TLS interactions

■ golden rule:
$$\Gamma_1 \propto \left(\frac{\Delta}{\sqrt{\Delta^2 + \epsilon^2}} \right)^2 \cdot S(\omega_{10})$$

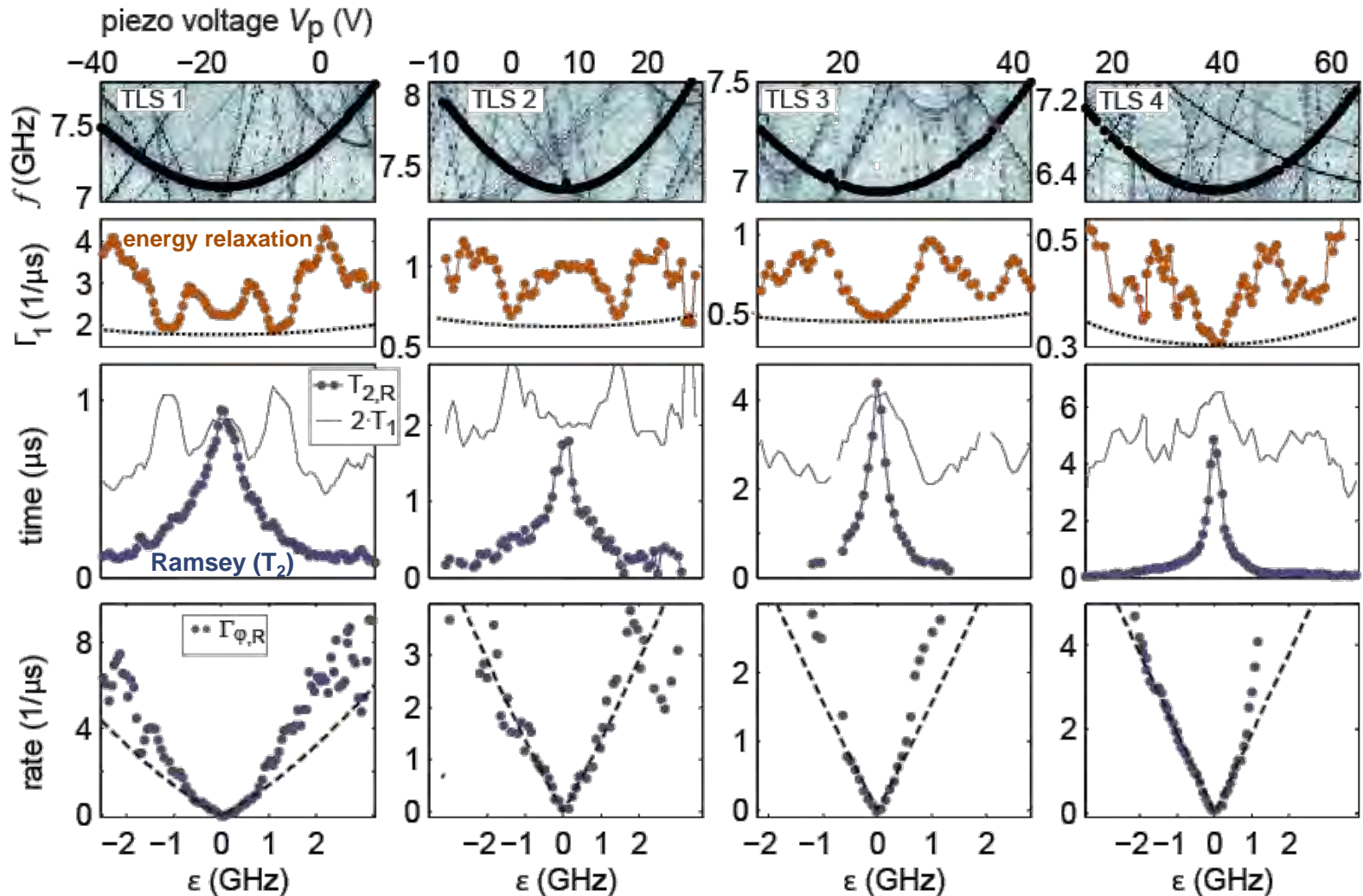
several TLS have a common maximum in Γ_1 around 7.4 GHz

➔ possibly coupling to same phonon mode



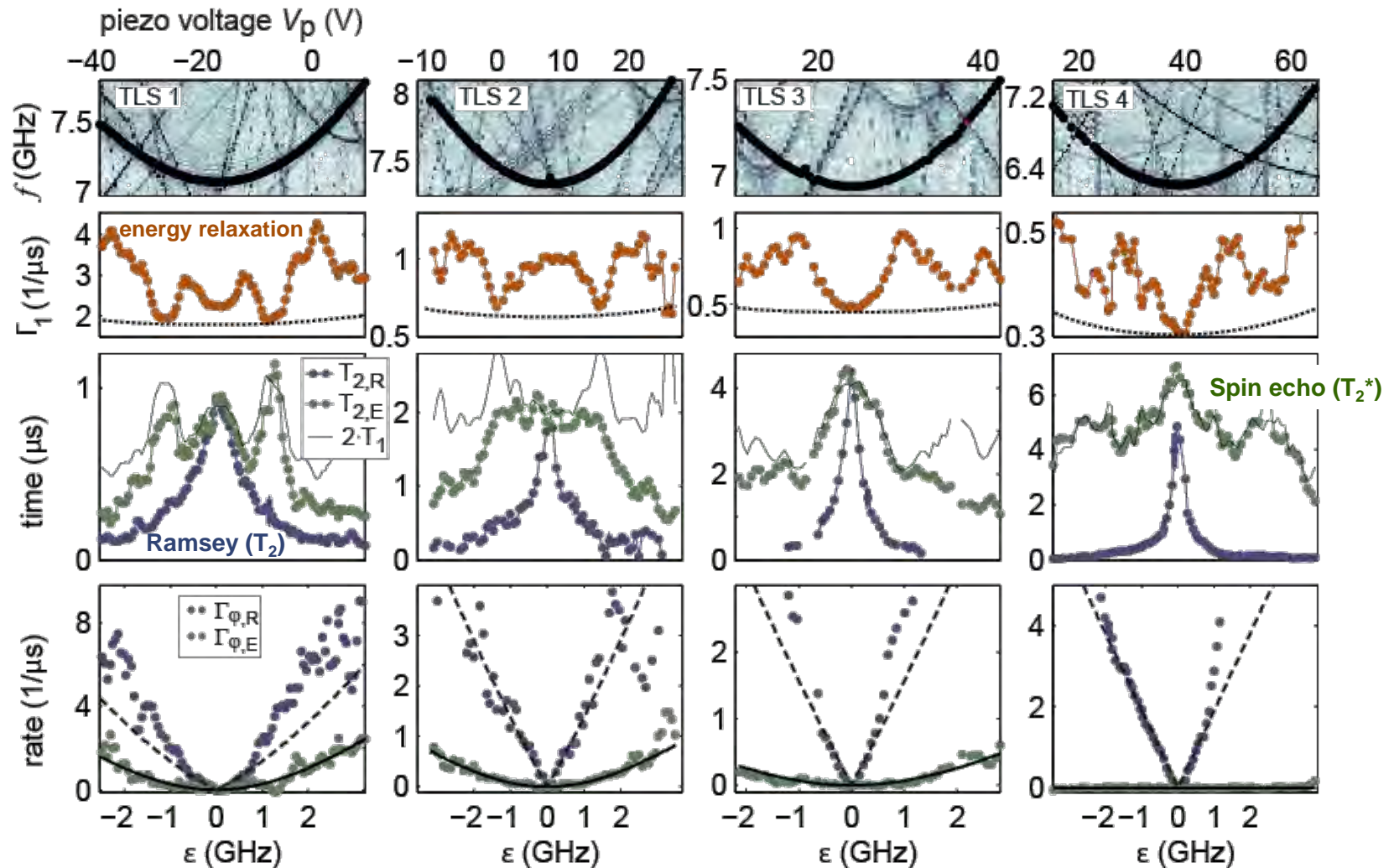
Strain-dependence of TLS coherence times

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Strain-dependence of TLS coherence times

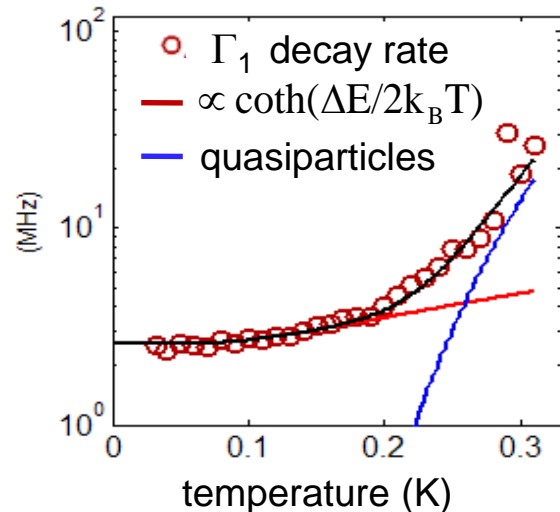
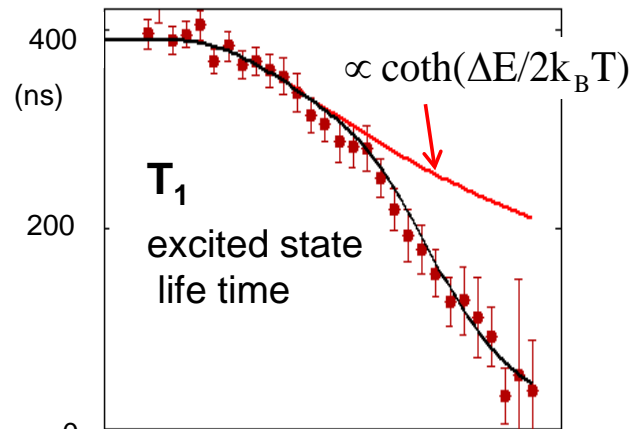
J. Lisenfeld, G. Schön, A. Shnirman, G. Weiss, A.V. Ustinov et al., Scientific Reports 6, 23786 (2015)



Temperature dependence of TLS coherence

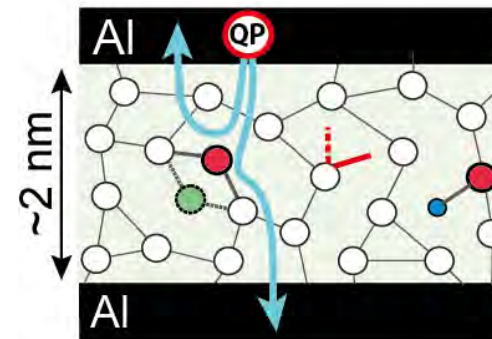
J. Lisenfeld et al., Phys. Rev. Lett. **105**, 230504 (2010)

TLS energy relaxation rate exceeds phonon contribution



TLS decay at higher temperatures [1] caused by quasiparticles ?

c.f. J. L. Black, Glassy Metals I, Topics in Applied Physics 46, 167 (1981).

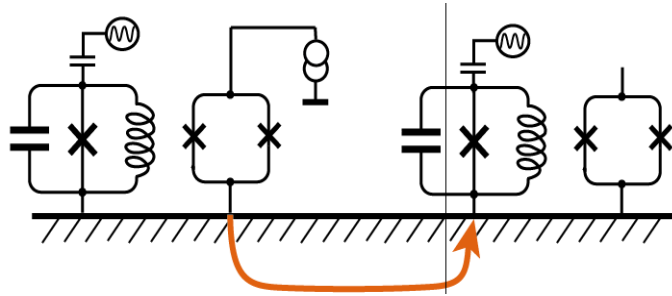


test :

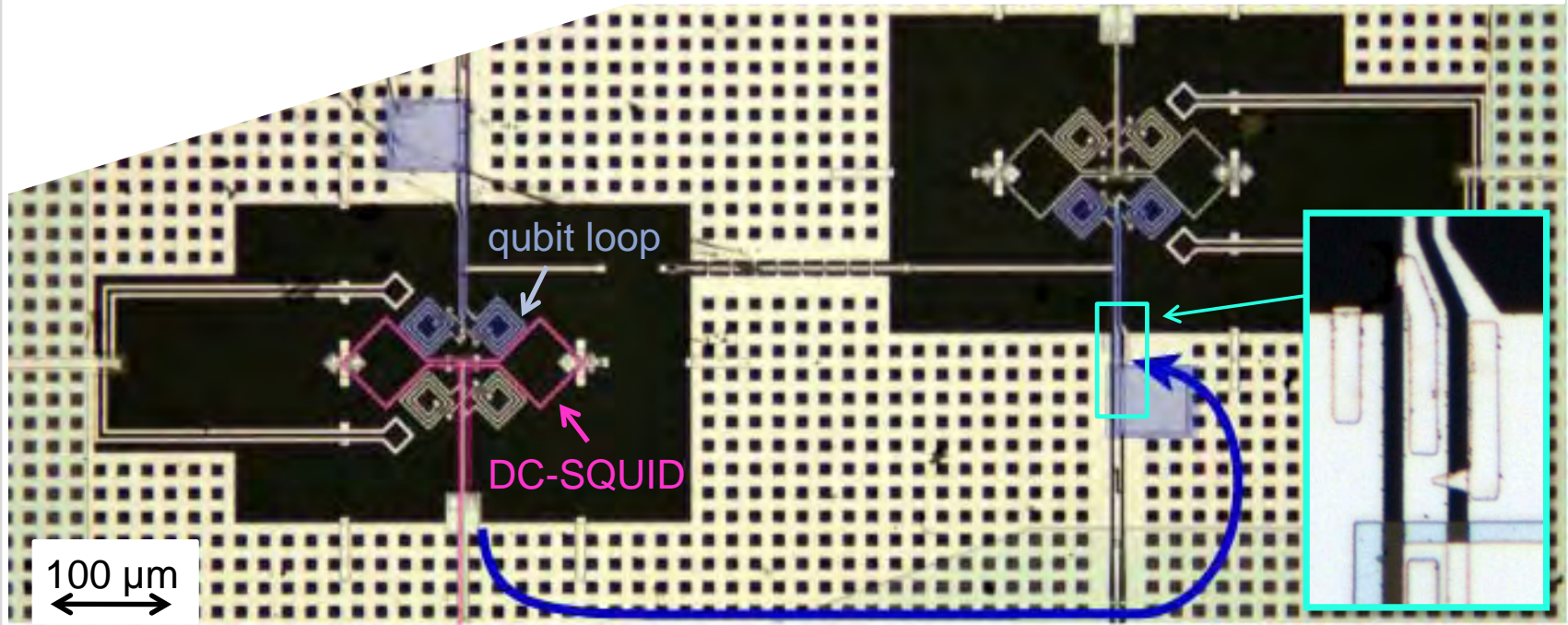
- 1) generate quasiparticles by injection or by raising the sample temperature
- 2) calibrate QP density using the qubit
- 3) measure TLS coherence times

injection of quasiparticles

cf. M. Lenander et al., PRB **84**, 024501 (2011)



- drive 2nd on-chip DC-SQUID with current $I_b > I_C$
- ➔ generated QPs diffuse to qubit's Josephson junction where they interact with TLS
- ➔ we expect a difference in QP density on the two JJ electrodes because of the sample layout



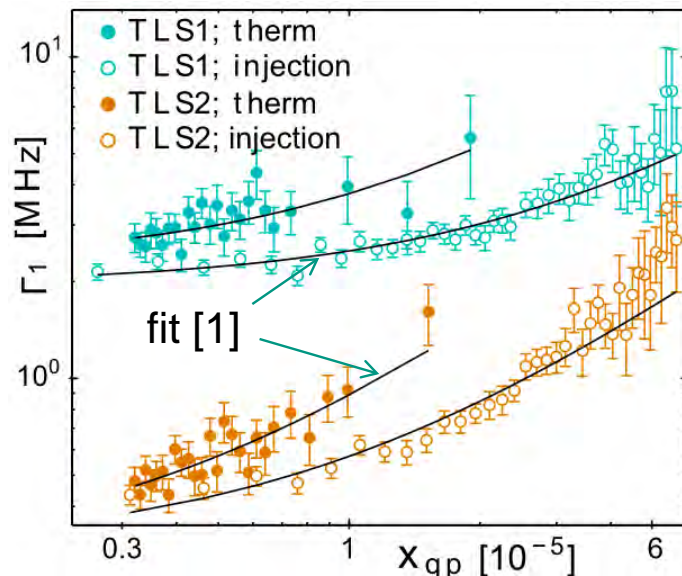
QP-induced decoherence of Two-Level Systems

A. Bilmes, J. Lisenfeld, G. Weiss, A.V. Ustinov et al., arXiv:1609.06173



- Korrington-like QP-TLS-interaction:

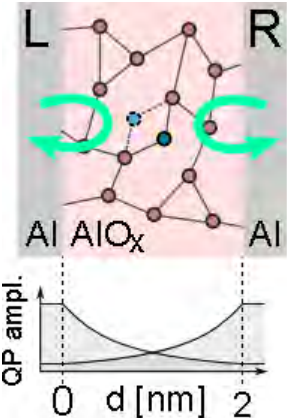
$$\Gamma_1 = S x_{qp} + \Gamma_1^{(0)}$$



- we observe:

$$\Gamma_1^{\text{therm}} > \Gamma_1^{\text{inject}}$$

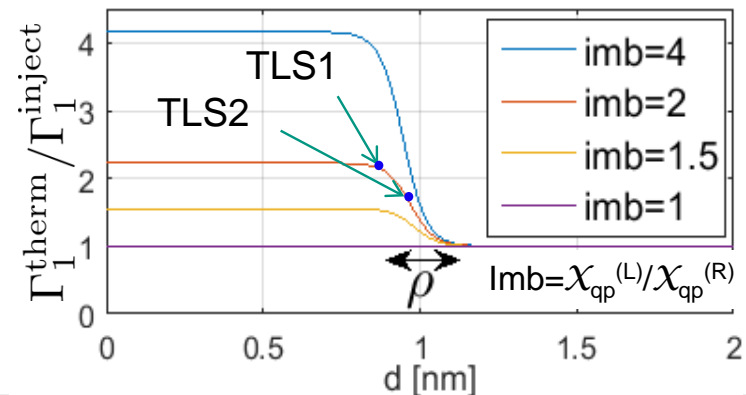
- QP-TLS coupling depends on TLS location



- QP densities differ in injection experiment:

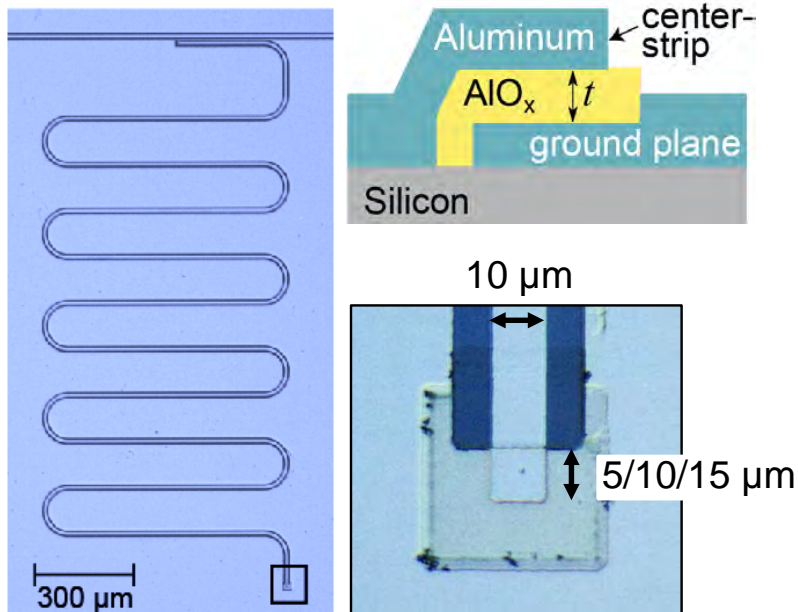
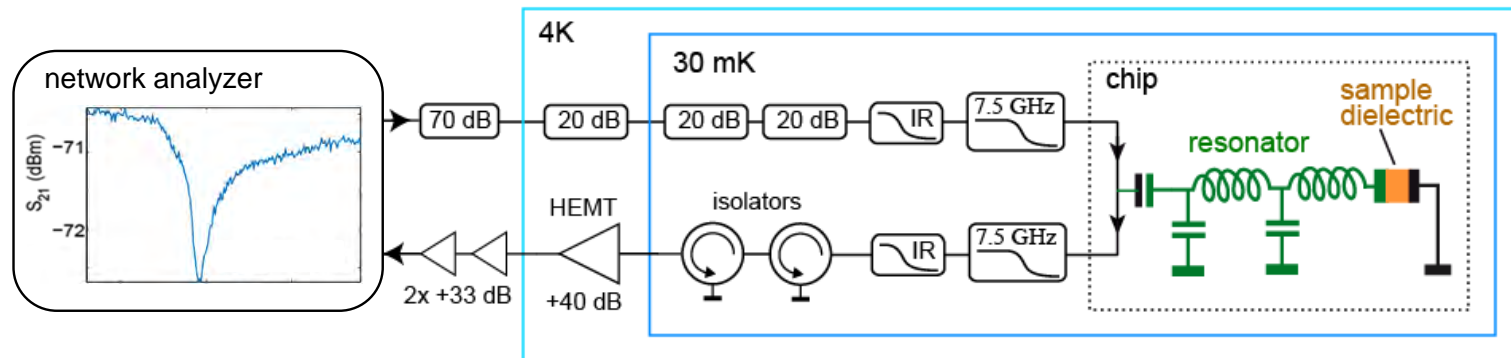
$$x_{qp}^{(L)} \approx x_{qp}^{(R)} / 2$$

- ➔ estimation of TLS location



coupling TLS to a transmission-line resonator

J. Brehm, A. Bilmes, J. Lisenfeld, to be published (2016)



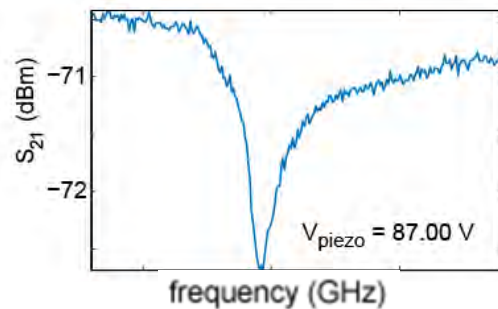
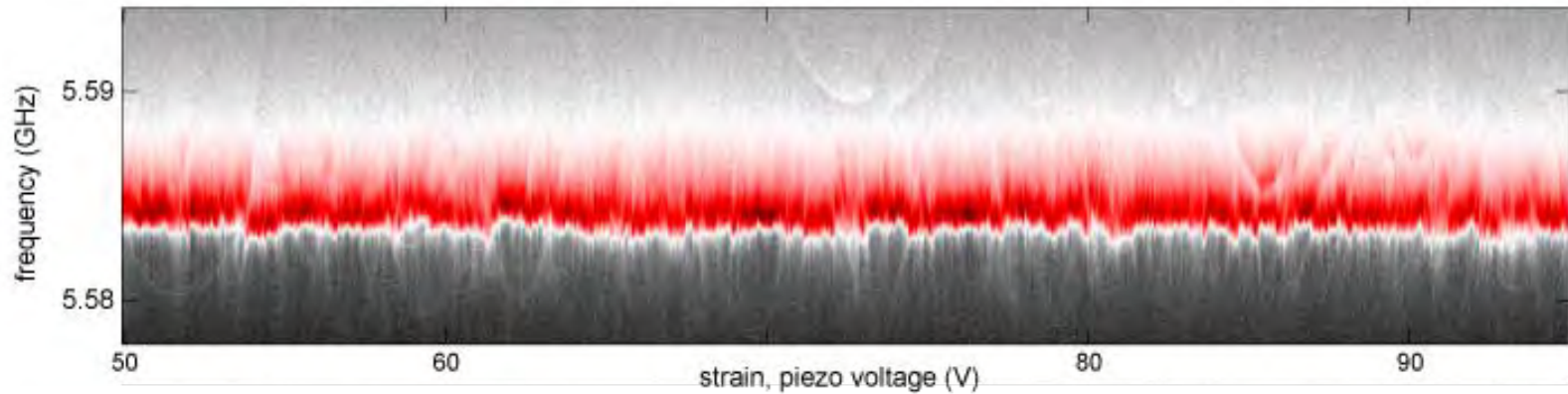
- AlOx fabricated using anodization
 $t = 50 \text{ nm}$, area $10 \mu\text{m} \times (5/10/15) \mu\text{m}$
- choose a dielectric volume to have
 ~ 1 TLS within 1-MHz-window
 $\rho \approx 100 \text{ TLS}/(\mu\text{m}^3 \cdot \text{GHz}) \Rightarrow V \approx 10 \mu\text{m}^3$
- coupling strength resonator-TLS g :

$$hg = p_{\parallel} |\vec{E}| \approx \frac{P_{\parallel}}{t} \sqrt{\frac{hf_{res}}{2(C_{res} + C_{cap})}}$$

$$C_{res} \approx 1.4 \text{ pF}, \quad C_{cap} \approx 0.2 \text{ pF}$$

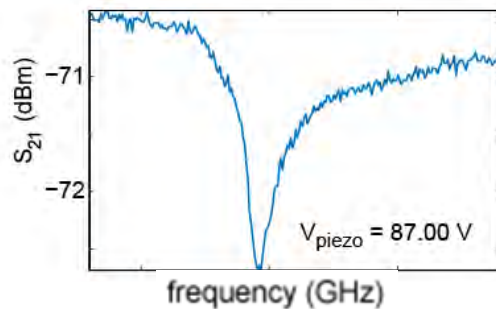
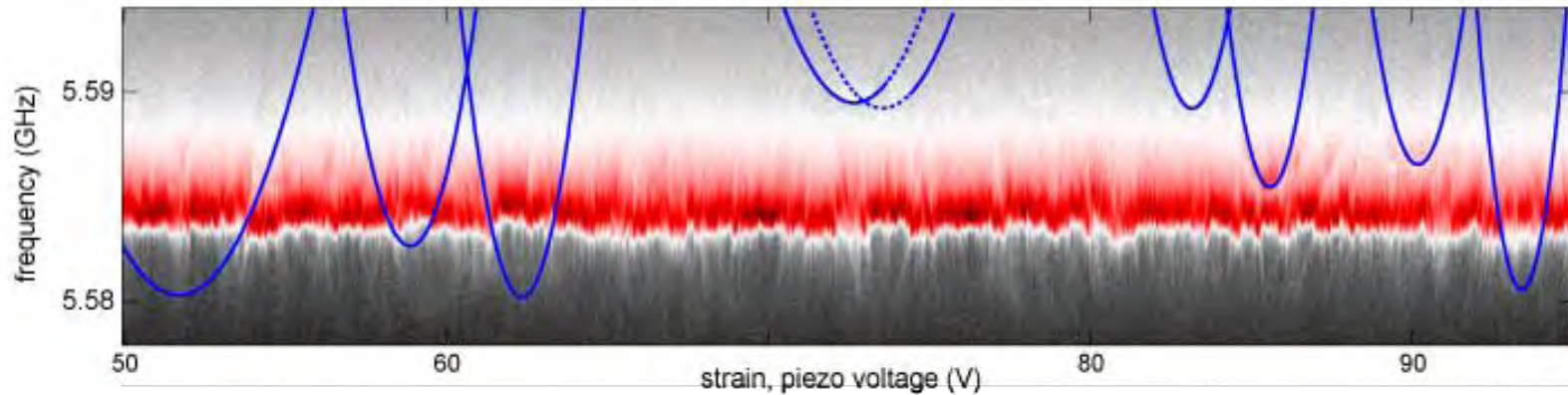
TLS coupled to a resonator: strain-spectroscopy

J. Brehm, A. Bilmes, J. Lisenfeld, to be published (2016)



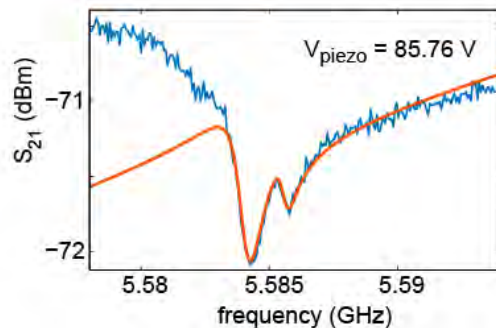
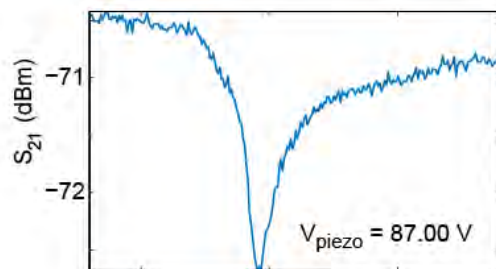
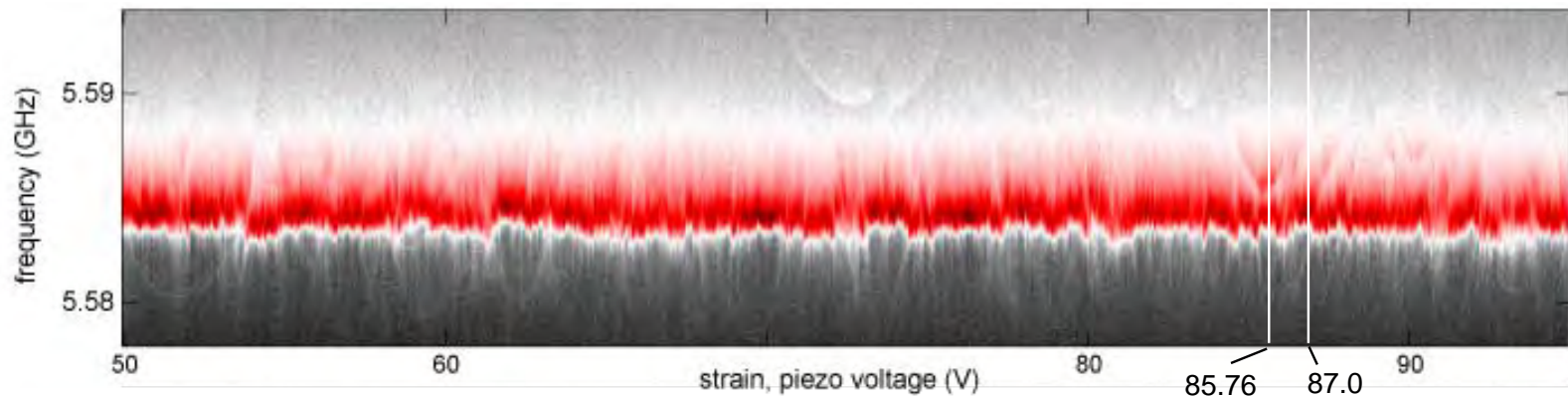
TLS coupled to a resonator: strain-spectroscopy

J. Brehm, A. Bilmes, J. Lisenfeld, to be published (2016)



TLS coupled to a resonator: strain-spectroscopy

J. Brehm, A. Bilmes, J. Lisenfeld, to be published (2016)



- individual TLS are resolved
- double-peaks in resonance with strongly coupled TLS

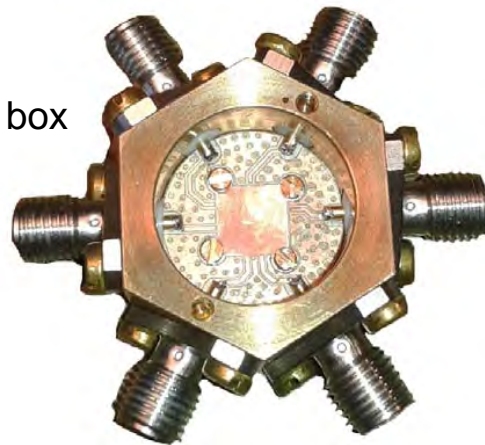
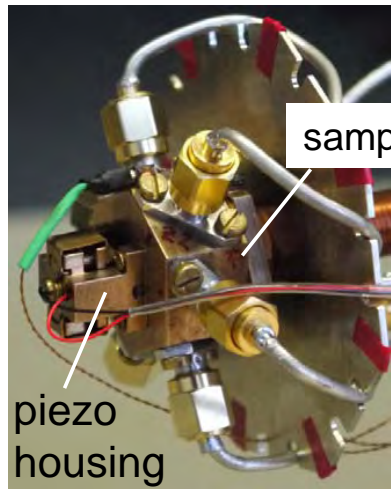
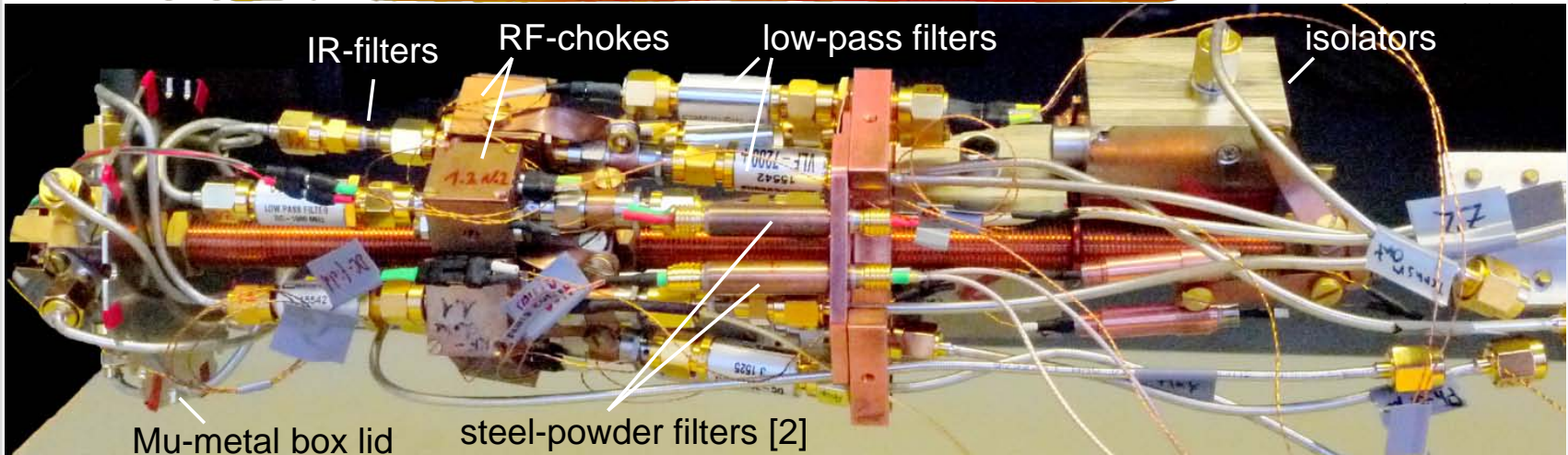
- fit to model provides TLS parameters [1]:

example:

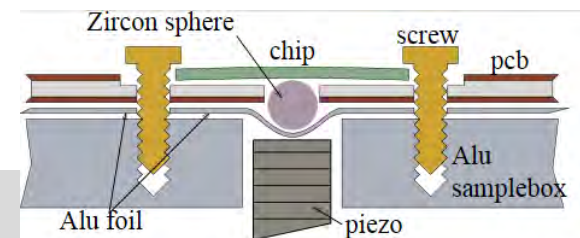
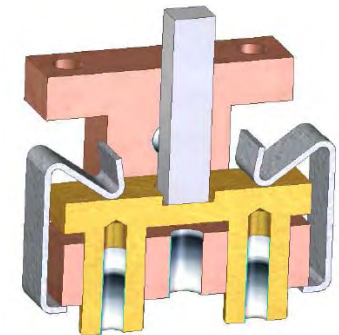
- $T_1 = 1.6 \mu\text{s}$, energy relaxation rate
- $g = 540 \text{ kHz}$, coupling strength
- $p_{\parallel} = 1.0 \text{ e}\text{\AA}$, dipole moment

[1] E. Rephaeli et al., 'Few-photon single-atom cavity QED with input-output formalism in Fock space, arXiv:1208.6053 (2012)

mechanical-strain tuning of TLS: setup^[1]



piezo housing

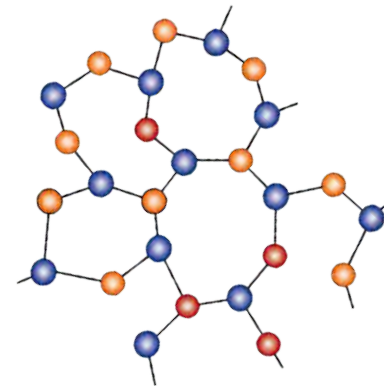


[1] design and assembly by A. Bilmes, PhD Thesis (KIT, 2018)

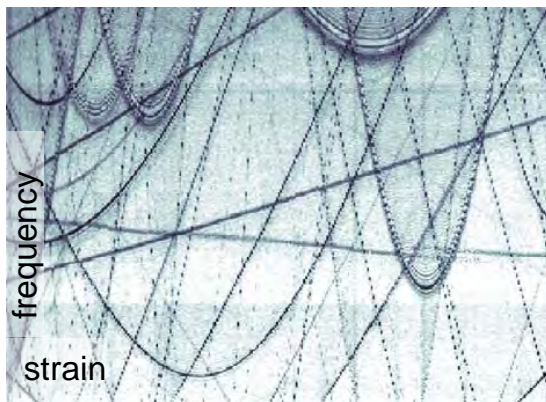
[2] A. Lukashenko and A.V. Ustinov, Rev. Sci. Instr. 79, 014701 (2008)

Summary: exploring TLS with superconducting Qubits

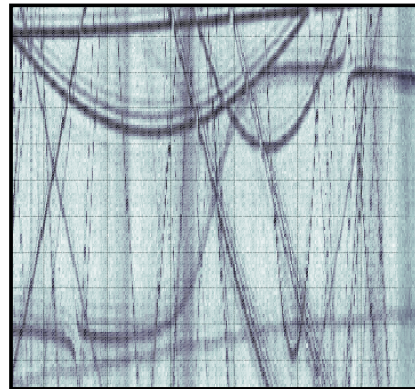
- **TLS** are a **major decoherence source** which affects various microfabricated devices
 - ➔ it is vitally important to understand emergence of TLS in fabrication
- it is now possible to **address single TLS coherently** with superconducting circuits



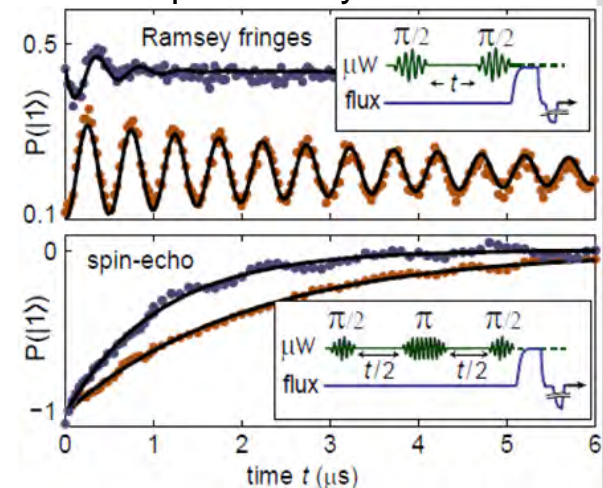
TLS strain spectroscopy



coherently coupled TLS



TLS quantum dynamics



- ➔ Superconducting qubits (and resonators) are **ideal tools** for the **characterization of materials** and **defect properties**.